# The first variation of the total mass of log-concave functions and related inequalities 

Andrea Colesanti ${ }^{\text {a,* }}$, Ilaria Fragalà ${ }^{\text {b }}$<br>${ }^{a}$ Dipartimento di Matematica "U. Dini", Università degli Studi di Firenze, Viale Morgagni 67/A, 50134 Firenze, Italy<br>${ }^{\mathrm{b}}$ Dipartimento di Matematica, Politecnico di Milano, Piazza Leonardo Da Vinci 32, 20133 Milano, Italy

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#### Abstract

On the class of log-concave functions on $\mathbb{R}^{n}$, endowed with a suitable algebraic structure, we study the first variation of the total mass functional, which corresponds to the volume of convex bodies when restricted to the subclass of characteristic functions. We prove some integral representation formulae for such a first variation, which suggest to define in a natural way the notion of area measure for a log-concave function. In the same framework, we obtain a functional counterpart of Minkowski's first inequality for convex bodies; as corollaries, we derive a functional form of the isoperimetric inequality, and a family of logarithmic-type Sobolev inequalities with respect to log-concave probability measures. Finally, we propose a suitable functional version of the classical Minkowski's problem for convex bodies, and prove some partial results towards its solution.


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## Contents

1. Introduction. ..... 709
2. Preliminaries ................................................................................................................... 714

[^0]2.1. Notation and background ..... 714
2.2. Functional setting ..... 716
3. Differentiability of the total mass functional ..... 719
3.1. Existence of the first variation ..... 719
3.2. Computation of the first variation in some special cases ..... 722
4. Integral representation of the first variation ..... 724
5. A functional form of Minkowski's first inequality ..... 733
6. Isoperimetric and log-Sobolev inequalities for log-concave functions ..... 736
7. About the Minkowski's problem ..... 739
Acknowledgments ..... 742
Appendix. ..... 742
References ..... 748

## 1. Introduction

This article regards log-concave functions defined in $\mathbb{R}^{n}$, i.e. functions of the form

$$
f: \mathbb{R}^{n} \rightarrow \mathbb{R}, \quad f=e^{-u}
$$

where $u: \mathbb{R}^{n} \rightarrow \mathbb{R} \cup\{+\infty\}$ is convex.
In the last decades the interest in log-concave functions has been considerably increasing, strongly motivated by the analogy between these objects and convex bodies (convex compact subsets of $\mathbb{R}^{n}$ ).

The first breakthrough in the discovery of parallel behaviours of convex bodies and logconcave functions, was the Prékopa-Leindler inequality, named after the two Hungarian mathematicians who proved it in the seventies [20,25-27]. It states that, for any given functions $f, g, h \in L^{1}\left(\mathbb{R}^{n} ; \mathbb{R}_{+}\right)$which satisfy, for some $t \in(0,1)$, the pointwise inequality

$$
h((1-t) x+t y) \geq f(x)^{1-t} g(y)^{t} \quad \forall x, y \in \mathbb{R}^{n}
$$

it holds

$$
\begin{equation*}
\int_{\mathbb{R}^{n}} h \geq\left(\int_{\mathbb{R}^{n}} f\right)^{1-t}\left(\int_{\mathbb{R}^{n}} g\right)^{t} \tag{1.1}
\end{equation*}
$$

Moreover, it was proved by Dubuc in [9] that the equality sign holds in (1.1) if and only if the functions $f$ and $g$ are log-concave and translates, meaning that $f(x)=g\left(x-x_{0}\right)$ for some $x_{0} \in \mathbb{R}^{n}$.

If $K$ and $L$ are measurable subsets of $\mathbb{R}^{n}$ such that also their Minkowski's combination $(1-t) K+t L$ is measurable, by applying the Prékopa-Leindler inequality with $f, g$ and $h$ equal respectively to the characteristic functions of $K, L$ and $(1-t) K+t L$, one obtains

$$
V((1-t) K+t L) \geq V(K)^{1-t} V(L)^{t} .
$$

This is an equivalent formulation of the classical Brunn-Minkowski inequality

$$
\begin{equation*}
V((1-t) K+t L)^{1 / n} \geq(1-t) V(K)^{1 / n}+t V(L)^{1 / n} \tag{1.2}
\end{equation*}
$$

which holds with equality sign if and only if $K$ and $L$ belong to the class $\mathcal{K}^{n}$ of convex bodies in $\mathbb{R}^{n}$ and are homothetic, namely they agree up to a translation and a dilation.

