# ANTENNAS FOR WV BAND APPLICATIONS 

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Electrical and Computer Engineering
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This dissertation is approved, and it is acceptable in quality and form for publication: Approved by the Dissertation Committee:

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# ANTENNAS FOR WV BAND APPLICATIONS 

## by

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## DISSERTATION

Submitted in Partial Fulfillment of the
Requirements for the Degree of

## Doctor of Philosophy in Engineering

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Albuquerque, New Mexico

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## Dedication

To my dear father, Nazem

My loving mother, Safiya

And my beloved sister, Amal

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# Antennas for WV Band Applications 

by

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#### Abstract

This dissertation focuses on designing, fabricating and testing antennas that are suitable for operation within the V/W bands. In particular, this work focuses on the design of slotted rectangular waveguide antenna arrays and cross slotted waveguide fed horn antennas. These structures are known for their high efficiency and high circularly polarized gain that can be implemented in satellite and terrestrial communication links. In addition, such designs can be implemented in radar applications that operate infrequency bands around 72 GHz or 84 GHz bands. Such antenna structures are inexpensive to fabricate since they can simply be machined using high precision conventional methods (like milling) and laser cutting when suitable.

In particular, this dissertation discusses two designs involving cross-slotted waveguides. The first one consists of an array of cross slotted rectangular waveguide antennas exhibiting radiation beams with Left Hand or Right Hand circular polarization. The array is composed of $8 \times 16$ elements and generates a circularly polarized gain of 25 dB over the frequency range $84.2-85.7 \mathrm{GHz}$. The array also exhibits a cross polarization


discrimination of more than 20 dB and an isolation of more than 20 dB between the feeding ports. A new type of z -arm shaped cross slots is introduced that fits on the broad-wall of a conventional rectangular waveguide. The feeding network is also optimized to regulate the power and the phase of each rectangular waveguide element in order to increase the gain of the antenna array.

This dissertation also presents the theoretical analysis of the second design which is a cross slotted waveguide polarizer. The polarizer achieves a RHCP or LHCP by altering the feeding ports. The polarizer feeds different conical horns or pyramidal horns without affecting their characteristics. The efficiency of the polarizer is improved by combining the power of different crossed slots using square waveguide combiners. Other modes of operation of the cross slot polarizer are investigated for multiband operation and an antenna system operating at 72 GHz and 84 GHz simultaneously is designed.

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## INTRODUCTION

## - Motivation

New communication systems that are able to support a large number of users and high data rates require researchers to investigate the use of millimeter-wave frequency bands for communication. These frequency bands are either unassigned, or are less crowded than the lower end of the frequency spectrum. This availability of these frequency bands ensures high bandwidths, which in turn ensure high data rates in communication links.

For years, the solution for the congestion of the lower end of the frequency spectrum was the use of techniques, such as cognitive radio, that allow systems to look for frequency bands that are free and use them for communication [1]. Another solution was the use of frequencies in the K and Ka -band ranges, which was made possible by the technological advancement of RF systems [2]. Nowadays, the use of higher frequencies in the W/V bands is under research and investigation.

In the millimeter-wave frequency bands, the free space propagation losses (FSPL) under normal weather conditions are significant for long distance communication. This loss factor is directly proportional to the distance of the link and inversely proportional to the wavelength. At these ranges, not all frequencies are suitable for long distance point-topoint communication links due to the atmospheric absorptions that adds to the losses. For instance, between $55-65 \mathrm{GHz}$ there is an oxygen absorption band that causes wave
attenuations of $10 \mathrm{~dB} / \mathrm{Km}$. At such frequencies, only short distance communication links are feasible [3].

For longer distance communication links, in the order of kilometers, increasing the frequency above 65 GHz does not present any challenges in terms of propagation other than the propagation losses. The $71-76 \mathrm{GHz}$ and $81-86 \mathrm{GHz}$ in the W/V band frequency ranges constitute a suitable candidate for such communication links. Propagation at these frequencies is being studied by the W/V band terrestrial link experiment (WTLE) located in Albuquerque, NM, USA [4]. These new bands are able to provide ultra-wide communication bandwidth and can be adopted for space, as well as terrestrial communications [5],[6]. Recently, these frequency bands were assigned by the Federal Communications Commission (FCC) to satellite communication and other applications that do not interfere with this type of communication [7]. The wave propagation at these frequencies creates some problems under severe weather conditions, which could be an obstacle to satellite or terrestrial communication. This dissertation presents research that aims to tackle these problems by designing inexpensive antennas and microwave devices that can be used for communication or radar systems at these high frequencies.

## APPLICATIONS

The W/V frequency bands have short wavelengths that make the wave more sensitive to small objects and more details in objects. This property is very attractive to high resolution radars. Multiple radar systems are being deployed at these frequencies. A monopulse radar system at 94 GHz was developed for high resolution tracking in antimissile systems or in satellite communication tracking [8]. Another radar system
targeting the detection of airport runway debris was developed at 94.5 GHz [9]. Other radar systems at 77 GHz were developed for car collision avoidance systems [10],[11],[12]. This type of radars can be used for both long range and medium range car detection as shown in Figure I.1.

(b)

Figure I.1: (a) The traditional two radar systems used for long range and medium range detection,
(b) the combined radar system used for both detection types.[10]

Another type of applications developed at W-band frequencies is passive millimeterwave imaging [13]. This application takes advantage of the high resolution of the wavelength, around 3 mm , to detect objects that are unseen in bad visibility conditions. This technique can be used for both military and civilian purposes, such as detecting
concealed weapons, rescue missions under low visibility, low visibility navigation and oil spill detection. An example of these applications that use imaging is shown in Figure I.2.


Figure I.2: Imaging used to detect concealed weapons. [13]

Finally, due to the availability of large frequency bandwidths at W/V bands, high data rates communication links are possible. One potential application at 73 GHz is the 5 G cellular network [14],[15]. This application is expected to deliver $6 \mathrm{~Gb} / \mathrm{s}$ data rates at multiple frequency bands including a frequency band centered at 73 GHz . This increase in data rates is necessary due to the increased use of cellular networks and the increased demand of higher data rates by both the customers and the applications.

Indoor high data rates communication links, reaching Gbps speeds, were also developed [16],[17],[18]. These applications operate at $60,87.5$ and 94 GHz , producing 3
-25 Gbps data rates. A similar system is shown in Figure I.3. These high data rates ensure a better quality of service to the customer. These systems especially the one at 60 GHz , are used in short communication links.


Figure I.3: High data rate indoor wireless communication system. [16]

Since the frequency bands $71-76 \mathrm{GHz}$ and $81-86 \mathrm{GHz}$ are shown to be good candidates for long distance communications links, satellite communication and earth to satellite communication is another application at these frequency bands. Recently, the FCC has declared these frequency bands to be reserved for satellite communications and any other application that does not interfere with it [7]. Several studies are being conducted to check the feasibility of such communication systems [19],[20],[21], and a new satellite mission called WAVE is planned to be deployed in the W-band [20].

For the past two years, the University of New Mexico, in collaboration with the Air Force research Lab, have been studying the wave propagation at $\mathrm{W} / \mathrm{V}$ bands under the WTLE experiment, as a first step in the study of the feasibility of satellite communication
at these frequencies. In the next section a more detailed explanation of this experiment is presented.

## - WTLE EXPERIMENT

The W/V band Terrestrial Link Experiment (WTLE), is a combined research effort between the University of New Mexico, the Air Force Research Labs and NASA, aiming to study the wave propagation at $71-76 \mathrm{GHz}$ and $81-86 \mathrm{GHz}$ [22].

The transmitter is located at Sandia Peak, and the receiver is located on the roof top of the COSMIAC building in Albuquerque. The transmitter consists of a coherent 72 and 84 GHz continuous wave (CW) beacon with an EIRP of 40 dBm using two lens antennas with a $3^{\circ}$ half-power beam-width (HPBW) and 35 dBi of gain. The receiver has two 0.5 m Cassegrain reflector antennas with 52 dBi of gain each with one receiving at 72 GHz and the other at 84 GHz as shown in Figure I.4. The transmitter and receiver electronics are temperature controlled so that their power performance does not fluctuate. The receiver has also a tone injected right before the LNA to monitor the overall gain performance of the receiving system [23].

The receiver has a dynamic range of 70 dB at V-band and 68 dB at W -band. The polarization of both transmitter and receiver is LHCP. The receiver has a cross polarization discrimination (XPD) of 13 dB at V-band and 20 dB at W -band. The power measurements are taken at a 10 Hz rate.

The link extends over a 24 Km with a $3.9^{\circ}$ slant angle as shown in Figure I.4. The slant angle is aimed to represent the direction of arrivals at which satellite communication takes
place. The link provides a dynamic range of 70 dB that allows it to study the attenuation caused by the different weather conditions such as rain, snow, fog, haze and sand storms.

Multiple atmospheric measurement devices are also used to monitor the weather conditions. At the transmitter side, a disdrometer is installed to measure the rain drop size distribution and velocity and a weather station is installed with a wind monitor, temperature, pressure, and relative humidity sensors. At the receiver side, the same equipment used at the transmitter side are installed in addition to a radiometer that gives an estimate of the attenuation experienced. Other systems are used to monitor the weather conditions over the link path. One such system is a NEXRAD that allows locating the areas where the weather events are taking place and gives a better understanding of the condition of the link path. Another device that is used is a SODAR in multiple locations under the link path. All these instruments give a better image of the weather conditions that are affecting the wave propagation and help better correlating the weather conditions with the measurements of power and polarization at the receiver side.

Preliminary results show that in heavy rain events with a rain rate of $17.2 \mathrm{~mm} / \mathrm{hr}$ at the receiver side, the attenuation of the wave exceeds the dynamic range of the system on all channels. No significant depolarization is observed during such rain event. However, in a snow event, wave attenuation is observed, but more significantly, wave depolarization takes place for both V and W bands. At W band the wave depolarization is more significant, where the X-polarization power level exceeds the co-polarization power level. At that particular event Albuquerque experienced 7.6 mm of snow. Another finding of the experiment shows that clouds have a peak 12-15 dB attenuation effect on the wave but no depolarization results from that.


Figure I. 4 The transmitter on top of Sandia Peak (left); the V and W-band receivers in COSMIAC Albuquerque, NM (right); and an illustration of the elevation profile of the slant path covered by the link (bottom) [23].

## - Dissertation Goals

The preliminary results shown in the previous section imply that in order to have a good communication link at W and V bands, it is critical to have a high dynamic range for the system being employed and a high cross polarization discrimination. These two factors will guarantee not interrupting the communication link between a ground station and a satellite, or between two ground stations for more than $0.01 \%$ of the time. These two factors could
be achieved by designing antennas that have high gains and high isolation between the Copolarization component and the cross-polarization component.

Several antenna designs are illustrated in this dissertation that could be used in communication links or radar systems at 72 and 84 GHz .

## LITERATURE REVIEW

Few types of high gain circularly polarized antennas are designed to operate at W-band frequencies. Some of the commercially available antennas include conical horns, reflector and lens antennas. In reality, these antennas are not circularly polarized by nature. They require the integration of a polarizer unit for circular polarization excitation [24],[25].

In this section, a literature review of the different types of antennas and polarization techniques that could be designed at W/V band to operate in communication links is done. The literature review targets the design of high gain antennas, or the design of microwave polarizers that transform the linear polarization of high gain antennas into a circular polarization.

## - Slotted Rectangular Waveguides

Slotted rectangular waveguide antennas are known for their high power handling, high efficiency and their ability to produce high gain while having a flat small size. These types of antennas can be used in the transmitter side of a satellite communication link. There are two types of slotted waveguide antennas: leaky wave arrays and standing wave arrays. The standing wave arrays are narrow band, which makes them a less favorable option in communication links that are trying to achieve high data rates. The leaky wave configuration is the better option and it is the one discussed in the text to follow herein.

Circular polarization of slotted waveguides is achieved by appropriately designing slotted rectangular waveguide arrays. Such structures can be implemented by having tilted
slots along the waveguide centerline on the narrow-wall [26], [27]. However, the implementation of rectangular slots along the waveguide narrow walls is critical at W bands since the slot's length is usually larger than the wall width. This requires the slots to extend to the broad-wall and thus leading to a difficulty in fabrication and a degradation in the structure radiation performance.

Another approach in the design of circularly polarized waveguide arrays is to place the slots on the broad-wall. For example, by adding parasitic dipoles on top of the rectangular slots of a linearly polarized waveguide array, the polarization can be switched to circular [28], [29]. This technique reduces the efficiency of the antenna due to the use of a dielectric material where the parasitic dipoles are fabricated as shown in Figure II. 1 (a) and (b).

The addition of tilted, round edge rectangular shaped cavities around the rectangular slots, switches the array polarization from linear to circular [30]. This technique adds more complexity on the fabrication of such arrays, especially at W/V band, making the entire design more expensive to implement.

Rectangular tilted slots that are located on both sides of the waveguide centerline, in an alternating order as shown in Figure II.2, can also generate circular polarization [31], [32]. This type of arrays needs to be terminated at the other end of the rectangular waveguide in order to provide a good axial ratio. By replacing the conventional rectangular waveguide with a ridge gap waveguide as shown in Figure II.2, the bandwidth performance of this type of arrays is improved as reported in [33]. However, such type of arrays is complex to design and fabricate, which may result in some deterioration in performance especially at W-band.

(a)

(b)

Figure II.1: (a) The configuration of the dipole layer on top of the slot, (b) fabricated prototype of the parasitic dipoles on top of a slotted waveguide array.[29]


Figure II. 2 : Tilted alternating slots with a ridge waveguide. [33]

One of the popular and convenient methods to achieve circular polarization in slotted rectangular waveguide arrays is through the integration of cross-slots. Such cross-slots can be located off the centerline of the broad-wall [34]. This type of circularly polarized arrays have been studied for frequencies up to Ka band $(26.5-40 \mathrm{GHz})$ and commercially adopted for satellite to earth communication systems [35] - [39] as shown in Figure II.3. These arrays can have different feeding configurations. More specifically, if the array has two separate feeding ports, then it leads to either LHCP or RHCP radiations with a change in the main beam direction [35]. But the feeding network used in [35], has led to a low isolation (10-12 dB) between ports 1 and 2 which is one of the causes that lowered the efficiency of the system by $6-10 \%$.

By adding tilted slot pairs on the other side of the centerline, the beam direction can be corrected to be perpendicular to the array [37], [38]. Also, if the array is fed through one end and the other end of the structure is shorted, then either RHCP or LHCP can be achieved depending on the placement of the cross-slotted elements [36], [39]. Under this last scenario, the reflection from the shorting plate affects the polarization by increasing the array's axial ratio. One possible solution is to add matching slot pairs close to the shorting plate in order to allow most of the power to be radiated [36].

In this dissertation, a prototype of this cross slotted arrays is proposed, with an improvement to the efficiency and the design procedure.


Figure II.3: Cross slotted rectangular waveguide array used in satellite communication transmitter. [39]

## - WAVEGUIDE Polarizers

A waveguide wave polarizer is a microwave device that transforms a linearly polarized wave at its feeding port into a circularly or elliptically polarized wave at its output port. They are usually implemented in square or circular waveguides. Multiple techniques are used to do this transformation that are explained below:

## $\square$ Corrugated waveguide polarizers and Iris Polarizers:

This type of polarizers uses a square waveguide as shown in Figure II.4. In its simplest form, the $T E_{10}$ and the $T E_{01}$ modes are excited in the waveguide. The corrugations are on either the E-plane or the H-plane of the square waveguide. These corrugations will delay
either the $T E_{10}$ or the $T E_{01}$ modes propagation, creating a phase shift. If the corrugations are designed correctly, the phase shift can be $90^{\circ}$. [40] - [43].

(a)

(b)

Figure II.4: (a) Corrugations on the wall of a square waveguide, (b) cross section view of the waveguide. [40]

A solution for not exciting two modes in the waveguide is by using a square waveguide propagating $T E_{10}$ and then adding a section of circular waveguide rotary joint that is connected to another square waveguide where the corrugations are located. The rotary joint creates a rotation of $45^{\circ}$ with the initial waveguide, hence partitioning its fields into $T E_{10}$ and $T E_{01}$ modes. [43]

This type of polarizers is optimized to produce either two narrow bands of operation, or to produce a bandwidth of operation between $5-10 \%$ of the frequency band. The polarizer can create a very flat phase shift in its operating bandwidth with a phase
difference of $90^{\circ}-91^{\circ}$. It can either create a left hand circularly polarized (LHCP) wave or a right hand circularly polarized (RHCP) wave.

In circular waveguides, irises are placed at a $45^{\circ}$ angle from the direction of the $T E_{11}$ propagating mode. At the $45^{\circ}$ angle, the irises see the mode as two perpendicular components, one of them is parallel to the iris, and the other is perpendicular to the iris. The irises then delay the parallel component by $90^{\circ}$ resulting in a circularly polarized wave [44]. If the $T E_{11}$ is vertical, then the resulting polarization is RHCP and if it is horizontal, then the resulting wave is LHCP.


Figure II.5: Schematic of the irises inside a circular waveguide and the resulting circularly polarized wave. [44]

This type of polarizers has a bandwidth of $10 \%$ of the frequency band and it can only produce one polarization once it is installed in a system.

This particular type of polarizers shown in this section are realized at frequencies up to Ka-band. They are very expensive to achieve at W/V bands because the process of fabrication is very complicated and requires very high precision.

## Dielectric Septum Loading Polarizer:

This type of polarizer is shown in Figure II.6. As in the case of the iris polarizers, the septum is at a $45^{\circ}$ angle with the propagation of the $T E_{11}$ mode propagating in the circular waveguide. The dielectric septum delays one of the components of the electric field. The delay of $90^{\circ}$ results in a circularly polarized wave. Matching this type of polarizers at its input causes problems. Slots in the perpendicular plane to the septum are added to reduce the reflection at the input of the polarizer [45].


Figure II.6: Schematic of the polarizer using dielectric septum, (a) 3D, (b) cross section view.

This design only results in one polarization sense at the same configuration. It is realized at 60 GHz with a good performance between $60-62.5 \mathrm{GHz}$. Other polarizers that use either dielectric septum or metallic septum are available commercially. These polarizers can produce a bandwidth of operation that could reach $20 \%$ of the frequency band. [46]

## Septum OMT polarizer:

This type of polarizers is a three port network, where the third port has two modes propagating in it and forming a circularly polarized wave. If the first port is fed, a RHCP wave is generated at port 3 , and if port 2 is fed, a LHCP wave is formed at port 3 as shown in Figure II. 7 [47]-[54].


Figure II.7: Schematic of the septum OMT polarizer showing its operation modes. [54]

This type of polarizers has an isolation of more than 20 dB between ports 1 and 2 and can produce axial ratio values of less than 1 dB over a frequency band of $15-18 \%$ of the
frequency band for low frequencies [47]. These characteristics, make the polarizer a good device for use in both transmission and reception simultaneously.

The basic model of the septum is a stepped septum that creates high isolation between ports 1 and 2, while creating two orthogonal modes in the square waveguide output with $90^{\circ}$ phase difference between them [47],[48],[49]. In some designs a corrugation is added in the square waveguide output in order to correct the phase and make it more flat and close to $90^{\circ}$ [50].

An optimization of the model is introduced in [51], [52]. This optimization makes the stepping also in the thickness of the septum. The designs done are in the Ka-band producing a good operation bandwidth of $12 \%$ of the frequency band. In [51], the design is made reconfigurable by introducing shorting plates controlled by piezo-motors in order to control which channel operates at a certain point for transmission as shown in Figure II.8.


Figure II.8: The different modes of operation of the reconfigurable OMT polarizer. [51]

This polarizer can also be implemented in circular waveguides without the need of use of a square to circular waveguide transition [53]. In this particular design, the polarizer is used to feed a conical horn feeding a reflector antenna.

At high frequencies, due to fabrication inaccuracies, this design can produce a $10 \%$ bandwidth of operation [54]. If used in the WTLE experiment, two of these polarizers are needed in order to operate at 72 GHz and 84 GHz . The Multi-band polarizer, that is introduced in this dissertation, makes the design simpler, less expensive and operating at 72 and 84 GHz by using the same device.

# CROSS SLOTTED RECTANGULAR WAVEGUIDE ARRAY 

## - INTRODUCTION

The initial work done at W/V band frequencies is the design of reconfigurable arrays using liquid crystal in its nematic form. This type of arrays is hard to fabricate, and has low efficiency at 76 GHz . This work is discussed in Appendix A in more details. For antennas that are more efficient the work presented hereafter is done.

The implementation of a planar cross-slotted rectangular waveguide array that is operating with a dual feed is presented here [55], [56]. This type of arrays was used previously in transmitter systems for satellite communication at lower frequencies (X-band and Ka-band) [35], [39]. Here the dual feed results in the generation of either RHCP or LHCP radiations at $84-86 \mathrm{GHz}$. The integration of $z$-shaped arm slots on the broad-wall of the waveguide is also implemented without affecting the total circular polarization performance and to improve design flexibility and overall radiation efficiency. The new cross-slot concept, proposed in this work, can be designed for any frequency in the operating $\mathrm{TE}_{10}$ fundamental mode of the waveguide. In what follows, a detailed description of the new array design, its fabrication process, and all analytical and experimental results are presented in details and discussed.

## - Design of Z-Shaped Cross-Slot

In this section, the design of a $z$-shaped cross-slot along the broad-wall of a WR-10 rectangular waveguide operating in the $\mathrm{TE}_{10}$ mode is investigated. For the fundamental mode and along the broad-wall of the waveguide, the magnetic field has two components: $H_{\mathrm{x}}$ and $H_{\mathrm{z}}$. As seen from Eq. 1 and Eq. 2, the two components of the magnetic field are $90^{\circ}$ out of phase at all times.

$$
\begin{align*}
& H_{x}=-\frac{A}{Z_{10}} \sin \frac{\pi x}{a} e^{-j \beta z}  \tag{III.1}\\
& H_{z}=\frac{j \pi A}{\beta a Z_{10}} \cos \frac{\pi x}{a} e^{-j \beta z} \tag{III.2}
\end{align*}
$$

At a certain range of positions $s$, far from the narrow-wall of the waveguide, the magnitudes of the magnetic field components are equal for certain frequencies [34]. These positions are given in Eq. 3.

$$
\begin{equation*}
s= \pm \frac{a}{\pi} \tan ^{-1} \frac{1}{\sqrt{(2 f a / c)^{2}-1}} \tag{III.3}
\end{equation*}
$$

Circularly polarized waves can therefore be radiated by having a slot centered at a position $s$. For each frequency there are two positions $s$ where circular polarization is achieved. Also, by changing the $z$-direction in which the wave is propagating, the radiated wave can be changed from RHCP to LHCP. This is physically done by changing the feeding in the waveguide from one port to the other port.

In order to radiate the maximum amount of power, a cross- slot is usually adopted. Figure III.1(a) depicts a conventional cross-slot placed at a distance $s$ away from one of the
waveguide narrow-walls. The cross-slot has two arms with length $\mathrm{L}_{\mathrm{a}} \approx \lambda / 2$ and width $\mathrm{W}_{\mathrm{a}}<\lambda / 10$.

At certain frequencies, the slot arm projection on the $x$-axis $\left(A_{\mathrm{p}}\right)$ is larger than the position $s$. This causes the slot to overlap with the narrow-wall of the rectangular waveguide. Figure III.1(b) shows the difference $X$ between $A_{\mathrm{p}}$ and $s$ for a span of frequencies from 75 GHz up to 110 GHz for a WR-10 waveguide. From this plot, one concludes that for a WR-10 waveguide, the design of a cross-slot is feasible only for frequencies less than or equal to 84 GHz . Thus, for our frequency range of interest (84-86 GHz ), the conventional cross-slot cannot be adopted. To solve this problem, the position $s$ can be increased in order to accommodate the arms within the broad-wall. By increasing $s$, the isolation between the co- and cross-polarization in the direction of the highest gain is reduced. This is due to the fact that the magnitudes of the components of the magnetic field are not equal at these positions.

Another solution is to use $z$-shaped arms for the cross-slot by bending the tips of the arms as shown in Figure III.1(c). The new arms will fit within the borders of the broadwall for frequencies beyond 84 GHz while maintaining the same positions $s$. By optimizing the dimensions of the $z$-shaped arms, the slot is able to radiate at the same frequencies than the cross-slot while having a shorter arm projection on the $x$-axis.

To investigate the effect of changing the shape of the slot, a WR-10 rectangular waveguide ( $a=2.54 \mathrm{~mm} \times b=1.27 \mathrm{~mm}$ ) is first designed with the cross-slot shown in Figure III.1(a). The cross-slot has $L_{\mathrm{a}}=1.83 \mathrm{~mm}, W_{\mathrm{a}}=0.15 \mathrm{~mm}$ and is at a distance $s=$ 0.67 mm from the narrow-wall. At 84 GHz , the slot radiates $73.5 \%$ of the power. Figure
III. 2 shows that the achieved gain is 4 dB with an isolation of 28 dB between LHCP and RHCP at the direction of maximum gain.

The next step is to replace the cross-slot by the one shown in Figure III.1(c). The new corresponding dimensions are $L_{z 1}=1.61 \mathrm{~mm}$ with the short arm length $L_{z 2}=0.38 \mathrm{~mm}$. The slot width is $W_{\mathrm{a}}=0.15 \mathrm{~mm}$ and centered at $s=0.82 \mathrm{~mm}$ from the narrow-wall of the WR10 waveguide. At $84 \mathrm{GHz}, 70.5 \%$ of the power is radiated by the slot. The achieved gain is 4 dB with an isolation of 28 dB between the co- and cross-polarized radiations at the direction of maximum gain, as shown in Figure III.2.


Figure III.1: (a) Conventional cross-slot dimensions, (b) The difference between the arm projection Ap and the slot position s over frequencies, (c) The $z$-shaped arm cross-slot structure.


Figure III.2: The co- and cross-polarization radiation patterns for the conventional and z-arm slots in the XZ plane.


Figure III.3: (a) A single WR-10 waveguide element with 4 and 16 slots, (b) The gain variation as a function of the number of slots.

The comparison between the two designs at 84 GHz shows that the $z$-shaped slot preserves the same polarization isolation as the conventional slot. However, the $z$-shaped arm cross-slot ensures that the full slot topology fits on the broad-wall without any overlap with the narrow-wall for all the frequencies of the W-band.

## - Design of A Waveguide Array

The waveguide topology discussed in Section II exhibits a gain of 4 dB . For the gain to be improved, an array of slots is introduced along the WR-10 waveguide broad-walls.

## Single Waveguide with Multiple Slots:

The number of $z$-shaped arm cross-slots determine the maximum gain that can be produced from a single WR-10 waveguide element. Thus, it is essential to determine the optimal number of slots that can produce the maximum achievable gain levels. Figure III. 3 (a) shows the corresponding structure with 4 and 16 slots. For the different scenarios, the $z$-shaped arm cross-slots have the same "long arm" length $L_{z 1}=1.67 \mathrm{~mm}$ and the same "short arm" length $L_{\mathrm{z} 2}=0.433 \mathrm{~mm}$. The slot width is chosen to be $W_{\mathrm{a}}=0.15 \mathrm{~mm}$ and centered at a distance $s=0.715 \mathrm{~mm}$ from the wall of the WR-10 waveguide. The mutual coupling resulting from introducing several slots in close proximity to each other necessitates a change in the dimensions of the slots. At first, the waveguide was designed with 4 slots. The gain exhibited is 8.5 dB at 85 GHz . Doubling the number of slots results in a gain increase by 2 to 3 dB for operation across the frequency range $84-86 \mathrm{GHz}$. This can be inspected in Figure III.3(b). The gain reaches 13.5 to 14 dB with $N=16$ slots. If the number of slots increases to 32 , the increase in gain is only about $0.5-1 \mathrm{~dB}$. Thus, it is found
that $N=16$ slots is a good tradeoff between array size and maximum achievable gain. For this number of slots, the length of the waveguide is 39.3 mm .

(b)

Figure III.4: The co- and cross-polarized radiation patterns at 85.5 GHz with $\mathrm{N}=16$ slots in the (a) elevation (XZ) plane, (b) azimuthal $\left(\theta=40^{\circ}\right)$ plane.

The feeding phase difference between the various slots can cause high side lobe levels if not designed appropriately. Thus, by optimizing the distance of separation between the slots, the side lobe level can be minimized. A separation distance of $\mathrm{D}_{\mathrm{x}}=2.05 \mathrm{~mm} \approx \lambda / 2$
was determined to be the optimal choice through an iterative analysis. The phase difference causes the main beam to switch from $\theta=0^{\circ}$ (one slot) to $\theta=40^{\circ}$. Figure III.4(a) shows the antenna gain pattern for $N=16$ slots at $f=85.5 \mathrm{GHz}$. It is noticed that the half power beamwidth (HPBW) in the elevation plane ( $X Z$ plane) is $8.7^{\circ}$ with an isolation of more than 18 dB between the co- and cross-polarized gain patterns. A drop in the HPBW is achieved as compared to a single slot due to the increase in the gain levels. The HPBW and isolation values vary slightly over the entire bandwidth.


Figure III.5: (a) The proposed $8 \times 16$ cross-slotted rectangular waveguide array, (b) The gain variation with the number of waveguides.

The increase in the number of slots in one direction ( $x$-direction) while preserving the same number in the other direction ( $y$-direction) makes the main beam narrower in one plane and wider in the other. For example, the HPBW in the azimuthal plane $\left(\theta=40^{\circ}\right)$
increases to $55^{\circ}$ as summarized in Figure III.4(b). One possible solution to maintain almost the same HPBW along both planes is to transform the proposed structure into a two dimensional array as detailed next.

## 2-D Waveguide Array with 16 Slots:

A 2-D array is designed by increasing the number of WR-10 waveguides having 16 slots. The new structure with 8 waveguide elements is shown in Figure III.5(a). The same slots used in a single waveguide array, are used for the 2-D array. The waveguides of the 2-D array are placed next to each other. All of the waveguides have the cross-slots on the same side of the wall. The separation between the slots of the same waveguide is kept the same as in the case of a single waveguide $D_{x}=2.05 \mathrm{~mm}$. The separation between the slots of the same row in adjacent waveguides is taken to be $D_{y}=3.04 \mathrm{~mm}$. This separation is limited by the WR-10 waveguide width $(a=2.54 \mathrm{~mm})$ as well as the thickness $\left(t_{w 1}=\right.$ 0.5 mm ) of the walls between any two adjacent waveguides.

The change in the array gain with the increase in the number of the waveguide elements is plotted in Figure III. 5 (b). The gain of the array improves by 2.5 to 3.5 dB as the number of waveguides is doubled. An increase from around 14 dB to around 24 dB is obtained for the case of eight waveguides. These results are obtained by forcing all the waveguides in the simulation environment to be fed in phase and with the same amount of power.

Feeding the various waveguides in phase across the entire frequency band of operation $(84-86 \mathrm{GHz})$ is a challenging task. The phase effect on the performance of the array is analyzed with 8 waveguides. If the consecutive waveguides are fed with a constant phase
difference $\Delta \phi_{\text {in }} \neq 0^{\circ}$ and with the same amount of power, the main beam of the array becomes directed in a plane $\phi \neq 0^{\circ}$. The following four scenarios are investigated:

Case 1:

If the phase difference is $0^{\circ}<\Delta \phi_{\text {in }}<20^{\circ}$ between the waveguide elements, the main beam is entirely circularly polarized and exhibits a gain of 24 dB , with an isolation between the co- and cross-polarized components of more than 22 dB . The main beam direction is in a plane $0^{\circ}<\phi<5^{\circ}$ as shown in Figure III.6.

Case 2:

If the phase difference is $20^{\circ}<\Delta \phi_{\text {in }}<45^{\circ}$, the array produces a main beam with a gain of 23.5 dB and the isolation between the co- and cross-polarized components is between 15 and 20 dB . This phase causes circular polarization to be lost across some parts of the HPBW as shown in Figure III.6.

## Case 3:

If the phase difference is $\Delta \phi_{\text {in }} \geq 45^{\circ}$, the array does not produce a circularly polarized radiation at the main beam, and the direction of the main beam is in a plane $\phi>14^{\circ}$ as shown in Figure III.6. In the case of a single slot, the main beam is circularly polarized along the $z$-axis. The further away from the $z$-axis, the less the radiated E-fields are circularly polarized. The phase difference discussed in this case causes the electric fields far from the $z$-axis to add up, and it causes the circularly polarized fields close to the $z$-axis to cancel each other. This effect causes the main beam to be elliptically polarized and directed in a plane $\phi>14^{\circ}$.

## Case 4:

In order to direct the main beam in the plane $\phi=0^{\circ}$, the phase differences between the consecutive waveguides should be $0^{\circ}<\Delta \phi_{\text {in }}<20^{\circ}$. The various waveguides must also be fed with the same amount of power. In addition, if the pairs of waveguides $(1,8),(2,7)$, $(3,6)$ and $(4,5)$, shown in Figure III.5(a), have the same phase at their input then the main beam is directed in the desired plane.


Figure III.6: The radiation pattern of the array in the azimuthal $\left(\theta=40^{\circ}\right)$ plane for scenarios 1,2 and 3.

Figure III. 7 compares the gain of the array when the various elements are fed in phase with the case when the phase difference between the elements is $20^{\circ}$. It is found that the same gain response is obtained. The isolation between the co- and cross-polarized gain patterns is preserved and the gain is higher than 22 dB in the direction of the highest gain and over the frequency band of operation. More specifically, the array has the main beam directed in the plane $\phi=0^{\circ}$ and at an angle $\theta=40^{\circ}$. The array's HPBW is $\Theta_{e p}=8^{\circ}$ in the elevation plane $\left(\phi=0^{\circ}\right)$, and $\Theta_{\mathrm{ap}}=12.5^{\circ}$ in the azimuthal plane $\left(\theta=40^{\circ}\right)$. The maximum gain achieved is 24 dB over the frequency span of $84-86 \mathrm{GHz}$.


Figure III.7: The gain patterns of the array in the azimuthal $\left(\theta=40^{\circ}\right)$ plane for the case where all waveguides are fed in phase and for case 4.

## Feeding Network Design

In the previous section, it is concluded that the optimum way to feed the different waveguide elements is by either having them all fed in phase or by allowing the symmetrical pairs of waveguides to be fed in phase, and the consecutive waveguides to have a small phase difference $\left(0^{\circ}<\Delta \phi_{\text {in }}<20^{\circ}\right)$. To achieve this, a separate feeding network must be designed.

The first proposed feeding network under investigation is depicted in Figure III.8(a). It is composed of one feeding waveguide that connects directly to the eight WR-10 waveguide elements. The connection is achieved by introducing rectangular slots on the narrow-wall of the feeding waveguide and shorting its other end. These slots are centered at the centerline of the narrow-wall and radiate the same power into the eight waveguide
elements of the array. The distance of separation between the slots is very crucial on the array performance since such distance sets the phase difference between the various eight waveguide elements.

The eight waveguide elements have a center to center separation $D_{y}=3.04 \mathrm{~mm}$. If a WR-10 rectangular waveguide is used directly as the feeding waveguide, the separation between the rectangular slots $(3.04 \mathrm{~mm})$ results in a $\Delta \phi_{\text {in }}=223^{\circ}$ phase difference between the consecutive waveguide elements at 85 GHz . This phase difference causes the main beam to be in the plane $\phi=32^{\circ}$. The gain of the array is reduced to 22 dB and the circular polarization is lost as shown in Figure III.9. The same behavior is obtained if the feeding waveguide is a WR-12 rectangular waveguide and the distance of separation between the rectangular slots is kept to be 3.04 mm . In this case, the phase difference is $\Delta \phi_{\text {in }}=254.5^{\circ}$ at 85 GHz . The main beam is directed in the plane $\phi=37^{\circ}$ and the array is no longer circularly polarized. The gain patterns for this case are also incorporated in Figure III.9. Based on these results, it is concluded that a transition must exist between the eight rectangular slots and the opening of the eight WR-10 waveguides.


Figure III.8: The proposed feeding with (a) direct connection between the feeding waveguide and the eight elements WR-10 waveguides, (b) The proposed transition to achieve the required phase difference between the various elements.


Figure III.9: The gain pattern of the array fed directly by a WR-10 or WR-12 rectangular waveguide in the azimuthal $\left(\theta=40^{\circ}\right)$ plane.

The modified layout of the proposed feeding network is depicted in Figure III.8(b). For this case, a WR-12 feeding waveguide with wall thickness $t_{w 2}=0.5 \mathrm{~mm}$ is adopted since during measurements, the feeding source has a WR-12 rectangular waveguide output. In order to achieve a phase difference of $\Delta \phi_{\mathrm{in}}=0^{\circ}+360^{\circ}$ between the consecutive waveguide elements at 85 GHz , the separation between the rectangular slots of the feeding waveguide should be increased from $D_{y}=3.04 \mathrm{~mm}$ to $D_{y f}=4.22 \mathrm{~mm}$. The feeding waveguide is shorted at its other end with a separation of $S_{\text {short }}=1.71 \mathrm{mmbetween}$ the last slot center and the shorting plate. The rectangular slots have a length $L_{s f}=1.91 \mathrm{~mm}$ and a width $\mathrm{W}_{\text {sf }}=0.25 \mathrm{~mm}$. These slot dimensions ensure that the feeding network has a good input match between $84-86 \mathrm{GHz}$ as shown in Figure III. 10 (a). The insertion loss $\left|\mathrm{S}_{\mathrm{i}}\right|$ [dB] between the input of the feeding waveguide and each of the rectangular slots is also shown in Figure III. 10 (a). These results show that each waveguide element of the array is able to receive approximately an equal amount of power from the rectangular slots at the
frequency bands of operation. The phase difference between the input of the feeding network and the eight rectangular slots is plotted in Figure III. 10 (b). A phase difference in the range $0^{\circ}<\Delta \phi_{\text {in }}<20^{\circ}$ is maintained.