The geometric inequality (1.2) is a cornerstone in Convex Geometry: it has many important consequences, such as the isoperimetric inequality for convex bodies, and the uniqueness issue in the solution of the Minkowski's problem (see the survey paper [12] for an overview). On the other hand, in view of its functional form, inequality (1.1) is somehow more "flexible", and finds many applications in different fields, such as convex geometry, probability, mass transportation; we refer the reader to [2,3,32] for more information on Prékopa-Leindler inequality, including proofs and bibliographical references.

In the same way as (1.1) paraphrases (1.2) into the realm of functions, recently analytic versions of other geometric inequalities have been studied. In particular, we mention the socalled Blaschke-Santaló inequality, involving the product of the volume of a convex body and its polar: functional versions of it have been achieved in [2,1,11,19,10]. Let us also emphasize that a suitable notion of mean width for log-concave functions has been introduced by Klartag and Milman in [17], where some related Urysohn-type inequality are also proved; a short time ago, these topics have been further developed by Rotem in [29,30]. We also refer to the papers [24,23,5,7], which contain further developments of the results presented here, or investigations on related topics.

In the same spirit, the aim of this paper is to cast some more light upon the geometry of log-concave functions, and to propose functional counterparts of some classical quantities and inequalities in Convex Geometry, that we briefly recall below (for more details, we refer to [31]).

Going back to the Brunn-Minkowski inequality, it admits a sort of "differential version", the so-called Minkowski's first inequality, which reads

$$
\begin{equation*}
V_{1}(K, L):=\frac{1}{n} \lim _{t \rightarrow 0^{+}} \frac{V(K+t L)-V(K)}{t} \geq V(K)^{\frac{n-1}{n}} V(L)^{\frac{1}{n}} \quad \forall K, L \in \mathcal{K}^{n} . \tag{1.3}
\end{equation*}
$$

Inequality (1.3) can be easily obtained from (1.2), and it is in fact equivalent to it. Notice that, when $L$ is the unit ball, $V_{1}(K, L)$ is just the perimeter of $K$, up to a factor $n$, and (1.3) becomes the isoperimetric inequality in the class of convex bodies.

The term $V_{1}(K, L)$, which is one of the mixed volumes of $K$ and $L$, admits a very simple and elegant integral representation:

$$
\begin{equation*}
V_{1}(K, L)=\frac{1}{n} \int_{S^{n-1}} h_{L} d \sigma_{K}, \tag{1.4}
\end{equation*}
$$

where $h_{L}$ is the support function of $L$, and $\sigma_{K}$ is the area measure of $K$. In view of (1.4), the measure $\sigma_{K}$ is usually interpreted as the first variation of volume with respect to the Minkowski's addition. The classical Minkowski's problem consists in retrieving $K$ from its surface area measure, and it is well-known that it admits a unique solution up to translations. More precisely, given any measure $\eta$ on the unit sphere $S^{n-1}$ which satisfies the compatibility conditions of having barycentre at the origin and being not concentrated on an equator, there exists a convex body, unique up to translations, such that $\eta=\sigma_{K}$.

Our main goals are to provide a functional version of Minkowski's first inequality (1.3), of the representation formula (1.4), and of the Minkowski's problem. In this perspective, a crucial issue is to identify a good notion of "area measure" for a log-concave function. To that aim, we pursue a quite natural idea, namely we replace the volume of a convex body by the integral of a log-concave function: we set

$$
J(f)=\int_{\mathbb{R}^{n}} f d x,
$$

and we compute the first variation of $J$ at $f$ with respect to suitable perturbations.
Actually, log-concave functions can be equipped with two internal operations: a sum and a multiplication by positive reals, that will be denoted respectively by $\oplus$ and $\cdot$, and can be characterized as follows (see Section 2 for a more rigorous presentation). If $f=e^{-u}$ and $g=e^{-v}$ are log-concave functions and $\alpha, \beta>0$, then

$$
\begin{equation*}
\alpha \cdot f \oplus \beta \cdot g:=e^{-w}, \quad \text { where } w^{*}=\alpha u^{*}+\beta v^{*} . \tag{1.5}
\end{equation*}
$$

Here * denotes as usual the Fenchel conjugate of convex functions. In other words, if we write a generic log-concave function as $e^{-u}$, the operations introduced in (1.5) are linear with respect to $u^{*}$. In particular, since the Fenchel conjugate of the indicatrix $I_{K}$ of a convex body (see the definition in Section 2.1) is precisely its support function $h_{K}$, one has

$$
\alpha \cdot \chi_{K} \oplus \beta \cdot \chi_{L}=\chi_{\alpha K+\beta L}
$$

Therefore, definition (1.5) can be seen as a natural extension to the class of log-concave functions of the Minkowski's structure on convex bodies.