Since the separation distances between the feeding rectangular slots $D_{y f}$ and the radiating cross-slots $D_{y}$ are different, serpentine WR-10 rectangular waveguides are used as shown in Figure III.8(b). The serpentine waveguides are symmetrical for the pairs of waveguides $(1,8),(2,7),(3,6)$ and $(4,5)$. This design ensures the same phase at the output of each pair of waveguides and a small phase difference between the consecutive waveguides. These waveguides are connected to the feeding slots by WR-10 waveguide sections of length $\mathrm{L}_{\mathrm{ext}}=5 \mathrm{~mm}$ in order to minimize the wave reflection caused by the curvature of the serpentine shape. The serpentine waveguides have a curvature radius $\mathrm{R}_{\text {serp }}=50 \mathrm{~mm}$ and a total length $\mathrm{L}_{\text {serp }}=58.15 \mathrm{~mm}$.


Figure III.10: The feeding network S-parameters magnitude in dB and phase in degrees.

## - Full Array Design

Two feeding networks, as the one detailed in the previous section, are used to feed the two sides of the eight elements WR-10 waveguide array. The various parts of the full array design are detailed in Figure III.11(a). Two WR-12 waveguides (parts 1 and 1a) constitute the input and the output ports of the whole array structure. A metallic sheet (parts 2 and 2a) with eight rectangular slots connect each of the WR-12 waveguide to the input of the serpentine shaped connections of the feeding network. Part 3 is composed of the eight rectangular WR-10 waveguides along with the serpentine shaped connections. A metallic sheet with the $8 \times 16 z$-shaped arm cross-slots constitutes part 4 . The last part (part 5) is used as a metallic enclosure around the $z$-shaped arm cross-slots in order to ensure a bound connection between the various parts of the array and to hold the array together.

The metallic sheet (part 4) has a thickness $\mathrm{t}_{\mathrm{w} 3}=0.4 \mathrm{~mm}$. It is noticed that for $\mathrm{t}_{\mathrm{w} 3}=$ 0.5 mm , the best isolation between the co- and cross-polarized gain patterns is attained. However, $\mathrm{t}_{\mathrm{w} 3}=0.4 \mathrm{~mm}$ is the limit imposed by the fabrication process that is implemented for this particular array. The antenna assembly technique is also highlighted in Figure III.11(a) and the physical layout of the assembled array is shown in Figure III.11(b).


Figure III.11: The full array design (a) disassembled, (b) assembled and (c) radiation pattern for each port.

Table III.1: Dimensions of the different parameters of the array.
z-shaped arm cross-slot parameters

| Parameter | Description | Value |
| :---: | :---: | :---: |
| $L_{z 1}$ | long arm length | 1.65 mm |
| $L_{z 2}$ | short arm length | 0.433 mm |
| $W_{a}$ | arm width | 0.15 mm |
| $S$ | slot center position | 0.7 mm |
| $\alpha_{2}$ | arm angle with y-axis | $43^{\circ}$ |
| $D_{x}$ | slots separation along x-axis | 2.05 mm |
| $D_{y}$ | slots separation along y-axis | 3.05 mm |
| $t_{w 3}$ | slot thickness | 0.4 mm |

Feeding Network Parameters

| Parameter | Description | Value |
| :---: | :---: | :---: |
| $L_{s f}$ | feeding slot length | 1.914 mm |
| $W_{s f}$ | feeding slot width | 0.25 mm |
| $D_{y f}$ | feeding slots separation | 4.22 mm |
| $S_{\text {short }}$ | serpentine waveguide curvature radius | 1.71 mm |
| $R_{\text {serp }}$ | serpentine waveguides length | 50 mm |
| $L_{\text {serp }}$ | feeding slot thickness | 58.15 mm |
| $t_{w 2}$ |  | 0.5 mm |

The final design dimensions are optimized using the CST built-in optimization tool in order to deliver a gain of $23-24 \mathrm{~dB}$ over the frequency band of operation ( $84-86 \mathrm{GHz}$ ). The main beam in the elevation plane is directed in $\theta=40^{\circ}$ when fed through port 1 and $\theta=-40^{\circ}$ when fed through port 2 as shown in Figure III. 11 (c). The final design parameters are shown in Table III.1. It is essential to note that the feeding network and the metallic enclosure have no effect on the radiation pattern of the array.

## Measured Results

The entire array structure is fabricated as shown in Figure III.12(a). The antenna operating bandwidth is $84.1-85.8 \mathrm{GHz}$ for port 1 , and $84.35-86 \mathrm{GHz}$ for port 2 as depicted in Figure III.13. The measurement results show an acceptable agreement with the simulated data. In fact, port 1 exhibits a slightly different yet tolerable performance. The slight discrepancy is due to some of the fabrication imprecisions ( $5-10 \mu \mathrm{~m}$ ) in terms of milling the various slots in addition to the assembly misalignment. Both ports share a 1.45 GHz operating bandwidth extending between $84.35-85.8 \mathrm{GHz}$. The antenna has a great isolation between both ports since the measured isolation between the ports is more than 20 dB over the frequency of operation as shown in Figure III.13.

The antenna's radiation pattern is measured in the azimuthal $\left(\theta=40^{\circ}\right)$ and elevation ( $\phi=0^{\circ}$ ) planes over different frequencies in the operating bandwidth. The measurement setup for both planes is shown in Figure III.12(b) and (c). A waveguide rotary joint is used for the azimuthal plane $\left(\theta=40^{\circ}\right)$ measurement. When port 1 is fed, port 2 is loaded with a WR-12 waveguide matched load and vice versa. The comparison between the simulated and measured gain patterns at $f=85 \mathrm{GHz}$ and when feeding either port 1 or 2 are summarized in Figure III. 14 for the elevation plane $\left(\phi=0^{\circ}\right)$ and in Figure III. 15 for the azimuthal plane $\left(\theta=40^{\circ}\right)$. The measured gain patterns at 85 GHz prove that when the antenna is fed through port 1 , it creates a RHCP main beam in the elevation plane in the direction $\theta=40^{\circ}$. When it is fed through port 2 , the antenna creates a LHCP main beam in the elevation plane in the direction $\theta=-40^{\circ}$. These results show that the change in port feeding causes the polarization to switch from LHCP to RHCP.

(a)


$$
\emptyset=0^{\circ} \text { Plane Setup }
$$

(b)

$$
\theta=40^{\circ} \text { Plane Setup }
$$

(c)

Figure III.12: (a) The fabricated antenna prototype, (b) The radiation pattern measurement setup in the (b) elevation plane $\left(\phi=0^{\circ}\right)$ and (c) azimuthal $\left(\theta=40^{\circ}\right)$.

The antenna's calculated RHCP gain is between 23.26 dB and 24.09 dB and the calculated
LHCP gain is between 23.2 dB and 24.35 dB over the operating bandwidth as shown in

Table III.2. The measured results at 85 GHz show that in the elevation plane, the antenna's measured HPBW is $\Theta_{\mathrm{ep}}=8^{\circ}$ and it is entirely circularly polarized for both polarizations. The side lobe level is around 12 dB and the maximum grating lobe level is around 13.5 dB for both polarizations as shown in Figure III.14. In the azimuthal plane, the antenna's HPBW is $\Theta_{\mathrm{ap}}=12^{\circ}$ and it is entirely circularly polarized for both polarizations. The antenna's side lobe level in this plane cut is around 11.5 dB for RHCP and 10.5 dB for LHCP as shown in Figure III. 15.


Figure III.13: (top) The simulated and measured results for the reflection coefficient at both ports, (bottom) The simulated and measured isolation between the two ports.

Table III.2: The array calculated gain for both ports at different frequencies.

| Frequency | RHCP Gain |  | LHCP Gain |  |
| :---: | :---: | :---: | :---: | :---: |
|  | Simulated | Measured | Simulated | Measured |
| 84.35 GHz | 21.5 dB | 23.72 dB | 21.5 dB | 23.2 dB |
| 84.5 GHz | 23.1 dB | 24.09 dB | 23.1 dB | 23.7 dB |
| 85 GHz | 23.31 dB | 23.76 dB | 23.31 dB | 24.33 dB |
| 85.3 GHz | 24.07 dB | 23.83 dB | 24.07 dB | 24.35 dB |
| 85.5 GHz | 24.4 dB | 23.79 dB | 24.4 dB | 23.92 dB |
| 85.7 GHz | 24.5 dB | 23.26 dB | 24.5 dB | 23.86 dB |

Table III.3: Comparison between different planar slotted waveguides working at millimeter wave frequencies.

| Reference | Frequency | Gain | Polarization | Efficiency | Size |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $(\mathrm{GHz})$ | $(\mathrm{dB})$ |  |  |  |  |
| SIW 2D Array[57] | 81.75 | 14 | Linear | $50 \%$ | $2 \times 8$ |  |
| 2D Slot array [58] | 94 | 26.8 | Linear | $81.9 \%$ | $8 \times 8$ |  |
| Ka-band X-slot array [39] | 30 | 22.9 | RHCP | $74 \%$ | $16 \times 1$ |  |
| Dual linear Polarization | $30 / 35$ | $24.8 / 25.4$ | VP/HP | N/A | $8 \times 10 / 9 \times 10$ |  |
| [59] |  |  |  |  | $82 \%$ | $8 \times 16$ |
| This work | 85 | 24 | LH/RH CP |  |  |  |

The isolation between RHCP and LHCP radiations in the direction of the maximum gain for both ports is measured. In the frequency of operation of both ports, the isolation level is found to be more than 18 dB and reaching 26 dB in certain cases as shown in Figure
III.16. For the frequency band $84.4-85.55 \mathrm{GHz}$, both ports produce an isolation of more than 20 dB .


Figure III.14: Simulated and measured gain patterns of the array in the elevation plane ( $\phi=0^{\circ}$ )
at 85 GHz when (a) port 1 is fed, (b) port 2 is fed.


Figure III.15: Simulated and measured gain patterns of the array in the azimuth plane $\left(\theta=40^{\circ}\right)$

$$
\text { at } 85 \mathrm{GHz} \text { when (a) port } 1 \text { is fed, (b) port } 2 \text { is fed. }
$$

Table III. 3 compares the proposed design in this work with the ones available in the literature. In general, the slotted rectangular waveguide arrays designed at millimeter-wave can all generate a gain close to 25 dB with an efficiency higher than 70\% [39], [58], [59]. These characteristics are close to the results presented in this work. However, the arrays operating at W-band are linearly polarized [57], [58] while the proposed design is circularly
polarized with the ability to generate both LHCP and RHCP radiations. The array presented in [39] is designed to operate at 30 GHz and can generate RHCP radiation but does not have the ability to change the polarization sense to LHCP. Another array that can generate a change in polarization sense is designed at 30 GHz and 35 GHz [59]. It incorporates two different sets of slots in different waveguides, it also operates at different frequencies and it is able to generate different schemes of linear polarization [59]. However, in this work, the same waveguides can create this change in polarization when fed from different ports. Another type of slotted waveguide arrays at W-band is the slotted SIW with an efficiency of $50 \%$ [57].


Figure III.16: The isolation between the co- and cross-polarization in the direction of the maximum gain.

## - Conclusion

A new high gain circularly polarized waveguide antenna array is presented. The waveguide antenna array has eight elements with $z$-shaped cross-slots across each element's broad-wall. The rectangular waveguide array operates between $84.35-85.8$ GHz . The array can produce a RHCP radiation with a main beam in the direction $\theta=40^{\circ}$
if fed through port 1 and a LHCP radiation with the main beam in the direction $\theta=-40^{\circ}$ if fed through port 2 . The gain achieved for both ports is between 23 and 24.5 dB . The HPBW has an elliptical shape and is entirely circularly polarized in both elevation and azimuthal planes. The isolation between the co- and cross-polarized gain patterns in the direction of the maximum gain is more than 20 dB for both ports. The array is proven to be a good candidate for point to point terrestrial communication systems as well as satellite to earth communication systems.

## CROSS SLOT POLARIZER

## I. Introduction

The array of cross slotted waveguides presented in the previous chapter can generate a RHCP beam in the direction $\theta=+40^{\circ}$ when fed through port 1 , and LHCP beam in the direction $\theta=-40^{\circ}$ when fed through port 2 . This requires the array to be rotated in order for the two beams to be pointing at the same target. This means that the LHCP and RHCP properties of the array cannot be used for the same target at the same time. A solution for this problem is to extract the power from one slot into a waveguide and radiate it into $\theta=$ $0^{\circ}$ for both LHCP and RHCP radiations. This device can be used as a polarizer for horn antennas when it is used in the transmitting mode, or it could be used as an orthomode transducer in the receiving mode. In sections II and III, the theory of a slot radiating into a circular waveguide is presented and its performance while feeding a conical horn is shown. Sections IV and V detail the theory of a cross slot radiating into a square waveguide and feeding a pyramidal horn. Section VI details the design of a serpentine combiner aiming to improve the performance of a single slot while section VII details the design of a square waveguide power combiner that improves the performance of the square waveguide polarizer design. Finally, this chapter concludes with section VIII where the design of multi band cross slot polarizer systems is illustrated.

## II. Slot Radiating Into a Circular Waveguide

In this section, a slot radiating into a circular waveguide is discussed. The position of the slot on the broad-wall that excites a circularly polarized wave in a circular waveguide is derived.

## Theory

Let's assume that the slot is located on the broad-wall of the rectangular waveguide with its center at $\mathrm{x}=\mathrm{s}, \mathrm{y}=\mathrm{b}$ and $\mathrm{z}=0$ as shown in Figure IV. 1 (a).

In a rectangular waveguide, the field equations for $\mathrm{TE}_{10}$ mode are:

$$
\begin{align*}
& E_{y}=A \sin \frac{\pi x}{a} e^{-j \beta z}  \tag{IV.1}\\
& H_{x}=-\frac{A}{Z_{10}} \sin \frac{\pi x}{a} e^{-j \beta z}  \tag{IV.2}\\
& H_{z}=\frac{j \pi A}{\beta a z_{10}} \cos \frac{\pi x}{a} e^{-j \beta z} \tag{IV.3}
\end{align*}
$$

where $Z_{10}=\frac{k_{0} \eta_{0}}{\beta}$

These fields are shown in Figure IV. 1 (b). At the slot, the electric and magnetic polarizabilities are given by the equations below:

$$
\begin{align*}
& \overrightarrow{P_{e}}=A \varepsilon_{0} \alpha_{e} \sin \frac{\pi s}{a} \delta(x-s) \delta(y-b) \delta(z) \hat{y}  \tag{IV.4}\\
& \overrightarrow{P_{m}}=-A \alpha_{m}\left[-\frac{\hat{x}}{z_{10}} \sin \frac{\pi s}{a}+\hat{z} \frac{j \pi}{\beta a z_{10}} \cos \frac{\pi s}{a}\right] \delta(x-s) \delta(y-b) \delta(z) \tag{IV.5}
\end{align*}
$$

where $\alpha_{e}$ and $\alpha_{m}$ are constants related to the slot shape and dimensions.


Figure IV.1: (a) The slot on the broad-wall of a rectangular waveguide, (b) $\mathrm{TE}_{10}$ mode fields on the broad-wall of a rectangular waveguide, (c) $\mathrm{TE}_{11}$ mode fields in a circular waveguide.

The induced electric and magnetic currents at the slot are:

$$
\begin{align*}
& \vec{J}=j \omega \overrightarrow{P_{e}}  \tag{IV.6}\\
& \vec{M}=j \omega \mu_{0} \overrightarrow{P_{m}} \tag{IV.7}
\end{align*}
$$

Since the slot is radiating the power into a circular waveguide, it is essential to examine the fields excited in a circular waveguide. The mode that can create a circularly polarized wave is the fundamental mode $\mathrm{TE}_{11}$ as shown in Figure IV. 1 (c). This mode has the potential to propagate two orthogonal components with a $90^{\circ}$ phase difference. The two following modes that can potentially propagate in a circular waveguide $\left(\mathrm{TM}_{01}-\mathrm{TE}_{21}\right)$, have their electric and magnetic fields symmetrical along the center of the waveguide, which makes them bad candidates for creating a circularly polarized wave propagating in the waveguide.

The fundamental mode $\mathrm{TE}_{11}$ has the following field equations:

$$
\begin{align*}
& E_{z}=0  \tag{IV.8}\\
& E_{\rho}=-\frac{j \omega \mu}{k_{c}^{2} \rho}[B \cos \phi-C \sin \phi] J_{1}\left(k_{c} \rho\right) e^{-j \beta z}  \tag{IV.9}\\
& \left.E_{\phi}=\frac{j \omega \mu}{k_{c}}[B \sin \phi+C \cos \phi]\right]_{1}^{\prime}\left(k_{c} \rho\right) e^{-j \beta z}  \tag{IV.10}\\
& H_{z}=[B \sin \phi+C \cos \phi] J_{1}\left(k_{c} \rho\right) e^{-j \beta z}  \tag{IV.11}\\
& H_{\rho}=-\frac{j \beta}{k_{c}}[B \sin \phi+C \cos \phi] J_{1}{ }^{\prime}\left(k_{c} \rho\right) e^{-j \beta z}  \tag{IV.12}\\
& H_{\phi}=-\frac{j \beta}{k_{c}^{2} \rho}[B \cos \phi-C \sin \phi] J_{1}\left(k_{c} \rho\right) e^{-j \beta z} \tag{IV.13}
\end{align*}
$$

These equations are in a coordinate system with the z axis aligned with the central line of the waveguide. Since the rectangular waveguide fields are in Cartesian coordinates and the circular waveguide fields are in cylindrical coordinates, a conversion of the circular waveguide fields to Cartesian is needed:

The circular waveguide fields in the Cartesian coordinate system can be written as follows:
$E_{z}=0$

$$
\begin{equation*}
H_{z}=H_{z} \tag{IV.14}
\end{equation*}
$$

$E_{x}=E_{\rho} \cos \phi-E_{\phi} \sin \phi$
$H_{x}=H_{\rho} \cos \phi-H_{\phi} \sin \phi$
$E_{y}=E_{\rho} \sin \phi+E_{\phi} \cos \phi$

$$
\begin{equation*}
H_{y}=H_{\rho} \sin \phi+H_{\phi} \cos \phi \tag{IV.18}
\end{equation*}
$$

with $\rho=\sqrt{x^{2}+y^{2}}$ and $\phi=\tan ^{-1} \frac{y}{x}$

This leads to the following field equations:

$$
\begin{align*}
& E_{z}=0  \tag{IV.20}\\
& E_{x}=-\frac{B j \omega \mu}{k_{c}}\left[\frac{\left(\cos \left(\tan ^{-1} \frac{y}{x}\right)\right)^{2}}{k_{c} \sqrt{x^{2}+y^{2}}} J_{1}\left(k_{c} \sqrt{x^{2}+y^{2}}\right)+\left(\sin \left(\tan ^{-1} \frac{y}{x}\right)\right)^{2} J_{1}^{\prime}\left(k_{c} \sqrt{x^{2}+y^{2}}\right)\right] e^{-j \beta z}+ \\
& \frac{C j \omega \mu}{k_{c}}\left[\sin \left(\tan ^{-1} \frac{y}{x}\right) \cos \left(\tan ^{-1} \frac{y}{x}\right)\left(\frac{J_{1}\left(k_{c} \sqrt{x^{2}+y^{2}}\right)}{k_{c} \sqrt{x^{2}+y^{2}}}-J_{1}{ }^{\prime}\left(k_{c} \sqrt{x^{2}+y^{2}}\right)\right)\right] e^{-j \beta z}  \tag{IV.21}\\
& E_{y}=\frac{j B \omega \mu}{k_{c}}\left[\sin \left(\tan ^{-1} \frac{y}{x}\right) \cos \left(\tan ^{-1} \frac{y}{x}\right)\left(J_{1}^{\prime}\left(k_{c} \sqrt{x^{2}+y^{2}}\right)-\frac{J_{1}\left(k_{c} \sqrt{x^{2}+y^{2}}\right)}{k_{c} \sqrt{x^{2}+y^{2}}}\right)\right] e^{-j \beta z}+ \\
& \frac{c j \omega \mu}{k_{c}}\left[\frac{\left(\sin \left(\tan ^{-1} \frac{y}{x}\right)\right)^{2}}{k_{c} \sqrt{x^{2}+y^{2}}} J_{1}\left(k_{c} \sqrt{x^{2}+y^{2}}\right)+\left(\cos \left(\tan ^{-1} \frac{y}{x}\right)\right)^{2} J_{1}^{\prime}\left(k_{c} \sqrt{x^{2}+y^{2}}\right)\right] e^{-j \beta z}  \tag{IV.22}\\
& H_{x}=\frac{j \beta B}{k_{c}}\left[\sin \left(\tan ^{-1} \frac{y}{x}\right) \cos \left(\tan ^{-1} \frac{y}{x}\right)\left(\frac{J_{1}\left(k_{c} \sqrt{x^{2}+y^{2}}\right)}{k_{c} \sqrt{x^{2}+y^{2}}}-J_{1}^{\prime}\left(k_{c} \sqrt{x^{2}+y^{2}}\right)\right)\right] e^{-j \beta z}- \\
& \frac{j \beta C}{k_{c}}\left[\frac{\left(\sin \left(\tan ^{-1} \frac{y}{x}\right)\right)^{2}}{k_{c} \sqrt{x^{2}+y^{2}}} J_{1}\left(k_{c} \sqrt{x^{2}+y^{2}}\right)+\left(\cos \left(\tan ^{-1} \frac{y}{x}\right)\right)^{2} J_{1}^{\prime}\left(k_{c} \sqrt{x^{2}+y^{2}}\right)\right] e^{-j \beta z}  \tag{IV.23}\\
& H_{y}=-\frac{j \beta B}{k_{c}}\left[\frac{\left(\cos \left(\tan ^{-1} \frac{y}{x}\right)\right)^{2}}{k_{c} \sqrt{x^{2}+y^{2}}} J_{1}\left(k_{c} \sqrt{x^{2}+y^{2}}\right)+\right. \\
& \left.\left(\sin \left(\tan ^{-1} \frac{y}{x}\right)\right)^{2} J_{1}^{\prime}\left(k_{c} \sqrt{x^{2}+y^{2}}\right)\right] e^{-j \beta z}+ \\
& \frac{c j \beta}{k_{c}}\left[\sin \left(\tan ^{-1} \frac{y}{x}\right) \cos \left(\tan ^{-1} \frac{y}{x}\right)\left(\frac{J_{1}\left(k_{c} \sqrt{x^{2}+y^{2}}\right)}{k_{c} \sqrt{x^{2}+y^{2}}}-J_{1}{ }^{\prime}\left(k_{c} \sqrt{x^{2}+y^{2}}\right)\right)\right] e^{-j \beta z}  \tag{IV.24}\\
& H_{z}=\left[B \sin \left(\tan ^{-1} \frac{y}{x}\right)+C \cos \left(\tan ^{-1} \frac{y}{x}\right)\right] e^{-j \beta z} \tag{IV.25}
\end{align*}
$$

The circular waveguide is centered at $\mathrm{x}=\mathrm{s}, \mathrm{y}=\mathrm{b}$ and $\mathrm{z}=0$. This requires certain transformations to be applied to the circular waveguide fields in order to match the case being solved in here and write everything in the initial coordinate system. The resulting fields of the $\mathrm{TE}_{11}$ mode propagating in the circular waveguide after the transformations are as follows:

$$
\begin{align*}
& E_{x}=\frac{j B \omega \mu}{k_{c}}\left[\operatorname { s i n } ( \operatorname { t a n } ^ { - 1 } \frac { x - s } { z } ) \operatorname { c o s } ( \operatorname { t a n } ^ { - 1 } \frac { x - s } { z } ) \left(J_{1}^{\prime}\left(k_{c} \sqrt{(x-s)^{2}+z^{2}}\right)-\right.\right. \\
& \left.\left.\frac{J_{1}\left(k_{c} \sqrt{(x-s)^{2}+z^{2}}\right)}{k_{c} \sqrt{(x-s)^{2}+z^{2}}}\right)\right] e^{-j \beta(y-b)}+\frac{j \omega \mu C}{k_{c}}\left[\frac{\left(\sin \left(\tan ^{-1} \frac{x-s}{z}\right)\right)^{2}}{k_{c} \sqrt{(x-s)^{2}+z^{2}}} J_{1}\left(k_{c} \sqrt{(x-s)^{2}+z^{2}}\right)+\right. \\
& \left.\left(\cos \left(\tan ^{-1} \frac{x-s}{z}\right)\right)^{2} J_{1}^{\prime}\left(k_{c} \sqrt{(x-s)^{2}+z^{2}}\right)\right] e^{-j \beta(y-b)}  \tag{IV.26}\\
& E_{y}=0  \tag{IV.27}\\
& E_{z}=-\frac{B j \omega \mu}{k_{c}}\left[\frac{\left(\cos \left(\tan ^{-1} \frac{x-s}{z}\right)\right)^{2}}{k_{c} \sqrt{(x-s)^{2}+z^{2}}} J_{1}\left(k_{c} \sqrt{(x-s)^{2}+z^{2}}\right)+\right. \\
& \left.\left(\sin \left(\tan ^{-1} \frac{x-s}{z}\right)\right)^{2} J_{1}^{\prime}\left(k_{c} \sqrt{(x-s)^{2}+z^{2}}\right)\right] e^{-j \beta(y-b)}+ \\
& \frac{c j \omega \mu}{k_{c}}\left[\operatorname { s i n } ( \operatorname { t a n } ^ { - 1 } \frac { x - s } { z } ) \operatorname { c o s } ( \operatorname { t a n } ^ { - 1 } \frac { x - s } { z } ) \left(\frac{J_{1}\left(k_{c} \sqrt{(x-s)^{2}+z^{2}}\right)}{k_{c} \sqrt{(x-s)^{2}+z^{2}}}-\right.\right. \\
& \left.\left.J_{1}{ }^{\prime}\left(k_{c} \sqrt{(x-s)^{2}+z^{2}}\right)\right)\right] e^{-j \beta(y-b)}  \tag{IV.28}\\
& H_{x}=-\frac{j \beta B}{k_{c}}\left[\frac{\left(\cos \left(\tan ^{-1} \frac{x-s}{z}\right)\right)^{2}}{k_{c} \sqrt{(x-s)^{2}+z^{2}}} J_{1}\left(k_{c} \sqrt{(x-s)^{2}+z^{2}}\right)+\right. \\
& \left.\left(\sin \left(\tan ^{-1} \frac{x-s}{z}\right)\right)^{2} J_{1}^{\prime}\left(k_{c} \sqrt{(x-s)^{2}+z^{2}}\right)\right] e^{-j \beta(y-b)}+ \\
& \frac{c j \beta}{k_{c}}\left[\operatorname { s i n } ( \operatorname { t a n } ^ { - 1 } \frac { x - s } { z } ) \operatorname { c o s } ( \operatorname { t a n } ^ { - 1 } \frac { x - s } { z } ) \left(\frac{J_{1}\left(k_{c} \sqrt{(x-s)^{2}+z^{2}}\right)}{k_{c} \sqrt{(x-s)^{2}+z^{2}}}-\right.\right. \\
& \left.\left.J_{1}{ }^{\prime}\left(k_{c} \sqrt{(x-s)^{2}+z^{2}}\right)\right)\right] e^{-j \beta(y-b)}  \tag{IV.29}\\
& H_{y}=\left[B \sin \left(\tan ^{-1} \frac{x-s}{z}\right)+C \cos \left(\tan ^{-1} \frac{x-s}{z}\right)\right] e^{-j \beta(y-b)} \tag{IV.30}
\end{align*}
$$

$$
\begin{gather*}
H_{z}=\frac{j \beta B}{k_{c}}\left[\operatorname { s i n } ( \operatorname { t a n } ^ { - 1 } \frac { x - s } { z } ) \operatorname { c o s } ( \operatorname { t a n } ^ { - 1 } \frac { x - s } { z } ) \left(\frac{J_{1}\left(k_{c} \sqrt{(x-s)^{2}+z^{2}}\right)}{k_{c} \sqrt{(x-s)^{2}+z^{2}}}-\right.\right. \\
\left.\left.J_{1}^{\prime}\left(k_{c} \sqrt{(x-s)^{2}+z^{2}}\right)\right)\right] e^{-j \beta(y-b)}-\frac{j \beta C}{k_{c}}\left[\frac{\left(\sin \left(\tan ^{-1} \frac{x-s}{z}\right)\right)^{2}}{k_{c} \sqrt{(x-s)^{2}+z^{2}}} J_{1}\left(k_{c} \sqrt{(x-s)^{2}+z^{2}}\right)+\right. \\
\left.\left(\cos \left(\tan ^{-1} \frac{x-s}{z}\right)\right)^{2} J_{1}^{\prime}\left(k_{c} \sqrt{(x-s)^{2}+z^{2}}\right)\right] e^{-j \beta(y-b)} \tag{IV.31}
\end{gather*}
$$

Since the magnitude of the components of the E-field are important in creating a circularly polarized wave, $B$ and $C$ should be calculated. In order to get $B$ and $C$, the following equations are solved:

$$
\begin{align*}
& B=\frac{1}{P_{11}} \int_{v} \overrightarrow{E_{1}} \cdot \vec{J} d v-\frac{1}{P_{11}} \int_{v} \overrightarrow{H_{1}} \cdot \vec{M} d v  \tag{IV.32}\\
& C=\frac{1}{P_{11}} \int_{v} \overrightarrow{E_{2}} \cdot \vec{J} d v-\frac{1}{P_{11}} \int_{v} \overrightarrow{H_{2}} \cdot \vec{M} d v \tag{IV.33}
\end{align*}
$$

In this case:
$\vec{E} \cdot \vec{J}=0$ and $\vec{H} \cdot \vec{M}$ is non-zero.

This leads the expressions of B and C to be reduced to the following:

$$
\begin{align*}
& B=H_{1 x}(s, b, 0) \times M_{x}+H_{1 z}(s, b, 0) \times M_{z}  \tag{IV.34}\\
& C=H_{2 x}(s, b, 0) \times M_{x}+H_{2 z}(s, b, 0) \times M_{z} \tag{IV.35}
\end{align*}
$$

where:

$$
\begin{gather*}
H_{1 x}=-\frac{j \beta}{k_{c}}\left[\frac{\left(\cos \left(\tan ^{-1} \frac{x-s}{z}\right)\right)^{2}}{k_{c} \sqrt{(x-s)^{2}+z^{2}}} J_{1}\left(k_{c} \sqrt{(x-s)^{2}+z^{2}}\right)+\right. \\
\left.\left(\sin \left(\tan ^{-1} \frac{x-s}{z}\right)\right)^{2} J_{1}^{\prime}\left(k_{c} \sqrt{(x-s)^{2}+z^{2}}\right)\right] e^{-j \beta(y-b)} \tag{IV.36}
\end{gather*}
$$

$$
\begin{align*}
& H_{2 x}=\frac{j \beta}{k_{c}}\left[\operatorname { s i n } ( \operatorname { t a n } ^ { - 1 } \frac { x - s } { z } ) \operatorname { c o s } ( \operatorname { t a n } ^ { - 1 } \frac { x - s } { z } ) \left(\frac{J_{1}\left(k_{c} \sqrt{(x-s)^{2}+z^{2}}\right)}{k_{c} \sqrt{(x-s)^{2}+z^{2}}}-\right.\right. \\
& \left.\left.J_{1}^{\prime}\left(k_{c} \sqrt{(x-s)^{2}+z^{2}}\right)\right)\right] e^{-j \beta(y-b)}  \tag{IV.37}\\
& H_{1 z}=\frac{j \beta}{k_{c}}\left[\operatorname { s i n } ( \operatorname { t a n } ^ { - 1 } \frac { x - s } { z } ) \operatorname { c o s } ( \operatorname { t a n } ^ { - 1 } \frac { x - s } { z } ) \left(\frac{J_{1}\left(k_{c} \sqrt{(x-s)^{2}+z^{2}}\right)}{k_{c} \sqrt{(x-s)^{2}+z^{2}}}-\right.\right. \\
& \left.\left.J_{1}^{\prime}\left(k_{c} \sqrt{(x-s)^{2}+z^{2}}\right)\right)\right] e^{-j \beta(y-b)}  \tag{IV.38}\\
& \quad H_{2 z}=-\frac{j \beta}{k_{c}}\left[\frac{\left(\sin ^{\left.\left(\tan ^{-1} \frac{x-s}{z}\right)\right)^{2}}\right.}{k_{c} \sqrt{(x-s)^{2}+z^{2}}} J_{1}\left(k_{c} \sqrt{(x-s)^{2}+z^{2}}\right)+\right. \\
& \left.\left(\cos \left(\tan ^{-1} \frac{x-s}{z}\right)\right)^{2} J_{1}^{\prime}\left(k_{c} \sqrt{(x-s)^{2}+z^{2}}\right)\right] e^{-j \beta(y-b)} \tag{IV.39}
\end{align*}
$$

To simplify these expressions, the small argument approximation for the Bessel function is used:

$$
\begin{aligned}
& J_{1}(l)=\frac{l}{2} \\
& J_{1}^{\prime}(l)=\frac{1}{2}
\end{aligned}
$$

These approximations stand because, in this scenario, the Bessel function is being evaluated in the proximity of zero. The resulting simplified fields are as follows:

$$
\begin{align*}
& \left.H_{1 x}(s, b, 0)=-\frac{j \beta}{k_{c}}\left[\frac{\left(\cos \left(\tan ^{-1} \frac{x-s}{z}\right)\right)^{2}}{2}+\left(\sin \left(\tan ^{-1} \frac{x-s}{z}\right)\right)^{2} \frac{1}{2}\right)\right]=-\frac{j \beta}{2 k_{c}}  \tag{IV.40}\\
& H_{2 x}(s, b, 0)=\frac{j \beta}{k_{c}}\left[\sin \left(\tan ^{-1} \frac{x-s}{z}\right) \cos \left(\tan ^{-1} \frac{x-s}{z}\right)\left(\frac{1}{2}-\frac{1}{2}\right)\right]=0 \tag{IV.41}
\end{align*}
$$

$$
\begin{align*}
& H_{1 z}(s, b, 0)=\frac{j \beta}{k_{c}}\left[\sin \left(\tan ^{-1} \frac{x-s}{z}\right) \cos \left(\tan ^{-1} \frac{x-s}{z}\right)\left(\frac{1}{2}-\frac{1}{2}\right)\right]=0  \tag{IV.42}\\
& H_{2 z}(s, b, 0)=-\frac{j \beta}{k_{c}}\left[\frac{\left(\sin \left(\tan ^{-1} \frac{x-s}{z}\right)\right)^{2}}{2}+\left(\cos \left(\tan ^{-1} \frac{x-s}{z}\right)\right)^{2} \frac{1}{2}\right]=-\frac{j \beta}{2 k_{c}} \tag{IV.43}
\end{align*}
$$

Solving for B and C using the simplified expressions leads to the following equations:

$$
\begin{align*}
& B=-\frac{\beta \omega \mu_{0} \alpha_{m} A}{P_{11} k_{c} Z_{10}} \sin \frac{\pi s}{a}  \tag{IV.44}\\
& C=\frac{j \omega \mu_{0} \alpha_{m} A \pi}{P_{11} a Z_{10} k_{c}} \cos \frac{\pi s}{a} \tag{IV.45}
\end{align*}
$$

In order to get circular polarization, we need $|\mathrm{B}|=|\mathrm{C}|$ and $\angle B=\angle C \pm \frac{\pi}{2}$. The expressions of $B$ and $C$ show that $B$ is real and $C$ is imaginary. This means that they are $90^{\circ}$ out of phase, hence satisfying the phase condition.

In order for $|B|=|C|$, the following should be satisfied:
$\left|-\frac{\beta \omega \mu_{0} \alpha_{m} A}{P_{11} k_{c} Z_{10}} \sin \frac{\pi s}{a}\right|=\left|\frac{j \omega \mu_{0} \alpha_{m} A \pi}{P_{11} a Z_{10} k_{c}} \cos \frac{\pi s}{a}\right|$
$\Rightarrow \tan \frac{\pi s}{a}=\frac{\pi}{\beta a}$

Hence,

$$
\begin{equation*}
s=\frac{a}{2} \pm \frac{a}{\pi} \tan ^{-1} \sqrt{\left(\frac{2 f a}{c}\right)^{2}-1} \tag{IV.46}
\end{equation*}
$$

## - Design

In order to verify the theoretical calculations presented in the previous section, a WR12 rectangular waveguide with width $a=3.1 \mathrm{~mm}$ and height $b=1.55 \mathrm{~mm}$ is used with
a circular slot on its broad-wall as shown in Figure IV. 1 (a). By using the formula for $s$ that was derived in the previous section, the slot should be centered at $s=0.7272 \mathrm{~mm}$. At this position, the electric field radiated into the circular waveguide is circularly polarized at $f=72 \mathrm{GHz}$. The slot feeds a circular waveguide in its center as shown in Figure IV. 1 (a). The design is simulated using CST MWS.

An analysis of the slot dimensions and shape is performed in order to determine an efficient extraction of power from the rectangular waveguide. When the slot radius is increased, the efficiency of the extraction increases, and more power reaches the end of the circular waveguide (port 3) as shown in Figure IV.1. It is noticed that the efficiency of circular slots at the frequency of interest is very low. For a slot diameter $D_{s}=1.4544 \mathrm{~mm}$ the maximum efficiency of the system is $0.88 \%$ at 72 GHz and it increases to $8.4 \%$ at 90 GHz . This means that circular slots are an inefficient way to extract the power from the rectangular waveguide.


Figure IV.2: Plot of the transmission between ports 1 and 3 for different circular slots and a cross slot operating at 72 GHz .

In order to improve the efficiency of the extraction, cross slots are used as can be seen in Figure IV.2. In what follows, the Z-shaped arm cross slot that is discussed in the previous chapter is used. The frequency of operation of the slot extraction relies on the physical dimensions of the slot. The total arm length of the slot has to be $L_{a} \approx \lambda_{g} / 2$ and the arm width has to be chosen so that the maximum power is extracted from the waveguide while taking into account the restrictions in fabrication. For the design to operate at 72 GHz , the slot dimensions are chosen as stated in Table III.1. It is noticed that the diameter of the circular waveguide affects the efficiency of the extraction because it changes the wave impedance and hence the loading on the slot. By changing the diameter of the circular waveguide, the wave impedance inside the waveguide is changed. If the diameter is chosen in a way where the cutoff frequency is less than 72 GHz and the wave impedance at the center of the waveguide matches the output impedance of the slot extraction, high
efficiency values can be achieved and more power reaches port 3 while minimum power reaches port 2 of the extraction. By using a circular waveguide diameter of $D_{w}=2.5 \mathrm{~mm}$, maximum power reaches port 3 when a cross slot is used.