In this framework, for a pair of log-concave functions $f$ and $g$, we study the quantity

$$
\begin{equation*}
\delta J(f, g):=\lim _{t \rightarrow 0^{+}} \frac{J(f \oplus t \cdot g)-J(f)}{t} . \tag{1.6}
\end{equation*}
$$

Let us point out that, red within this formalism, the above quoted works [17,29,30] are concerned precisely with the limit in (1.6), in the special case when $f$ is equal to $\gamma_{n}$, the density of the Gaussian measure in $\mathbb{R}^{n}$. In fact, to some extent, $\gamma_{n}$ plays the role of the unit ball in the class of log-concave functions. Thus, according to [17], the mean width of a log-concave function $g$ is given by $\delta J\left(\gamma_{n}, g\right)$, by analogy with the mean width of a convex body $K$ which is given, up to a constant depending on the dimension, by $V_{1}(B, K)$. We also mention the paper [16] by Klartag (see in particular Section 3), where a limit similar to (1.6) is considered, in the class of $s$-concave functions endowed with the appropriate algebraic operations, in order to derive several functional inequalities.

When $f$ and $g$ are arbitrary log-concave functions, the limit in (1.6) exists under the fairly weak condition $J(f)>0$. In Section 3.1 we give a rigorous proof of this fact, already pointed out in [17], and we show that the condition $J(f)>0$ is not necessary in the one dimensional case. Moreover we give simple examples which reveal that $\delta J(f, g)$ may become negative or $+\infty$ (indeed, whereas $V(K+t L)$ is a polynomial in $t$ for every $K$ and $L$ in $\mathcal{K}^{n}$, this is no longer true in general for $J(f \oplus t \cdot g)$ ). Then in Section 3.2 we compute $\delta J(f, g)$ in some special cases: the case when $f=g$, which brings into play the entropy of $f$ :

$$
\operatorname{Ent}(f)=\int_{\mathbb{R}^{n}} f \log f d x-J(f) \log J(f)
$$

and the case when the logarithms of $f$ and $g$ are powers of support functions of convex bodies. In the latter case, in order to give the explicit expression of the first variation, we exploit an integral representation formula for the derivative of $p$-mixed volume due to Lutwak (see [21,22]).

To go farther than these special cases, in Section 4 we come to the problem at the core of the paper, namely the problem of giving some general integral representation formula for $\delta J(f, g)$. We are able to achieve such a representation in two distinct settings: when the finiteness domains
of $u=-\log f$ and $v=-\log g$ are the whole space $\mathbb{R}^{n}$, and when such domains are smooth strictly convex bodies. In both cases we have to assume further properties on $u$ and $v$, concerning regularity, growth at the boundary of their domain, and strict convexity. To be more precise, our integral representation formulae are settled in the classes $\mathcal{A}^{\prime}, \mathcal{A}^{\prime \prime}$ of log-concave functions $f=e^{-u}$ such that $u$ belongs respectively to

$$
\begin{aligned}
\mathcal{L}^{\prime}:= & \left\{u \in \mathcal{L}: \operatorname{dom}(u)=\mathbb{R}^{n}, u \in \mathcal{C}_{+}^{2}\left(\mathbb{R}^{n}\right), \lim _{\|x\| \rightarrow+\infty} \frac{u(x)}{\|x\|}=+\infty\right\}, \\
\mathcal{L}^{\prime \prime}:= & \left\{u \in \mathcal{L}: \operatorname{dom}(u)=K \in \mathcal{K}^{n} \cap \mathcal{C}_{+}^{2}, u \in \mathcal{C}_{+}^{2}(\operatorname{int}(K)) \cap \mathcal{C}^{0}(K),\right. \\
& \left.\lim _{x \rightarrow \partial K}\|\nabla u(x)\|=+\infty\right\} .
\end{aligned}
$$

Here the notation $\mathcal{C}_{+}^{2}$, used for functions and sets, has the following standard meaning: when it is referred to a function $u$, it means that $u \in \mathcal{C}^{2}$ and the Hessian matrix of $u$ is positive definite at each point; when it is referred to a convex body $K$, it means that $\partial K \in \mathcal{C}^{2}$ and the Gauss curvature is everywhere strictly positive.