Table IV.1: Dimensions of the Z-shaped arm cross-slot at 72 GHz .

| Parameter | Description | Value |
| :---: | :---: | :---: |
| $L_{z 1}$ | long arm length | 1.85 mm |
| $L_{z 2}$ | short arm length | 0.46 mm |
| $W_{a}$ | arm width | 0.15 mm |
| $S$ | slot center position | 0.69 mm |
| $\alpha_{2}$ | arm angle with y-axis | $43.5^{\circ}$ |
| $t_{w 3}$ | slot thickness | 0.4 mm |

The use of the cross slot increases the efficiency of the extraction to more than $70 \%$ in the frequency band $71.71-72.8 \mathrm{GHz}$, reaching its peak efficiency of $83.86 \%$ at 72.16 GHz . The metal used for this design is Aluminum and it causes around $15 \%$ of the total power loss in the system at 72.16 GHz . The ohmic losses can be reduced to $12 \%$ by reducing the lengths of the rectangular waveguide and the circular waveguide. The system is well matched over the entire E-band, with an $S_{11}<-18 d B$ for the frequency band 68-88 GHz as shown in Figure IV.3. The isolation between the two feeding ports of the rectangular waveguide reaches $\left|S_{21}\right|=20 \mathrm{~dB}$ at 72.16 GHz and is more than 10 dB for the frequency range $71.8-72.6 \mathrm{GHz}$ as can be seen from the $S_{21}$ plot in Figure IV.3. At the end of the circular waveguide (port 3), two orthogonal $\mathrm{TE}_{11}$ modes are detected. These two modes are excited by the cross slot and have almost equal powers in the frequency range
71.5-88 GHz as seen in the $\left|S_{31}\right|$ plot in Figure IV.4. However, the power received by port 3 is not flat for the entire frequency range. At least half of the power fed through port 1 is received by port 3 in the frequency range $71.43-73.35 \mathrm{GHz}$. The maximum power that can be received by port 3 coming from port 1 is at 72.16 GHz where $\left|S_{31(1)}\right|=$ $-3.88 d B$ and $\left|S_{31(2)}\right|=-3.67 d B$. The phase difference between $S_{31(1)}$ and $S_{31(2)}$ is in the range of $89^{\circ}<\angle S_{31(2)}-\angle S_{31(1)}<95.4^{\circ}$ in the frequency range $71-74 \mathrm{GHz}$ as shown in Figure IV.5.

These results show that this device can work as a wave polarizer with an efficiency of more than $70 \%$ for the frequency range of $71.71-72.6 \mathrm{GHz}$ with an isolation between the two feeding ports of more than 10 dB . The polarizer will create a LHCP wave in the circular waveguide if fed by port 1 , and a RHCP wave if fed by port 2 . If both ports 1 and 2 are fed with equal powers, the resulting wave is linearly polarized. In the next section, the testing of this device with a horn antenna is illustrated.


Figure IV.3: Plot of the $S_{11}$ and $S_{21}$ over frequency for the Z-arm shape cross slot extraction.


Figure IV.4: Plot of the power received at the end of the circular waveguide in the form of $\mathrm{TE}_{11}$ mode.


Figure IV.5: Plot of the unwrapped phase of the two $\mathrm{TE}_{11}$ modes at the end of the circular waveguide.

## III. A Single Slot Feeding a Conical Horn Antenna

In order to test its performance, a polarizer is used to feed a horn antenna that can radiate a circularly polarized wave. Two versions of the polarizer at 72 GHz and 84 GHz are designed and modified to be matched with the input of the conical horn.

## $\square$ Horn Antenna Design

In order to test the polarizer, a conical horn antenna is used. This antenna has a symmetrical nature that allows it to support a circularly polarized wave. A polarizer is used to excite two $\mathrm{TE}_{11}$ modes with orthogonal electric fields at the feeding waveguide input with a phase difference of $90^{\circ}$.

The design of conical horns should follow the guidelines set by A. P. King in [60]. These guidelines determine an optimum horn line that will lead to the best gains of the horn. This line identifies the best gain that could be achieved as a function of different horn aperture diameters, for different horn lengths L . The horn parameters L and $d_{m}$ along with the optimum gain achieved are illustrated in Figure IV.6. A numerical equation relating the optimum directivity of the horn to its geometrical dimensions is also derived and a relationship between directivity $D_{c}$, horn length L and aperture diameter $d_{m}$ is established leading to the following equations [61]:

$$
\begin{equation*}
\left(D_{c}\right)_{o p t} \approx 15.9749\left(\frac{L}{\lambda}\right)+1.7209 \tag{IV.47}
\end{equation*}
$$

and

$$
\begin{equation*}
L \approx 0.3232\left(\frac{d_{m}}{\lambda}\right)^{2}-0.0475\left(\frac{d_{m}}{\lambda}\right)+0.0052 \tag{IV.48}
\end{equation*}
$$

Since the applications at W/V-band that are targeted in this dissertation are either satellite communication links, point-to-point terrestrial links, or radar systems, and since the free space path loss at these frequencies is very high, a minimum gain of 20 dB is required for the horn. As shown in Figure IV.6, for a conical horn to have a gain $\left(\mathrm{G}_{\mathrm{h}}\right)$ more than 20 dB , the length of the horn (L) should be around $10 \lambda$, which makes the gain curve for the different diameters well above 20 dB , and the aperture diameter $\left(\mathrm{d}_{\mathrm{m}}\right)$ should be between $5 \lambda$ and $6 \lambda$. These dimensions at 72 GHz translate to $\mathrm{L}=41.67 \mathrm{~mm}$ and $\mathrm{d}_{\mathrm{m}}=25$ mm .

An off the shelf conical horn with dimensions close to the ones derived earlier is used to test the polarizer. The horn has $L=42.6 \mathrm{~mm}$ and $d_{m}=25 \mathrm{~mm}$. The circular waveguide feed of the horn has a diameter $d=3.175 \mathrm{~mm}$ in order for the horn to be operational over the mid-range of the E-band. Due to the dimensions of the circular waveguide, the cutoff frequency for $\mathrm{TE}_{11}$ inside the waveguide is at:

$$
\begin{equation*}
f_{c, T E_{11}}=\frac{1.841 \times c}{d \pi}=55.37 \mathrm{GHz} \tag{IV.49}
\end{equation*}
$$

At the same time, the cutoff frequency of $\mathrm{TM}_{01}$ is at:

$$
\begin{equation*}
f_{c, T M_{01}}=\frac{2.405 \times c}{d \pi}=72.33 \mathrm{GHz} \tag{IV.50}
\end{equation*}
$$

This cutoff frequency is very close to the frequency of operation of the system being designed. When the design is fabricated, the fabrication errors might cause the $\mathrm{TM}_{01}$ mode to be excited. If the antenna is directly connected to the cross slot, the excited $\mathrm{TM}_{01}$ causes the output wave to lose its circular polarization and causes the entire design to have a low efficiency. In order to avoid that, a circular waveguide with narrower diameter is used to
feed the horn. The design details are explained in the next sub-section. The horn's input is well matched from 60 GHz to 90 GHz if fed without adding the extra waveguide as shown in Figure IV.7.


Figure IV.6: The absolute gain of a conical horn as a function of aperture diameter $\left(d_{m} / \lambda\right)$ for a series of axial lengths, L. [60]


Figure IV.7: Conical horn reflection coefficient.

## - Full System Design

The conical horn is fed using the cross polarizer as shown in Figure IV.8. In order to suppress higher order modes and only propagate the fundamental $\mathrm{TE}_{11}$ mode inside the circular waveguide feed of the horn, another circular waveguide with a narrower diameter is added between the cross slot and the horn system. This new waveguide plays a dual role of making the extraction of power from the rectangular waveguide more efficient, and limiting the modes being fed to the conical horn.

In order to have a good matching in the system and good efficiency, a transition between the two circular waveguides is added. The transition is made of a tapered circular waveguide. Two different designs, one at 72 GHz and one at 84 GHz are done. The 72 GHz design has a cross slot on the broad-wall of a WR-12 waveguide $(\mathrm{a}=3.1 \mathrm{~mm}$ and $\mathrm{b}=1.55 \mathrm{~mm}$ ), while the 84 GHz design has the cross slot on the broad-wall of a WR-10
waveguide $(\mathrm{a}=2.54 \mathrm{~mm}$ and $\mathrm{b}=1.27 \mathrm{~mm})$. The different design dimensions for both frequencies are indicated in Table IV.2.

Table IV.2: Design dimensions for the polarizer at 72 GHz and 84 GHz .

| Parameter | Description | Value at 72 GHz | Value at 84 GHz |
| :---: | :---: | :---: | :---: |
| $L_{z 1}$ | long arm length | 1.85 mm | 1.65 mm |
| $L_{z 2}$ | short arm length | 0.46 mm | 0.36 mm |
| $W_{a}$ | arm width | 0.15 mm | 0.15 mm |
| $S$ | slot center position | 0.69 mm | 0.69 mm |
| $\alpha_{2}$ | arm angle with y-axis | $43.5^{\circ}$ | $43.5^{\circ}$ |
| $t_{w 3}$ | slot thickness | 0.4 mm | 0.4 mm |
| $D_{c}$ | circular waveguide diameter | 2.5 mm | 2.2 mm |
| $L_{c 1}$ | circular waveguide length | 8 mm | 8 mm |
| $L_{T}$ | tapered waveguide length | 2 mm | 8 mm |
| diameter of the horn input | waveguide | 3.175 mm | 3.175 mm |
|  |  |  |  |

- 72 GHz model:

The system designed at 72 GHz has an input reflection coefficient of $\mathrm{S}_{11}<-15 \mathrm{~dB}$ over the frequency range $70-86 \mathrm{GHz}$ for both feeding ports. The isolation $\mathrm{S}_{21}$ between ports 1 and 2 reaches -20 dB at 71.76 GHz which is the main resonance frequency of the slot. Another resonance for the slot caused by the different design components is at 73 GHz where the isolation is $\left|S_{21}\right|=9.5 d B$ as shown in Figure IV.9. The two resonances in the design have increased the bandwidth over which the efficiency of the design is higher than $70 \%$.


Figure IV.8: Single slot feeding a conical horn antenna.

The system produces a LHCP gain if fed through port 1, and a RHCP gain if fed through port 2. The system has a total CP gain between 21.25 dB and 21.8 dB over the frequency range $71.6-73.2 \mathrm{GHz}$ for both LHCP and RHCP . In this frequency range, the total efficiency of the system $70 \%<\varepsilon_{h}<80 \% . \varepsilon_{h}$ reaches its peak of $80 \%$ at 71.8 and 73 GHz which are the two resonant frequencies of the slot, and the system gain is 21.8 dB at these frequencies as shown in Figure IV. 10 (a). The axial ratio $\left(A R_{72}\right)$ of the system in the direction of the maximum gain, $\theta=0^{\circ}$, is $A R_{72}<2 d B$ over the frequency range $71.6-$ 74 GHz and it reaches a minimum of $A R_{72}=0.125 d B$ at 73 GHz as shown in Figure IV. 10 (b). These results show that the system has a good efficiency, gain and axial ratio
over the frequency range $71.6-73.2 \mathrm{GHz}$ giving it a 1.6 GHz bandwidth of operation as a polarizer that can produce both LHCP and RHCP.

The system produces a $H P B W=11.5^{\circ}$ in both planes $\phi=0^{\circ}$ and $\phi=90^{\circ}$. The entire HPBW of the system is circularly polarized with an $A R<3 d B$ for the frequency range $71.8-73.2 \mathrm{GHz}$. At 73 GHz , the system produces an $A R<2 d B$ for the entire HPBW. The cross polarization discrimination is higher than 18 dB for the entire HPBW as can be seen in Figure IV. 11 (a) and (b) for both planes $\phi=0^{\circ}$ and $\phi=90^{\circ}$. It can be noticed that the polarizer did not affect the HPBW of the horn, in addition, in the plane $\phi=90^{\circ}$, the system has reduced the SLL of the horn from -19 dB to -23.9 dB .

- 84 GHz model:

The system designed at 84 GHz has an input reflection coefficient $S_{11}<-19 \mathrm{~dB}$ for all frequencies between $70-86 \mathrm{GHz}$ for both feeding ports. The isolation $\left|S_{21}\right|>20 \mathrm{~dB}$ for a bandwidth of 250 MHz over the frequency range of $83.78-84.03 \mathrm{GHz}$. The maximum isolation between ports 1 and 2 reaches $\left|S_{21}\right|=33 d B$ at 83.906 GHz where the resonant frequency of the slot is, as shown in Figure IV.9. By having one resonance for the slot, the system has a narrower bandwidth where the efficiency, $\varepsilon_{h}>70 \%$, but the isolation between the ports is improved for a wider frequency range.

The system efficiency, $\varepsilon_{h}$, in this case is higher than $70 \%$ for the frequency band 83.4 -84.65 GHz as shown in Figure IV. 10 (a). The maximum efficiency that the system can reach is $\varepsilon_{h}=83 \%$ at 84 GHz with a CP gain of 22.57 dB . The axial ratio $\left(A R_{84}\right)$ of the system in the direction of the maximum gain, $\theta=0^{\circ}$, is $A R_{72}<1 d B$ over the frequency range $83-85 \mathrm{GHz}$ and it reaches a minimum of $A R_{84}=0.2 \mathrm{~dB}$ at 83 GHz as shown in

Figure IV. 10 (b). These results show that the system has a good efficiency, gain and axial ratio over the frequency range of $83.4-84.65 \mathrm{GHz}$ giving it a 1.25 GHz bandwidth of operation as a polarizer that can produce both LHCP and RHCP. In this case the bandwidth of operation is less than in the case of the 72 GHz polarizer, due to the system having only one resonant frequency.


Figure IV.9: Reflection coefficient and the isolation between both ports for both 72 and 84 GHz models.

(a)

(b)

Figure IV.10: (a) The maximum gain of the 72 and 84 GHz designs vs frequency, compared to the linear gain of the horn, (b) the axial ratio of the 72 and 84 GHz designs vs frequency.


Figure IV.11: Gain pattern at 73 GHz , (a) in $\phi=0^{\circ}$, (b) in $\phi=90^{\circ}$.

## - Fabrication Results

The 72 GHz system is fabricated and tested. The system assembly is shown in Figure IV.12. The cross slot is fabricated on two sheets of copper with thickness 0.2 mm each using LPKF ProtoLaser S4 laser cutting machine. The copper sheets are stacked and shown
as (2) in Figure IV.12. The circular waveguides along with the tapered waveguide are fabricated in a separate aluminum block, component (1) in Figure IV.12. This block contains an area where the sheets with the cross slot can fit and be flat with the bottom of the block. Finally, the WR-12 waveguide channel is fabricated using mechanical milling onto a separate aluminum block (3) in Figure IV.12. The two different aluminum blocks sandwich the copper sheet, and the whole design is aligned using alignment pins. The whole design is screwed together and the gaps between the different blocks are suppressed to the order of $1 \mu \mathrm{~m}$ which doesn't affect the functionality of the fabricated polarizer.


Top View


Bottom View

Figure IV.12: Polarizer assembling technique.

The polarizer is tested for s-parameters while a conical horn is connected to it and while the circular waveguide output is left open as shown in Figure IV. 13 (a). Both scenarios lead to the same results, proving that the polarizer parameters are barely affected by the device connected to the circular waveguide output. The measured results show that the
resonant frequency of the slot is shifted down by 500 MHz and occurs at 71.25 GHz as shown in Figure IV.14. This shift is caused by the fabrication errors caused by the laser prototyping machine that is used to cut out the slots leaving a slot width larger on the top level than the bottom level of the copper sheet. The reflection coefficient of both feeding ports is less than -14 dB for the entire operating frequency range.


Figure IV.13: (a) The conical horn attached to the polarizer, (b) The measurement setup for the sparameters of the polarizer.

In an attempt to measure the E-field power and phase of the two orthogonal $\mathrm{TE}_{11}$ modes at the output of the circular waveguide, a rectangular waveguide with a waveguide bend is connected to the circular waveguide output directly as shown in Figure IV. 13 (b). The connection shown allows the measurement of one of the modes and then the rectangular waveguide is rotated by $90^{\circ}$ to measure the other mode. The direct connection causes reflections at the output of the circular waveguide, which makes the power reading
inaccurate, but still allows to measure the power difference and phase difference between the two excited modes.


Figure IV.14: Plot of the measured s-parameters of the polarizer.

The power measured by the two different setups of the rectangular waveguide shows that both modes have very close powers over the operating frequency range as shown in Figure IV. 15 and predicted by the simulated results. The power difference is at a maximum of 1.7 dB at around 72.6 GHz . The measured phase difference between the two orthogonal modes is between $90^{\circ}$ and $105^{\circ}$ in the operating frequency band as shown in Figure IV.16. The phase difference values might be affected by the connection of the rectangular waveguide to the circular waveguide without any transitions. These results show that a circularly polarized wave is present at the circular output of the polarizer.


Figure IV.15: Plots of the power received by port 3 from port 1 for two orthogonal $\mathrm{TE}_{11}$ modes.


Figure IV.16: Plots of the phase difference between the two orthogonal $\mathrm{TE}_{11}$ modes seen at port 3.

## IV. Slot Radiating Into a Square Waveguide

Square waveguides are another type of waveguides that could be used to extract circularly polarized waves from the rectangular waveguide. Square waveguides are symmetrical and can have $T E_{10}$ and $T E_{01}$ modes propagating in them. In this section the slot radiating into a square waveguide is studied.

## Theory

Assuming that a slot is on the broad-wall of a rectangular waveguide and coupling power into a square waveguide as shown in Figure IV. 17 (a).


Figure IV.17: (a) Cross slot on the broad-wall of a rectangular waveguide coupling power into a square waveguide, (b) $T E_{10}$ mode fields in a rectangular waveguide, (c) $T E_{10}$ and $T E_{01}$ modes Electric fields in a square waveguide.