After proving that $\mathcal{A}^{\prime}$ and $\mathcal{A}^{\prime \prime}$ are both closed with respect to the operations $\oplus$ and $\cdot$ (see Lemma 4.9), we state our main results, which are valid under the assumption that the perturbation $g$ is "controlled" by the perturbed function $f$ (see Definition 4.4 for the precise statement of this assumption, which is not necessary in the one dimensional case). In Theorem 4.5 we prove that, when $f, g \in \mathcal{A}^{\prime}, \delta J(f, g)$ is finite and is given by

$$
\begin{equation*}
\delta J(f, g)=\int_{\mathbb{R}^{n}} v^{*}(\nabla u(x)) f(x) d x \tag{1.7}
\end{equation*}
$$

In Theorem 4.6 we prove that, when $f, g \in \mathcal{A}^{\prime \prime}, \delta J(f, g)$ is finite and is given by

$$
\begin{equation*}
\delta J(f, g)=\int_{K} v^{*}(\nabla u(x)) f(x) d x+\int_{\partial K} h_{L}\left(v_{K}(x)\right) f(x) d \mathcal{H}^{n-1} \tag{1.8}
\end{equation*}
$$

where $K=\operatorname{dom}(u), v_{K}$ is the unit outer normal to $\partial K, L=\operatorname{dom}(v)$, and $h_{L}$ is the support function of $L$. The proof of these results is quite delicate and requires a careful analysis, see Section 4.

If we perform the change of variable $\nabla u(x)=y$ in (1.7), it becomes

$$
\begin{equation*}
\delta J(f, g)=\int_{\mathbb{R}^{n}} v^{*} d \mu(f), \quad d \mu(f):=f(y) e^{-\left\langle y, \nabla u^{*}(y)\right\rangle+u^{*}(y)} \operatorname{det}\left(\nabla^{2} u^{*}(y)\right) d y . \tag{1.9}
\end{equation*}
$$

Comparing (1.9) with (1.4), we are lead to identify the measure $\mu(f)$ as the area measure of a function $f$ in the class $A^{\prime}$. (Under this point of view, $v^{*}$ plays the role of the support function of $g$, as in [17]; this interpretation is quite natural in view of the fact that the algebraic structure we put on log-concave functions $e^{-u}$ is linear with respect to $u^{*}$, in the same way as the Minkowski's structure on $\mathcal{K}^{n}$ is linear with respect to support functions). Similarly, with the changes of variable $\nabla u(x)=y$ and $\nabla v_{K}(y)=\xi$, (1.8) becomes

$$
\begin{align*}
\delta J(f, g) & =\int_{\mathbb{R}^{n}} v^{*} d \mu(f)+\int_{S^{n-1}} h_{L} d \sigma(f), \quad d \mu(f) \text { as above } \\
d \sigma(f) & :=f\left(v_{K}^{-1}(\xi)\right) d \sigma_{K}(\xi) \tag{1.10}
\end{align*}
$$

Hence, within the class $\mathcal{A}^{\prime \prime}$, the notion of area measure of $f$ is provided by the pair $(\mu(f), \sigma(f))$ (notice that the former is a measure on $\mathbb{R}^{n}$, the latter on $S^{n-1}$ ).

Having the above representation formulae at our disposal, we then turn attention to functional inequalities involving $\delta J(f, g)$. Our approach is similar to the one used by Klartag in [16] for the class of $s$-concave functions. In Section 5, we prove the following functional form of Minkowski's first inequality (1.3) (see Theorem 5.1):

$$
\begin{equation*}
\delta J(f, g) \geq J(f)[\log J(g)+n]+\operatorname{Ent}(f) \tag{1.11}
\end{equation*}
$$

with equality sign if and only if there exists $x_{0} \in \mathbb{R}^{n}$ such that $g(x)=f\left(x-x_{0}\right) \forall x \in \mathbb{R}^{n}$. Loosely speaking, (1.3) can be proved taking the right derivative at $t=0$ of both sides of the Brunn-Minkowski inequality (1.2), and inequality (1.11) is obtained by adapting this idea to the Prékopa-Leindler inequality, and using Dubuc's characterization of the equality case.

In Section 6 we show that, by combining the abstract inequality (1.11) with the above representation formulae for $\delta J(f, g)$, further functional inequalities come out.