In a rectangular waveguide, the fundamental mode propagating is $T E_{10}$ as shown in Figure IV. 17 (b). The field equations for this mode are:

$$
\begin{equation*}
E_{y}=A \sin \frac{\pi x}{a} e^{-j \beta z} \tag{IV.51}
\end{equation*}
$$

$$
\begin{align*}
& H_{x}=-\frac{A}{Z_{10}} \sin \frac{\pi x}{a} e^{-j \beta z}  \tag{IV.52}\\
& H_{z}=\frac{j \pi A}{\beta a z_{10}} \cos \frac{\pi x}{a} e^{-j \beta z} \tag{IV.53}
\end{align*}
$$

where $\quad Z_{10}=\frac{k_{0} \eta_{0}}{\beta}$

Assuming that the slot is at a distance s from the narrow wall and at $\mathrm{z}=0$, at the slot the electric and magnetic polarizabilities are:

$$
\begin{align*}
& \overrightarrow{P_{e}}=A \varepsilon_{0} \alpha_{e} \sin \frac{\pi s}{a} \delta(x-s) \delta(y-b) \delta(z) \hat{y}  \tag{IV.54}\\
& \overrightarrow{P_{m}}=-A \alpha_{m}\left[-\frac{\hat{x}}{z_{10}} \sin \frac{\pi s}{a}+\hat{z} \frac{j \pi}{\beta a Z_{10}} \cos \frac{\pi s}{a}\right] \delta(x-s) \delta(y-b) \delta(z) \tag{IV.55}
\end{align*}
$$

In the square waveguide, the two fundamental modes that could be excited are the $T E_{10}$ and the $T E_{01}$ modes with their electric fields are shown in Figure IV. 17 (c). The fields equations for these modes are as follows:

For $T E_{10}$ :

$$
\begin{array}{lll}
E_{x}=0 & \text { (IV.56) } & H_{x}=-\frac{B}{Z_{10}} \sin \frac{\pi\left(x-s+\frac{a_{1}}{2}\right)}{a_{1}} e^{-j \beta_{1}(y-b)} \\
E_{y}=0 & (\mathrm{IV} .58) & H_{y}=\frac{j \pi B}{\beta_{1} a_{1} Z_{10}} \cos \frac{\pi\left(x-s+\frac{a_{1}}{2}\right)}{a_{1}} e^{-j \beta_{1}(y-b)} \\
H_{z}=0 & \text { (IV.60) } & E_{z}=B \sin \frac{\pi\left(x-s+\frac{a_{1}}{2}\right)}{a_{1}} e^{-j \beta_{1}(y-b)}
\end{array}
$$

For $T E_{01}$ :
$E_{x}=C \sin \frac{\pi\left(z+\frac{a_{1}}{2}\right)}{a_{1}} e^{-j \beta_{2}(y-b)}$

$$
\begin{equation*}
H_{x}=0 \tag{IV.62}
\end{equation*}
$$

$$
\begin{array}{lll}
E_{y}=0 & (\text { IV.64 ) } & H_{y}=\frac{j \pi C}{\beta_{2} a_{1} Z_{01}} \cos \frac{\pi\left(z+\frac{a_{1}}{2}\right)}{a_{1}} e^{-j \beta_{2}(y-b)} \\
E_{z}=0 & (\text { IV.66 }) & H_{z}=-\frac{c}{Z_{01}} \sin \frac{\pi\left(z+\frac{a_{1}}{2}\right)}{a_{1}} e^{-j \beta_{2}(y-b)}
\end{array}
$$

$B$ and $C$ in the previous equations determine the magnitude and phase of the fields in the square waveguide and their values are essential to excite a circularly polarized wave. $B$ and $C$ can be expressed as two integral equations as follows:

$$
\begin{align*}
& B=-\frac{1}{P_{10_{2}}} \int_{v} \overrightarrow{E_{10}} \cdot \vec{J} d v+\frac{1}{P_{10_{2}}} \int_{v} \overrightarrow{H_{10}} \cdot \vec{M} d v  \tag{IV.68}\\
& C=-\frac{1}{P_{01_{2}}} \int_{v} \overrightarrow{E_{01}} \cdot \vec{J} d v+\frac{1}{P_{01_{2}}} \int_{v} \overrightarrow{H_{01}} \cdot \vec{M} d v \tag{IV.69}
\end{align*}
$$

Solving the integral equations leads to the following expressions of B and C:

$$
\begin{align*}
& \quad B=\frac{1}{P_{10_{2}}}\left[\vec{H}_{x}(s, b, 0) \cdot \overrightarrow{P_{e}}\right]=\frac{1}{P_{10_{2}}}\left[-\frac{1}{Z_{10_{1}}} \sin \left(\frac{\pi\left(\frac{a_{1}}{2}\right)}{a_{1}}\right)\left(-\alpha_{m}\right) A\left(-\frac{1}{Z_{10}}\right) \sin \frac{\pi s}{a}\right]= \\
& \frac{1}{P_{10_{2}}}\left[-\frac{\alpha_{m} A}{Z_{10} Z_{10_{1}}} \sin \frac{\pi s}{a}\right] \tag{IV.70}
\end{align*}
$$

$$
\begin{align*}
& \quad C=\frac{1}{P_{01}}\left[\vec{H}_{Z}(s, b, 0) \cdot \overrightarrow{P_{m}}\right]=\frac{1}{P_{01}}\left[\frac{1}{Z_{01}} \sin \left(\frac{\pi\left(\frac{a_{1}}{2}\right)}{a_{1}}\right)\left(-\alpha_{m}\right) A\left(\frac{j \pi}{\beta a Z_{10}}\right) \cos \frac{\pi s}{a}\right]= \\
& -\frac{j \pi \alpha_{m} A}{P_{01} Z_{10} Z_{01} \beta a} \cos \frac{\pi s}{a} \tag{IV.71}
\end{align*}
$$

In order to have a circularly polarized wave in the square waveguide, we need the following identities between B and C to hold true:

$$
|B|=|C| \text { and } \angle E_{Z_{10}}=\angle E_{x_{01}} \pm \frac{\pi}{2}
$$

if we have $a=a_{1}$ then:
$\beta=\beta_{1}=\beta_{2}$
$Z_{10_{1}}=Z_{01}$
$P_{10_{2}}=P_{01}$

Hence, solving the magnitudes equality leads to the following:

$$
\begin{align*}
& \frac{1}{P_{10_{2}}}\left[\frac{\alpha_{m} A}{Z_{10} Z_{10_{1}}} \sin \left(\frac{\pi s}{a}\right)\right]=\frac{\pi \alpha_{m} A}{P_{01} Z_{10} Z_{01} \beta a} \cos \left(\frac{\pi s}{a}\right) \\
& \sin \frac{\pi s}{a}=\frac{\pi}{\beta a} \cos \frac{\pi s}{a} \\
& \tan \frac{\pi s}{a}=\frac{\pi}{\beta a} \\
& \quad s=\frac{a}{2} \pm \frac{a}{\pi} \tan ^{-1} \sqrt{\left(\frac{2 f a}{c}\right)^{2}-1} \tag{IV.72}
\end{align*}
$$

Now, by examining $\angle E_{Z_{10}}=\angle E_{X_{01}} \pm \frac{\pi}{2}$ we notice that the $\angle E_{Z_{10}}$ is dependent on the $\angle B$ and the $\angle E_{x_{01}}$ is dependent on the $\angle C$. From the expressions of B and C it can be noticed that B is a real number and C is an imaginary number, which means that:

$$
\angle B=0 \text { and } \angle C=\frac{\pi}{2}
$$

This means that $\angle B-\angle C=-\frac{\pi}{2}$ hence the two electric fields in the square waveguide are in $90^{\circ}$ phase difference. The conditions above allow a slot in the broad-wall to radiate a circularly polarized wave into a square waveguide. It is worth mentioning that the condition
on the slot location is the same for both a square waveguide extraction and a circular waveguide extraction.

## Design

In this section, the conditions found in the theoretical analysis are put to test by designing a cross slot radiating into a square waveguide. The same z -arm shaped cross slot used in the circular waveguide case is used in here and the design parameters of the slot are shown in Table IV.1. A square waveguide with side width of $W_{s q}=2.12 \mathrm{~mm}$ is used to extract the power as shown in Figure IV. 17 (a). The cutoff frequency for both $T E_{01}$ and $T E_{10}$ modes is $f_{c 10}=70.7 \mathrm{GHz}$. For a design operating at 72 GHz , these dimensions for a square waveguide allow a very good match between the square waveguide and the cross slot extraction, at the same time it is ensured that only the fundamental modes are propagating. It is noticed again that when the width of the waveguide is increased, the efficiency of the extraction decreases.

The extraction with the dimensions stated above has a good input match with $S_{11}<$ $-19 d B$ for both ports in the frequency band $68-88 \mathrm{GHz}$. The polarizer has an isolation between ports 1 and 2 of $\left|S_{21}\right|>10 d B$ in the frequency band $71.825-72.66 \mathrm{GHz}$. This isolation reaches $\left|S_{21}\right|=18.2 d B$ at 72.2 GHz as shown in Figure IV.18. At the end of the square waveguide, a $T E_{10}$ and $T E_{01}$ modes are excited by the cross slot. The two modes have almost the same magnitude as shown in Figure IV. 19 and have a phase difference $88^{\circ}<\angle S_{31(2)}-\angle S_{31(1)}<97.5^{\circ}$ in the frequency range $71-77 \mathrm{GHz}$ as shown in Figure IV.20. As in the case of a circular waveguide extraction waveguide, the square waveguide extraction with a cross slot can reach efficiencies of more than $70 \%$. This design has a
peak efficiency of $83.2 \%$ at a frequency of 72.2 GHz when using aluminum as the conducting material. The efficiency of the design is more than $70 \%$ for the 71.76-72.85 GHz.

These properties of the polarizer make it a good device to be used to polarize the wave from linear to circular, both LHCP and RHCP, in the frequency range $71.76-72.85 \mathrm{GHz}$. In the next section, the performance of the polarizer is tested with a pyramidal horn.


Figure IV.18: The plots of the input reflection coefficient and the isolation between ports 1 and 2 for a square waveguide extraction design.


Figure IV.19: Plot of the power received at the end of the square waveguide in the form of $\mathrm{TE}_{01}$ and $\mathrm{TE}_{10}$ modes.


Figure IV.20: Plot of the phase difference between the $\mathrm{TE}_{01}$ and $\mathrm{TE}_{10}$ modes at the end of the square waveguide.

## V. A Single Slot Feeding a Pyramidal Horn Antenna

For this section, the polarizer designed in section IV is tested by feeding a pyramidal horn antenna.

## Horn Antenna Design

For the applications targeted by this dissertation, the pyramidal horn should have a high gain and yield circular polarization. For the circular polarization to be able to propagate, the pyramidal horn should have both identical E and H planes.

In order to achieve a maximum stable gain of around 23 dB , the iterative process described in [61] is used. The iterative process holds the following equality true:

$$
\begin{equation*}
\left(\sqrt{2 \chi}-\frac{b}{\lambda}\right)^{2}(2 \chi-1)=\left(\frac{G_{0}}{2 \pi} \sqrt{\frac{3}{2 \pi}} \frac{1}{\sqrt{\chi}}-\frac{a}{\lambda}\right)^{2}\left(\frac{G_{0}}{6 \pi^{3}} \frac{1}{\chi}-1\right) \tag{IV.73}
\end{equation*}
$$

Where $\chi=\rho_{e} / \lambda$ and $\frac{\rho_{h}}{\lambda}=\frac{G_{0}^{2}}{8 \pi^{3}} \frac{1}{\chi}$, and a, b, $\rho_{e}$ and $\rho_{h}$ are defined in Figure IV. 21 and $G_{0}$ is the dimensionless gain of the antenna.

Once the correct value of $\chi$ is found, all the other parameters of the design can be found. In this case, we have $=b, a_{1}=b_{1}$ and $\rho_{e}=\rho_{h}$. The dimensions of the pyramidal horn that lead to a gain of around 23 dB are stated in Table IV.3. For the horn to have a good match with the polarizer, a square waveguide with $a=b=2.12 \mathrm{~mm}$ is used to feed it.


Figure IV.21: Illustration of a pyramidal horn structure. [62]

Table IV.3: Pyramidal horn dimensions operating at E-band.

| Parameter | Description | Value |
| :---: | :---: | :---: |
| $\rho_{e}$ | flare edge length | 55.21 mm |
| $a_{1}$ | aperture length | 26.2 mm |
| $p_{e}$ | horn length without feeding waveguide | 49.3 mm |

## Full System Design

The horn is connected to the polarizer as shown in Figure IV.22. The polarizer used is the same one designed in section IV with cross slot dimensions designed to work at 72 GHz . Since the square waveguide of the horn and the cross slot extraction waveguide have the same dimensions, there is no need for a tapered waveguide transition. The full design has an operation at 71.8 GHz with an input reflection coefficient of $S_{11}<-17 \mathrm{~dB}$ over the frequency range $68-88 \mathrm{GHz}$ as shown in Figure IV.23. The isolation between ports 1
and 2 is $\left|S_{21}\right|>10 d B$ over the frequency range $71.7-72.1 \mathrm{GHz}$ reaching a maximum of 20 dB at 71.85 GHz . As in the case of the circular waveguide, the slot results in two resonances, in this case the resonance is at 73.2 GHz .


Figure IV.22: Layout of the pyramidal horn fed by a cross slot polarizer.


Figure IV.23: The plots of the input reflection coefficient and the isolation between ports 1 and 2 for a cross slot feeding a pyramidal horn.

The system produces a LHCP main beam if fed through port 1and a RHCP main beam if fed through port 2 . The system produces a flat gain between $21.3-21.55 \mathrm{~dB}$ over the frequency band of $71.8-73 \mathrm{GHz}$ with a system total efficiency $70 \%<\varepsilon_{h}<76 \%$. The peak gain the system attains is 21.55 dB at 71.9 GHz where the main resonant frequency of the slot is as shown in Figure IV.24. The axial ratio produced in the direction of the highest gain is $A R<2 d B$ over the frequency range $71.5-76 \mathrm{GHz}$, reaching values less than 1 dB over the frequency range $72-76 \mathrm{GHz}$ as shown in Figure IV. 26 .

The radiation pattern characteristics for both LHCP and RHCP are the same. The main beam produced has a $H P B W=11^{\circ}$ in all plane cuts over the frequency range $71.8-73$ GHz. In the plane $\phi=90^{\circ}$ the $S L L_{90}=-15 d B$ and is at its lowest value, whereas in the plane $\phi=0^{\circ}$ the $S L L_{0}=-23.5 d B$ and is at its highest value as shown in Figure IV.25. The cross polarization discrimination between both polarizations is of more than 20 dB over the frequency where the efficiency is higher than $70 \%$ as shown in Figure IV.27.

The cross slot polarizer is also tested with another pyramidal horn with gain of 17 dB . The horn has $a_{1}=10 \mathrm{~mm}$ and $p_{e}=15 \mathrm{~mm}$ and the same square waveguide is used to feed it. The simulation results show that the scattering parameters results are the same as in the case of the 23 dB gain horn as shown in Figure IV.23. In addition, the same results of axial ratio, efficiency and cross polarization discrimination are obtained as shown in Figure IV.24Figure IV. 26 and Figure IV.27. This proves that if the horn input is matched to the output of the polarizer, then the performance of the polarizer will not change. In this case, the easy solution for matching is to use the same square waveguide dimensions as the polarizer, or to use tapered square waveguides.


Figure IV.24: Plot of the Gain over frequency for the 17 dB and 23 dB designs.


Figure IV.25: Gain pattern in the planes $\phi=0^{\circ}$ and $\phi=90^{\circ}$ for both RHCP and LHCP patterns at 72 GHz .


Figure IV.26: Plot of the axial ratio for the 17 dB and 23 dB designs.


Figure IV.27: Cross polarization discrimination plot over frequency.

As a conclusion, this polarizer has an efficiency $70 \%<\varepsilon_{h}<76 \%$ over the frequency range $71.8-73 \mathrm{GHz}$ with a good axial ratio of less than 1.5 dB . The efficiency values can be improved by changing the square waveguide dimensions. In the next sections, both the square waveguide extraction and the circular waveguide extraction are improved in terms
of bandwidth. For that matter, the 17 dB horn is used to make the design and simulations faster.

## VI. Dual Slot Polarizer Using Serpentine Power

## COMBINER

The use of one slot limits the bandwidth of operation of the polarizer to $1.5 \%$ of the frequency band. In the case of a good isolation of more than 20 dB of the polarizer, for a good transmitting/receiving response, the bandwidth is around $0.4 \%$ of the frequency band. In order to increase the bandwidth of the polarizer, multiple slots are created in the broad-wall of the rectangular waveguide, their power is extracted through either a circular waveguide or a square waveguide, then combined and fed to the horn. By doing so, the second slot helps the first slot to extract more power, hence making the bandwidth of the polarizer wider. In this section, the use of 2 slots with a serpentine circular waveguide combiner is explored at 84 GHz .

## - Slots Design

The first step in the optimization of the polarizer is the design of the cross slots. The cross slots determine the frequency of operation of the polarizer and the amount of power that could be extracted from the rectangular waveguide. As a starting point, two slots with the dimensions stated in Table IV. 2 are used for a design operating at 84 GHz . In order to minimize the effect of coupling between the slots on the operating frequency, the distance between the two slots is increased and set at $L_{\text {sep }}=6 \mathrm{~mm}$. A WR-10 waveguide ( $a=$ 2.54 mm and $b=1.27 \mathrm{~mm}$ ) is used for this design. The extraction waveguides for each
slot have a cross section diameter of $D_{\text {ext }}=2.5 \mathrm{~mm}$ and a length minimized to $L_{\text {ext }}=$ 5 mm . When adding the extraction waveguides a mismatch between the slots and the waveguides takes place. This mismatch can be solved by changing the dimensions of the slots in order to extract the maximum power from the rectangular waveguide. In addition, when the extraction waveguides are connected using the serpentine power combiner, more loading is added on the slots, which changes their operating frequency and efficiency. The final slots dimensions are shown in Table IV.4.

Another configuration of the slots, that could increase the bandwidth of operation, is by having the resonant frequencies of the slots close to each other, but not the same. This configuration is studied in the section of square waveguides combiners.

Table IV.4: Dimensions of the Z-shaped arm cross-slots at 84 GHz used in the serpentine power combiner.

| Parameter | Description | Value |
| :---: | :---: | :---: |
| $L_{z 1}$ | long arm length | 1.6 mm |
| $L_{z 2}$ | short arm length | 0.45 mm |
| $W_{a}$ | arm width | 0.15 mm |
| $S$ | slot center position | 0.7 mm |
| $\alpha_{2}$ | arm angle with y-axis | $43^{\circ}$ |
| $t_{w 3}$ | slot thickness | 0.4 mm |

## Serpentine Power Combiner Design

In order to combine the power extracted by the two cross slots, a serpentine power combiner is used. The serpentine power combiner has two circular waveguide inputs and one circular waveguide output. Each section of the input waveguides has a convex curvature followed by a concave curvature or vice versa as shown in Figure IV. 28 (a). These curvatures are arcs of a circle connected smoothly. These two sections intersect at the output waveguide. The radius of all waveguides used in the serpentine design is optimized to be $R_{w s}=1.25 \mathrm{~mm}$ with a cutoff frequency $f_{c}=70.33 \mathrm{GHz}$ for the fundamental mode $T E_{11}$ and a cutoff frequency for $T M_{01}, f_{c}=91.9 \mathrm{GHz}$, ensuring that only the fundamental mode is traveling inside the waveguide at 84 GHz .

The curvatures of the serpentine waveguides have a radius $R_{\text {serp }}=45 \mathrm{~mm}$ as shown in Figure IV. 28 (b). The curvature radius ensures that at the intersection plane of the two sections of the serpentine, the geometry is smooth with respect to the operating wavelength. The smoothness ensures that the wave is not scattered when it hits the intersection plane and the phase of each $T E_{11}$ mode is preserved. This ensures that the $90^{\circ}$ phase difference is intact and the orthogonal components of the electric field have close magnitudes.

The separation between the slots $L_{s s}$ along with $R_{\text {serp }}$, shown in Figure IV. 28 (b), are chosen in a way that makes each section of the serpentine input waveguides well matched with the serpentine output waveguide. If this matching is not done correctly, a standing wave is seen in the section of the waveguide that extracts the least power, causing the efficiency of the combiner to drop. In this design, the separation between the slots is optimized to be $L_{s s}=10 \mathrm{~mm}$ for $R_{\text {serp }}=45 \mathrm{~mm}$. All these factors are essential in
making sure that the wave extracted from the rectangular waveguide is circularly polarized and that the extraction process has minimal losses.

The performance of the serpentine is discussed in what follows and a comparison with the single slot polarizer is performed.

(b)

Figure IV.28: (a) Illustration of the serpentine combiner geometry, (b) Illustration of the different parts of the system combined.

## - Full System Design and Performance

The serpentine power combiner is used to feed the conical horn from section III producing a gain of 23 dB Figure IV. 28 (b). A tapered circular waveguide of length $L_{\text {trans }}=8 \mathrm{~mm}$ is used to match the polarizer to the conical horn.

By examining the E-fields inside the system it can be noticed that the slot that is closer to the feeding port extracts the most power from the rectangular waveguide as seen in Figure IV.29. The intensity of the E-field decreases after passing slot 1 and it decreases
again after passing slot 2 . In the circular waveguides that are used for extracting the power and feeding it to the serpentine, two orthogonal $T E_{11}$ modes can be seen. These modes are combined at the end of the serpentine and the intensity of the E-field is shown to have increased.


Figure IV.29: The E-field inside the different parts of the system at 84 GHz .


Figure IV.30: The s-parameters of the system with a serpentine power combiner.

The s-parameters of the system, shown in Figure IV.30, show that the system is well matched for the frequency range $80-86 \mathrm{GHz}$. The system shows two different resonances
at 82 and 83.5 GHz caused by the coupling between the slots. Those resonances create two different frequency bands where the isolation between ports 1 and 2 is $\left|S_{21}\right|>10 d B$. These frequency bands are $B W_{1}=[81.882 .35] G \mathrm{GHz}$ and $B W_{2}=[82.983 .9] \mathrm{GHz} . B W_{2}$ has a range where the isolation reaches values higher than 20 dB over a frequency band of 250 MHz .


Figure IV.31: (a) Plot of the comparison of the maximum gain performance; (b) plot of the comparison between the axial ratios of the serpentine vs the single slot polarizer.

The system produces a LHCP main beam if fed by port 1 and a RHCP main beam if fed by port 2. The maximum gain generated by the system, shown in Figure IV. 31 (a), is improved for the frequency range of $80-83 \mathrm{GHz}$ in comparison to the single slot polarizer and it reaches values between 20 and 21 dB . In the frequency range of $83-86 \mathrm{GHz}$, the gain has decreased in comparison to the single slot polarizer by 1 dB . This decrease is due to the conduction losses in the serpentine combiner due to its large dimensions. The system has an efficiency $50 \%<\varepsilon_{s}<70 \%$ over the frequency band $80-86 \mathrm{GHz}$.

The axial ratio of the wave radiated by the serpentine fed horn has values higher than the one fed by the single slot polarizer. This deterioration in the axial ratio is due to the serpentine structure that causes a change in the phase difference between the E-fields of the different $T E_{11}$ modes. However, the axial ratio, in the direction of the maximum gain, is less than 3 dB over the frequency range of $80-85.2 \mathrm{GHz}$ as shown in Figure IV. 31 (b). The axial ratio reaches a minimum value of 0.65 dB at 81 GHz .

The serpentine polarizer does not affect the radiation pattern of the conical horn. The HPBW of the horn is still $12^{\circ}$ in all plane cuts as shown in Figure IV.32. The SLL is improved and reaches a value of -27.9 dB for all plane cuts at 84 GHz . The XPD is around 22 dB at 84 GHz and the entire HPBW is circularly polarized.

In conclusion, the large size of the serpentine combiner drops its efficiency due to conduction losses. The polarizer still improves the gain performance of the system at frequencies close to the resonant frequencies of the slots.


Figure IV.32: The LHCP and RHCP radiation pattern of the serpentine fed horn in different plane cuts at 84 GHz .

## VII. Dual Slot Polarizer Using A Square Waveguide POWER COMBINER

In this section, an improvement to the performance of a single slot polarizer with a square waveguide extraction is investigated at 72 GHz . The improvement relies on the use of two slots in the rectangular waveguide with their powers combined using square waveguide combiners. The design and improvements are discussed in what follows.

## $\square$ Dual Slot Design

Same as in the case of the serpentine combiner, the first step in designing this system is the design of the cross slots. By determining the right configuration of the slots, the bandwidth and efficiency of the system can be improved. Two different scenarios are considered for this matter and explained below:

Scenario 1. The slots have the same resonant frequency:
For this scenario, the slots used have the same exact dimensions which lead to the same resonant frequency. The aim of this design is to increase the isolation between ports 1 and 2 to very high values and to have the best axial ratio values between the different components of the electric field. The slots used in this design have their dimensions stated in Table IV.5.

## Scenario 2. The slots have close resonance frequencies

In this scenario, the slots used have dimensions that are different but very close. This configuration of the slots results in a wider bandwidth of operation of the polarizer and in higher efficiencies of the system over a wider frequency range than in scenario 1 . The dimensions of the slots are state in Table IV.5.

The slots are located on the broad-wall of a WR-12 waveguide and have a separation distance of $L_{s s}=6 \mathrm{~mm}$ for both scenarios. For both scenarios, the frequency of operation is close to 72 GHz . The coupling between the slots along with the dimensions of the slots contributed to the frequency shift between the different designs.

Table IV.5: Dimensions of the Z-shaped arm cross-slots at 72 GHz used in the square waveguide power combiner.

| Parameter | Description | Value |  |  |
| :---: | :---: | :---: | :---: | :---: |
|  |  | Scenario 1 | Scenario 2 |  |
|  |  |  | Slot 1 | Slot 2 |
| $L_{z 1}$ | long arm length | 1.85 mm | 1.85 mm | 1.80 mm |
| $L_{z 2}$ | short arm length | 0.46 mm | 0.46 mm | 0.46 mm |
| $W_{a}$ | arm width | 0.15 mm | 0.15 mm | 0.15 mm |
| $S$ | slot center position | 0.69 mm | 0.69 mm | 0.69 mm |
| $\alpha_{2}$ | arm angle with $y$-axis | $46^{\circ}$ | $46^{\circ}$ | $46^{\circ}$ |
| $t_{w 3}$ | slot thickness | 0.4 mm | 0.4 mm | 0.4 mm |

## - Full System Design

The square waveguide combiner, shown in Figure IV. 33 (a), has two square waveguide inputs and one square waveguide output. The square waveguide inputs have a cross section width of $a_{\sin 1}=a_{\sin 2}=2.4 \mathrm{~mm}$ while the square waveguide output has a cross section width $a_{\text {sout }}=2.7 \mathrm{~mm}$. The input waveguide cross section width plays a role in the efficiency of the extraction of power from the rectangular waveguide. On the other hand, the output waveguide cross section width plays a role in the matching of the power combiner. The output waveguide is located at the same distance from the two input waveguides. This position ensures that the power combined from each input waveguide does not bounce upon reaching the output waveguide. It also ensures that the power from both waveguides does not add up off the centerline of the output waveguide, causing beam squinting in the radiation pattern.

The power combiner has its input and output waveguides connected by tapered tilted waveguides. The taper in the connecting waveguides plays an essential role in better matching the entire design, since the waveguides it is connecting have different cross section widths. The taper causes minimal reflection to occur at the connection with the output waveguide of the power combiner. The tilting angle $\alpha_{t}$ of the waveguides is optimized in a way that minimizes the scattering at the intersection plane of the waveguides. This scattering might reflect back into the input waveguides and cause a standing wave. This phenomenon reduces the efficiency of the power combiner. For a power combiner working at 72 GHz , the tilting angle is $\alpha_{t}=16.7^{\circ}$. The separation between the output of the input waveguides and the input of the output waveguides is $L_{c}=$ 10 mm .


Figure IV.33: (a) The layout and the different parameters of the square waveguide combiner, (b) The E-field inside the different sections of the full system at 72 GHz when a pyramidal horn is connected.

## - Results

The power combiner feeds a pyramidal horn as shown in Figure IV. 33 (b). The pyramidal horn used here is the same as the one in section V producing a gain of 17 dB . As Figure IV. 33 (b) shows, when the wave fed passes the first slot, its power in the rectangular waveguide is reduced. When it passes the second slot, its power becomes insignificant in the rectangular waveguide. The power transfer into the square waveguides done by the slots causes this power reduction. In the square waveguides, the E-fields of the two $T E_{10}$ modes can be seen. These E-fields are in two orthogonal directions all over the entire combiner. It can be noticed that there is a phase difference of around $\pm 90^{\circ}$ between them. This phase difference is carried out through the entire square power combiner and can be seen in the horn. At the intersection plane between the titled waveguides, it can be noticed that the scattering is minimal and the power is being combined successfully resulting in higher intensities of both E-fields in the power combiner's output waveguide. This is true for both scenarios and for both feeding ports.

For scenario 1, the design resonant frequency is at 72.4 GHz as shown in Figure IV.34. The design is well matched for the frequency range $70.5-76 \mathrm{GHz}$. The different loading that is imposed on the slots, in comparison to the single slot design, causes the shift in the operating frequency. In addition, the coupling between the two slots contributes to this frequency shift. For a better matching that can improve the efficiency of the slots, a rectangular ridge is added in the rectangular waveguide, on the opposite side of each slot as shown in Figure IV.35. The ridge has a width $W_{R}=0.2 \mathrm{~mm}$ and a ridge length $L_{R}=$ 2 mm . The ridge improves the resonant frequency and the 10 dB and 20 dB isolation
between ports 1 and 2 as shown in Figure IV.34, along with the efficiency of the system. The ridge does not affect the input matching bandwidth.


Figure IV.34: S-parameters comparison between a regular combiner and a combiner with a ridge.


Without Ridge

Figure IV.35: Illustration of the design with and without a ridge in the rectangular waveguide.

The design of scenario 1 has an isolation of $\left|S_{21}\right|>10 \mathrm{~dB}$ in the frequency range 71.7 - 73.1 GHz with or without a ridge installed in the rectangular waveguide. The $\left|S_{21}\right|>$ 20 dB for a bandwidth of 400 MHz in the case of a ridge, and 300 MHZ when the ridge is not added. The same s-parameter results are obtained if port 2 is fed instead of port 1.

Simulation indicates that scenario 1 can produce a maximum gain of 17.1 dB at 72.4 GHz with or without the ridge. At this frequency, the system has a total efficiency of $\varepsilon_{c, \max }=89 \%$. The system has a total efficiency of more than $\varepsilon_{c}>70 \%$ over the frequency range $71.5-73.1 \mathrm{GHz}$ as shown in Figure IV. 36 (a). The improvement of the gain that the two slot combiner introduced over the single slot fed horn is also shown in Figure IV. 36 (a). The gain improvement is noticed over the frequency band of $70-76 \mathrm{GHz}$. The axial ratio in the direction of the highest gain is shown in Figure IV. 36 (b). For both combiner designs, with or without a ridge, the axial ratio in the direction of the highest gain is less than 3 dB over the bandwidth over which the system has an efficiency of more than $70 \%$. The introduction of the ridge improves the axial ratio between $72-73 \mathrm{GHz}$. The radiation pattern of the horn is not affected by the introduction of the combiner as shown in Figure IV.37. The HPBW is circularly polarized in all plane cuts over the frequency range $71.7-73 \mathrm{GHz}$. The system exhibits a LHCP main beam when fed through port 1 and RHCP main beam when fed through port 2 .

(b)

Figure IV.36: (a) Gain vs frequency for different designs, (b) Axial ratio vs frequency for the rectangular waveguide with and without ridge.

For scenario 2, the same ridge is also used since it improves the performance of the system. The resonant frequency of scenario 2 is at $\mathrm{f}=72.65 \mathrm{GHz}$. The shift in resonant frequency in comparison to scenario 1 is caused by the slot 2 size, which is smaller than the size of slot 1 . The system is well matched in the frequency bandwidths of $71.4-73.4$ GHz and $74-76 \mathrm{GHz}$ as shown in Figure IV.38. Scenario 2 shows two different frequency ranges where the isolation between ports 1 and 2 is good, showing a good slot power extraction in these frequency bands which promises a good polarizer efficiency in these ranges. The isolation between ports 1 and 2 is $\left|S_{21}\right|>10 d B$ for the frequency range 72.2-
73.2 GHz. The isolation $7 \mathrm{~dB}<\left|S_{21}\right|<9.5 d B$ in the frequency range $74.2-75.2 \mathrm{GHz}$. In these two bands the system works well as a polarizer. This is shown in the analysis of the radiation properties of the horn in scenario 2.


Figure IV.37: Gain pattern of scenarios 1 and 2 for both LHCP and RHCP when port 1 is fed at 72.5 GHz .


Figure IV.38: S-parameters of the different scenarios vs frequency.

The maximum gain achieved in scenario 2 is 17.1 dB at 71.65 GHz . The system total efficiency is $\varepsilon_{c, \max }=84 \%$ at this frequency. This scenario exhibits two frequency ranges, $B W_{1}=[71.973 .1] \mathrm{GHz}$ and $B W_{2}=[74.375 .2] \mathrm{GHz}$ where the system total efficiency is $\varepsilon_{c}>70 \%$ as shown in Figure IV. 39 (a). The axial ratio over these two frequency bandwidths is less than 3 dB as shown in Figure IV. 39 (b). It can be noticed that over $B W_{1}$, the axial ratio of scenario 2 is higher than scenario 1 . This deterioration in axial ratio is caused by slot 2 since it has a resonance at a higher frequency (Section IV shows that the axial ratio is higher before the resonant frequency of the slot). In $B W_{2}$, both the gain and the axial ratio of scenario 2 are improved in comparison to scenario 1 . The radiation pattern of the horn is not affected by scenario 2 as shown in Figure IV.37. The SLL and HPBW are the same as in the case of a linear gain horn. The HPBW of the system is entirely circularly polarized, exhibiting a LHCP when fed through port 1 and a RHCP when fed through port 2 and an axial ratio $A R>3 d B$ for all the angles and the plane cuts.

As a conclusion, scenario 1 improves the bandwidth over which the system has a good efficiency and a good isolation between both ports. The system works as an efficient polarizer for a bandwidth of $2.2 \%$ of the frequency band. Scenario 2 creates two separated, but close frequency bands where the system works well as a polarizer with a bandwidth of 1.4 and $1.6 \%$ of the frequency band.

The property illustrated in scenario 2 is further investigated in the following section in order to create a system that can work at frequency bands that are far apart from each other and where no single OMT polarizer can cover them.

(b)

Figure IV.39:(a) Gain comparison vs frequency, (b) Axial ratio comparison vs frequency for scenarios 1 and 2.

## VIII. Multi-Band Slot Polarizer Design

- Concept

In both, the serpentine power combiner and the square waveguide power combiner, the slots were capable of creating two different frequency bands where the polarizer can operate with high efficiency and provide a good axial ratio and good isolation between the two ports. This means that this technique can be used to create a system that can operate at
two frequency bands that are far away from each other and still provide the desired antenna properties. The large frequency gap between the slots that are used, removes the effect of coupling on the slots. Since the emphasis of this dissertation is on the wave propagation at 72 and 84 GHz , a system that operates at these two frequencies is designed using several scenarios for the design.

## Dual Horn Design

The first way of realizing such system is by using a slot that works at 72 GHz and another one that works at 84 GHz . The slots will have minimum coupling between them since they have resonances at frequencies that are far away from each other. The slot dimensions are indicated in Table IV.6. The slots feed two circular waveguides of radii $R_{72}=2.5 \mathrm{~mm}$ and $R_{84}=2.2 \mathrm{~mm}$, that in turn feed two different horns as shown in Figure IV.40. Tapered waveguides are used to connect the circular waveguides to the input of each horn. The tapered waveguide of the 72 GHz slot has a length of $L_{\text {trans } 72}=3 \mathrm{~mm}$ while the one for the 84 GHz slot has a length of $L_{\text {trans84 }}=8 \mathrm{~mm}$. The 72 GHz slot is on the broad-wall of a WR-12 waveguide while the 84 GHz slot is on the broad-wall of a WR10 waveguide. The two rectangular waveguides are connected by tapered waveguides and the feeding of the whole system is made of WR-12 rectangular waveguides.

The system has two resonances at 72 GHz and 84.3 GHz as shown in Figure IV.41. The system is well matched over the frequency range $71-90 \mathrm{GHz}$. Similar to the single slot design, the 72 GHz slot has two resonances. The 84 GHz slot is slightly different in this case and it operates at 84.3 GHz since its dimensions is smaller than the single slot design.

Table IV.6: Dimensions of the Z-shaped arm cross-slots at 72 and 84 GHz used in the multi-band system.

| Parameter | Description |  | Value |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | Dual Horn System |  | Single Horn System |  |
|  |  | 72 GHz | 84 GHz | 72 GHz | 84 GHz |
| $L_{z 1}$ | long arm length | 1.85 mm | 1.62 mm | 1.85 mm | 1.70 mm |
| $L_{z 2}$ | short arm length | 0.46 mm | 0.36 mm | 0.46 mm | 0.38 mm |
| $W_{a}$ | arm width | 0.15 mm | 0.15 mm | 0.15 mm | 0.15 mm |
| $S$ | slot center position | 0.69 mm | 0.69 mm | 0.69 mm | 0.69 mm |
| $\alpha_{2}$ | arm angle with y-axis | $43.5^{\circ}$ | $44^{\circ}$ | $46^{\circ}$ | $46^{\circ}$ |
| $t_{w 3}$ | slot thickness | 0.4 mm | 0.4 mm | 0.4 mm | 0.4 mm |



Figure IV.40: The full system with two horns.


Figure IV.41: S-parameters of the system with two horns operating at 72 and 84 GHz .

The gain performance of the system shows that the two-slot design delivers similar results to the case of a single slot design. The system has two frequency bands where it has a high efficiency. The first one is between $71.7-73 \mathrm{GHz}$ and the second one is between 84 and 86 GHz as shown in Figure IV.42. The efficiency of the design at 84 GHz is higher because in this case, the design acts as an array of two horns fed with different power levels. This occurs because the 72 GHz slot still radiates a wave with low power into its horn that is circularly polarized. This increases the gain of the horn but distorts the radiation pattern of the system only in the 84 GHz band as shown in Figure IV.43. In this band, the SLL increases in the plane cuts perpendicular to the feeding rectangular waveguide and the HPBW becomes narrower due to the horns forming a non uniform array.

The axial ratio of the system is less than 1.5 dB in the frequency bands where the antenna operates with high efficiency. This axial ratio ensures that the HPBW is entirely circularly polarized over the frequencies of operation of the system.


Figure IV.42: Gain comparison between linear gain of the horn and the system with two horns.


Figure IV.43: The co-polarization and cross-polarization radiation patterns of the system at 85 GHz in different plane cuts.


Figure IV.44: The axial ratio of the system with two horns.

The results of the system show that it operates well at 72 GHz but has a distorted radiation pattern at 84 GHz . This could be solved by using a single horn fed by the two slots. This solution solves the problem at 84 GHz , makes the design more concise, and less expensive by cutting down on the number of horns being used. This solution is illustrated in the following section.

## - Single Horn Full System Design

Instead of using two horns fed by two different channels, in this section the design of a single horn fed by two slots operating at 72 and 84 GHz is illustrated. Each slot radiates its power into a different channel. Each channel is a square waveguide with its width $W_{s 72}=2.17 \mathrm{~mm}$ for the 72 GHz channel and $W_{s 84}=1.9 \mathrm{~mm}$ for the 84 GHz channel.

The two channels are connected to the feed of the horn by tapered square waveguides as shown in Figure IV.45. The slot dimensions used in this design are indicated in Table IV.6.


Figure IV.45: The E-field in the system with two square waveguide channels for the different frequencies of operation.

The same combiner used in the case of the square waveguide combiner is used here with the width of the input square waveguides changed. Since the channel for the 84 GHz has a cutoff frequency higher than 72 GHz , the only channel where the wave propagates in at 72 GHz is the 72 GHz channel as seen in Figure IV.45. At 84 GHz , the 72 GHz slot couples low power into the 72 GHz channel, causing a wave with low power to propagate at 84 GHz in the 72 GHz channel. The power gets combined in the output waveguide of the combiner. The horn used here is the pyramidal horn with gain 17 dB used in previous sections.

The designed system has two resonant frequencies at 71.65 GHz and 84.15 GHz where there is a good input match as shown in Figure IV.46. The system is well matched between
$71-73 \mathrm{GHz}$ and $81-86 \mathrm{GHz}$. The isolation between ports 1 and 2 is $\left|S_{21}\right|>10 \mathrm{~dB}$ in the frequency bands $71.4-71.8 \mathrm{GHz}$ and $83.7-84.3 \mathrm{GHz}$. This isolation reaches values more than 20 dB at the frequencies of resonance with a bandwidth of $100-150 \mathrm{MHz}$ at each resonance.


Figure IV.46: The single horn full system s-parameters.


Figure IV.47: The gain and axial ratio of the system with a single horn.

The gain achieved by the system reaches its maximum at 71.65 GHz with a gain of 16.45 dB and another maximum at 84.15 GHz with a gain of 18.3 dB as shown in Figure IV.47. The higher efficiency at 84.15 GHz is due to the fact that the 72 GHz channel passes power at 84 GHz . The efficiency at 71.65 GHz is limited by the geometry of the system. It can be noticed that another resonance at 73.5 GHz is seen causing a bad input matching for the entire system. This resonance is not caused by the slot; instead the channels of the combiner cause it. The efficiency around this frequency drops because of the high reflections at the input. The frequency bands where the efficiency of the system is higher than $70 \%$ are between $71.4-72 \mathrm{GHz}$ and $83.6-84.8 \mathrm{GHz}$.


Figure IV.48: Gain pattern of the single horn system in different plane cuts.

The axial ratio in the direction of the highest gain is less than 3 dB in the frequency bands of $71.45-71.9 \mathrm{GHz}$ and of $84.85-84.6 \mathrm{GHz}$ as shown in Figure IV.47. At 71.65 GHz and 84.4 GHz , the axial ratio reaches a minimum with a value of 0.9 dB and 0.3 dB respectively. The system does not affect the radiation pattern of the horn. The HPBW is still the same at with HPBW $=24^{\circ}$. The SLL of the horn is at -27 dB in all plane cuts. The HPBW of the horn is entirely circularly polarized at 71.7 GHz and 84.15 GHz in all plane cuts as shown in Figure IV.48.

These results show that the system works at 71.7 and 84.15 GHz with a good bandwidth of operation of 400 MHz and 900 MHz , respectively. The system can be used to feed any horn and make it circularly polarized in those frequency bands.

## IX. CONCLUSION

This chapter illustrated the design of the cross slot fed waveguide polarizer. Both a circular waveguide and square waveguide were investigated as an extracting waveguide. The power of multiple cross slots was combined to form a more efficient polarizer that has a $70 \%$ efficiency bandwidth of $2.2 \%$ of the frequency band. This polarizer concept is also used to design a multi band circularly polarized antenna system capable of providing either LHCP or a RHCP by changing the feeding port.

## CONCLUSION \& FUTURE WORK

The design and fabrication of several antennas and microwave devices that are easy to fabricate, less expensive than the ones available in the market today are presented in this dissertation. These designs can be efficiently used for satellite communication, terrestrial communication and radar systems operating in the W/V bands

The slotted rectangular waveguide array presented in this dissertation was designed and fabricated to operate at 85 GHz . It consists of an eight rectangular waveguide elements array with each having sixteen cross slots. The array was shown to produce two beams with one RHCP when fed by port 1 and the other LHCP when fed by port 2 . The gain achieved by the array is around 25 dB in its operating frequency range and its cross polarization discrimination is more than 20 dB . The array has an improved isolation between its ports which in turn improves its efficiency by a factor of $8 \%$ over the design reported in [35]. The z-arm shaped cross-slot introduced in the array can fit on the broadwall of the rectangular waveguide for any frequency of operation and has an efficiency that is equal to the conventional cross slot.

The cross-slotted waveguide polarizer was designed for both a square waveguide extraction and a circular waveguide extraction and can produce either an LHCP if fed by port 1 or and RHCP if fed by port 2 . The polarizer can produce axial ratio values of less than 1 dB over its operating frequency band and a stable phase difference between the orthogonal generated modes of around $90^{\circ}$. The polarizer has a $1.2-1.5 \%$ bandwidth of operation with an efficiency higher than $70 \%$. The bandwidth is improved to $2.25 \%$ by the
use of square waveguide combiners that combine the power of different slots. The polarizer has the ability to operate simultaneously at multiple frequency bands that are far away from each other while producing a good axial ratio and an isolation between the feeding ports that is more than 20 dB making it goof for radar applications.

A system that allows the activation of the different feeding ports of the cross slotted rectangular waveguide array should be realized in the future. The fabrication of the square waveguide combiner that is used to improve the cross slot polarizer should be fabricated and tested. Finally, the design of the square waveguide combiner that operates at both 72 GHz and 84 GHz and the 23 dB pyramidal horn should be fabricated and tested.

## Appendix $A$.

## LIQUID CRYSTAL RECONFIGURABLE ANTENNA DESIGNS

This appendix illustrates the work done on liquid crystal reconfigurable antennas [63],[64]. A design operating at $76-77 \mathrm{GHz}$ is shown and the difficulties faced are detailed. This work is the initial work done at W/V band frequencies and it shows why a transition into slotted waveguide arrays and devices is chosen over liquid crystal devices for automotive radars.

## - Liquid Crystal Properties at Microwave Frequencies

The LC employed in this proposed antenna is in its nematic phase. That is, at microwave frequencies, the LC changes dielectric properties due to the different orientations of its molecules caused by an electrostatic field [64], as shown in Figure A.1.

The liquid crystal used for this antenna is the BL038 type obtained from Merck Laboratories. The BL038 liquid crystal has a dielectric constant of $\varepsilon_{r, \perp}=2.36$ with a $\tan \delta$ $=0.006$, and a parallel dielectric constant value of $\varepsilon_{r, \|}=3.03$ with a $\tan \delta=0.01$. By changing the DC voltage at the plates, different $\varepsilon_{\mathrm{r}}$ values are obtained.


Figure A.1: Illustration of liquid crystal transition as a bias voltage is applied.

## - Frequency Tunable Array Designed at X-Band

An antenna array is designed at X-band as a first step for the ease of fabrication and better understanding of how liquid crystal could be used in reconfigurable arrays. This antenna design is a step before moving to W/V band frequencies in order to understand the effect of switching on the antenna response and whether it could be replaced by the control of the voltage value applied to all the patches.

## Design

The proposed antenna array consists of sixteen patches in a $4 \times 4$ layout, fed by a single port. Liquid crystal serves as the substrate situated beneath each individual patch, and the structural substrate elsewhere is Rogers 3003 with $\varepsilon_{\mathrm{r}}=3$, thickness $=0.5 \mathrm{~mm}$ and $\tan \delta=$ 0.001 at 10 GHz . Each rectangular patch element is fed using an inset feed to match its input impedance to the impedance of the feeding line. A corporate feeding network is used to feed all the patches in phase. All patches in the same row are fed from either the top or the bottom, which creates a phase difference between patches in consecutive rows. A narrow gap is inserted in the quarter-wave transformer leading to the input of the patch as shown in Figure A.2. The gap in the feeding line blocks the DC supplied to this particular patch from flowing to other patches, hence securing the element by element DC control. The position of the gap is optimized to minimize the effect on the bandwidth of the array and the resonant frequency. The DC voltage reaches each patch by the means of a pin that goes through the LC cavity and the ground plane to connect the patch to the DC feeding network located underneath the array ground plane. The DC feeding network consists of sixteen feeding connections to the patches that are connected to the main DC line via
mechanical switches. An RF choke is placed between each pin and feeding connection to isolate the RF from the DC network as shown in Figure A.3.


Figure A.2: Illustration of the corporate feed and array layout.


Figure A.3: (a) $4 \times 4$ array of rectangular patches on top with the DC feeding network on the bottom, (b) The DC feeding network.

For an operation at 9.38 GHz , the patch has a length $\mathrm{L}=6.8 \mathrm{~mm}$, a width $\mathrm{W}=10.6 \mathrm{~mm}$, and an inset of 0.4 mm with a gap of 0.4 mm . The DC blocking gap is optimized to be 0.11 mm , the connecting pin has a diameter of 0.6 mm and a length of 2 cm extending underneath the ground. The position of the pin is optimized to minimize its effect on the array operation. The patches are aligned with the LC cavities underneath them.

## - Simulated Results

The design was simulated using CST MWS. When the DC feed is 0 V for all patches, the array resonates at 9.38 GHz with a bandwidth of 260 MHz the LC relative permittivity is 2.65 in this case. For different patch combinations fed with a DC voltage of 40 V , with an LC dielectric constant $\varepsilon_{r}=3$, the resonant frequency shifts to a lower value. As shown in Figure A.4, the resonant frequency shifts from 9.38 GHz to 9 GHz as the lowest value. In the case of two resonant frequencies, the main resonance has a value greater than 9 GHz while the second resonance has a value less than 9 GHz . The bandwidth covered by this frequency tuning is 700 MHz , extending between 8.8 GHz and 9.5 GHz .

The radiation pattern of the array showed 2 main beams in the direction of $54^{\circ}$ in the E plane. In most cases, there is no beam steering. A $4^{\circ}$ shift in the two main lobes is seen as shown in Figure A.5. The element combinations for this shift are: Case 1 - first row and first column. Case 2: first and last square blocks of 4 patches. Case 3: one of the diagonals.


Figure A.4: Frequency tuning caused by different combinations of patches with $\mathrm{DC}=40 \mathrm{~V}$.


Figure A.5: Gain pattern (E-plane) showing a small shift in the direction of the main beams.

## - Measured Results

The antenna is fabricated and tested as shown in Figure A. 6 and Figure A.7. The filling tubes are used to pump the LC into the array to fill the cavities. The measured results show that the array has a resonant frequency when no DC is applied at 9.35 GHz . The shift between simulated and measured results is attributed to the fabrication errors, such as the etching of the substrate and the glue added in order to seal the whole antenna layers and hold them together. When the DC voltage is changed between 0 and 40 V , the array resonant frequency is also tuned and the resonant frequency shifts to lower frequencies until it reaches 9.15 GHz at 40 VDC as seen in Figure A.8. When different switching configurations are tested, a similar behavior of the array resonant frequency is seen (Figure A.9). This work shows that the DC voltage control and the individual patch control can yield to similar results. The advantages of the voltage control is that it makes the array design much easier and simpler: instead of controlling switches that feed the voltage to the individual patches, and adds a level of complication to the design since the switches are integrated in it, it is easier to control the DC voltage source that feeds all the patches.


Figure A.6: Different Antenna Fabrication Layers.


Figure A.7: Assembled antenna.


Figure A.8: Antenna Return Loss caused by different applied DC voltages when all switches are ON.


Figure A.9: Different antenna Return Loss plots for different switch connections at 40 V DC.

## - Frequency tunable and Beam Steering Array Design at W/V BAND.

The next step in this work is to design an array operating at at $76-77 \mathrm{GHz}$. The following reconfigurable antenna array is realized [63].

## Design

The proposed antenna array is composed of a total of three layers. The top layer consists of eight rectangular patches in a 2 x 4 layout, fed by a single microstrip feed network. It is important to indicate that each individual patch is fed by a uniform meander line of impedance $Z_{0}=100 \Omega$. Since the antenna operates in the W/V-band, a grounded CPW to microstrip line transition (Figure A.10) or a rectangular waveguide to microstrip line transition (Figure A.11) are used. The middle layer is composed of a Rogers 3003 substrate with $\varepsilon_{\mathrm{r}}=3$, a thickness of 0.13 mm , and a $\tan \delta$ of 0.001 at 70 GHz . The thickness of the substrate is taken to be thin to ensure high speed switching of the incorporated LC in the cavities. Within this layer, LC cavities are positioned underneath each individual patch. Under each meander line feeding the individual patches there is an LC cavity. The meander lines are used to control the phase at the input of each patch.

All patches in the same row are fed from either the top or the bottom, which creates a phase difference of $180^{\circ}$ between patches in different rows. The bottom layer of the array is composed of the ground plane where a gap is inserted underneath each meandered line. The gap insures the separation of DC voltage supplied to each row of meander lines.

By introducing the gaps in the ground plane, the radiation pattern is affected. To impede such a consequence, a Rogers Ultralam Bondply 3850HT layer of thickness $25 \mu \mathrm{~m}$ is added. A high impedance line is added to the main feed to connect it to the DC voltage feeding network. A radial stub is also added to this line to ensure that RF signals do not reach the DC feeding network.


Figure A.10: (a) The different layers of the design, (b) the different feeds in the model and (c) the
LC cavities with the DC control.


Figure A.11: 2x8 Reconfigurable array fed using a rectangular waveguide to microstrip line transition.

The antenna array is designed to operate at 72 GHz with a patch length of $\mathrm{L}_{\mathrm{P}}=1.04$ mm , a width $\mathrm{W}_{\mathrm{P}}=1.35 \mathrm{~mm}$. The DC blocking gap on the ground plane is optimized to be 0.11 mm . The separation between the patches is $0.9 \lambda$. The meander line length is optimized to be $\mathrm{L}_{\mathrm{M}}=1.52 \mathrm{~mm}$, and the separation between the lines is $\mathrm{S}=0.25 \mathrm{~mm}$.

## - Results

## a) CPW feeding technique:

The array resonates between $71.5-73 \mathrm{GHz}$ when all LC cavities have an $\varepsilon_{\mathrm{r}}=3.03$, with a bandwidth of 5 GHz as shown in Figure A.12. The array shows two main beams in the $\mathrm{H}-$ plane with a gain of 12 dB at the resonant frequency. The beams' directions, when steering is not activated, are $\pm 30^{\circ}$. Another resonance is seen at 79 GHz which is caused by the DC feeding network and no proper radiation is seen at this frequency without the appropriate matching that could be insured by the LC.

For this array design, it is noticed that the LC cavities under each patch have the role of switching the frequency of operation of the array and securing a better match at the frequency where the beam steering is desired. The cavities under the meander lines have the role of changing the electric length of the meander line and consequently changing the phase at the input of each patch. It is noticed that the phases at the input of the patches of the same row should be the same in order for beam steering to take place.


Figure A.12: The Reflection Coefficient of the array in different configuration scenarios.


Figure A.13: The radiation pattern of the CPW fed array showing beam steering at 84.5 GHz in the H -Plane.

When $\varepsilon_{\mathrm{r}}=2.36$ for the top row and $\varepsilon_{\mathrm{r}}=3.03$ for the bottom row, beam steering of $+15^{\circ}$ in $\theta$ is achieved at the resonant frequency. When $\varepsilon_{\mathrm{r}}=3.03$ for the top row and $\varepsilon_{\mathrm{r}}=2.36$ for the bottom row, beam steering of $-15^{\circ}$ in $\theta$ is achieved at the resonant frequency as shown in Figure A.13. The change in the $\varepsilon_{\mathrm{r}}$ for the meander line cavities does not affect the resonant frequency of the array as shown Figure A.12.

## b) Rectangular waveguide feeding technique.

The array is well matched between $72.5-77 \mathrm{GHz}$ when all LC cavities have an $\varepsilon_{\mathrm{r}}=3.03$, with a bandwidth of 4.5 GHz as shown in Figure A.14. The array shows two main beams in the H -plane with a gain of 12 dB at the resonant frequency. The beams' directions, when
steering is not activated, are $\pm 30^{\circ}$. Another resonance is seen at 79 GHz which is caused by the DC feeding network and no proper radiation is seen at this frequency without the appropriate matching that could be insured by the LC.

Similar to the previous array design, it is noticed that the LC cavities under each patch have the role of switching the frequency of operation of the array and securing a better match at the frequency where the beam steering is desired. The cavities under the meander lines have the role of changing the electric length of the meander line and consequently changing the phase at the input of each patch. It is noticed that the phases at the input of the patches of the same row should be the same in order for beam steering to take place.


Figure A.14: The Reflection Coefficient of the array in different configuration scenarios


Figure A.15: The radiation pattern of the rectangular waveguide fed array showing beam steering at different frequencies in the H-Plane.

When $\varepsilon_{r}=2.36$ for the top row and $\varepsilon_{r}=3.03$ for the bottom row, beam steering of $+20^{\circ}$ in $\theta$ is achieved at the resonant frequency. When $\varepsilon_{\mathrm{r}}=3.03$ for the top row and $\varepsilon_{\mathrm{r}}=2.36$ for the bottom row, beam steering of $-20^{\circ}$ in $\theta$ is achieved at the resonant frequency as shown in Figure A.15. The change in the $\varepsilon_{r}$ for the meander line cavities does not affect the resonant frequency of the array as shown Figure A.14.

## - Fabrication and Measurement:

The array fed using a rectangular waveguide is fabricated. The different layers of the design are fabricated using mechanical milling (LPKF ProtoMat S63) and chemical etching as shown in Figure A.16. The different layers are assembled, aligned and glued together. The LC cavity is filled using a point by point filling technique. Then the whole designed is sealed together in order to make sure that the LC doesn't leak as shown in Figure A.17.


Figure A.16: Antenna fabrication: a) Ground plane layer, b) Feeding network layer and LC cavities, c) Rectangular patches.


Figure A.17: Assembled array design.


Figure A.18: Fabricated antenna array input reflection coefficient.

The input reflection coefficient of the array is measured. The results show that the array is operational between $71-79 \mathrm{GHz}$. When the biasing voltage at the LC cavities underneath the patches is changed, it can be noticed that the resonant frequency shifts from 77.5 to 76.2 GHz as shown in Figure A.18. This is an expected behavior from the array design that can also be seen in the simulated results where the frequency shifts from 77 to 75.8 GHz. A shift in frequency between simulated and measured results is noticed because of some fabrication errors and misalignment between the different layers.

This type of reconfigurable arrays relying on liquid crystal in its nematic form is hard to fabricate. The addition of the glue layers along with the liquid crystal reduce the efficiency of such antenna system.

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