Firstly, we define the perimeter of a function $f \in \mathcal{A}^{\prime}$ in the natural way, that is as $P(f):=$ $\delta J\left(f, \gamma_{n}\right)$, and we show that, under suitable assumptions, the following functional version of the isoperimetric inequality holds (see Proposition 6.2):

$$
P(f)=\frac{1}{2} \int_{\mathbb{R}^{n}} \frac{\|\nabla f\|^{2}}{f} d x+\left(\log c_{n}\right) J(f) \geq n J(f)+\operatorname{Ent}(f) .
$$

Here $c_{n}:=(2 \pi)^{-n / 2}$, and the inequality becomes an equality if and only if there exists $x_{0} \in \mathbb{R}^{n}$ such that $f(x)=\gamma_{n}\left(x-x_{0}\right) \forall x \in \mathbb{R}^{n}$.

Then we derive a family of inequalities of logarithmic Sobolev type for probability measures $v$ with a log-concave density $v$ : under suitable assumptions on $v, a$ and $h$, including the existence of a positive constant $c$ such that $\nabla^{2} v$ is bounded below by $c$ times the identity matrix, we obtain (see Proposition 6.3)

$$
\begin{align*}
& \int_{\mathbb{R}^{n}} a(h) \log a(h) d v-\left(\int_{\mathbb{R}^{n}} a(h) d v\right) \log \left(\int_{\mathbb{R}^{n}} a(h) d v\right) \\
& \quad \leq \frac{1}{c} \int_{\mathbb{R}^{n}} \frac{\left(a^{\prime}(h)\right)^{2}}{a(h)}\|\nabla h\|^{2} d v \tag{1.12}
\end{align*}
$$

In particular, by choosing $v=\gamma_{n} d x$ and $a(t)=t^{2}$, we recover Gross' logarithmic Sobolev inequality for the Gaussian measure. We point out that our approach allows much more general choices of $v$ and $a$; on the other hand, as a drawback, the validity of (1.12) is obtained under some further restrictions on $h$.

Finally, in Section 7 we move few steps towards the solution of the Minkowski's problem for log-concave functions. As a natural extension of the Minkowski's problem for convex bodies, such a problem can be formulated as follows: retrieve a log-concave function given the first variation of its total mass functional. Clearly, in view of (1.9) and (1.10), the prescribed first variation will consist of a single measure on $\mathbb{R}^{n}$ or of a pair of measures (the first on $\mathbb{R}^{n}$ and the second on $S^{n-1}$ ), depending on whether we want to solve the problem in the class $\mathcal{A}^{\prime}$ or $\mathcal{A}^{\prime \prime}$, respectively. We establish a uniqueness result for both these problems (see Proposition 7.4), and we find some necessary conditions for the existence of a solution, which are quite similar to those mentioned before for the classic Minkowski's problem (see Proposition 7.2). However, differently from the case of convex bodies, it turns out that such conditions are in general not
sufficient, as the analysis of the one dimensional case easily shows. Thus, at this stage, some substantial difference between the geometric and the functional setting emerges, which deserves in our opinion further investigation.
Added in proof. When the publication of the present paper was in final part, the authors learned about the preprint [8], by Dario Cordero-Erausquin and Bo'az Klartag, which is closely related to the Minkowski problem introduced in Section 7 of this paper.

## 2. Preliminaries

### 2.1. Notation and background

We work in the $n$-dimensional Euclidean space $\mathbb{R}^{n}, n \geq 1$, endowed with the usual scalar product $\langle x, y\rangle$ and norm $\|x\|$; we set $B_{r}:=\left\{x \in \mathbb{R}^{n}:\|x\| \leq r\right\}$.

For $m \leq n$, we indicate by $\mathcal{H}^{m}$ the $m$-dimensional Hausdorff measure; integration with respect to the Lebesgue measure $\mathcal{H}^{n}$ is abbreviated by $d x$.

We denote by $\mathcal{K}^{n}$ the class of convex bodies (compact convex sets) in $\mathbb{R}^{n}$, and by $\mathcal{K}_{0}^{n}$ the subclass of convex bodies $K$ whose relative interior $\operatorname{int}(K)$ is nonempty. We indicate by $V(K)=\mathcal{H}^{n}(K)$ the $n$-dimensional volume of $K \in \mathcal{K}^{n}$.

Given $K \in \mathcal{K}_{0}^{n}$, we denote by $\nu_{K}$ its Gauss map (i.e., the map which associates with every point $x \in \partial K$ the subset of $S^{n-1}$ given by the unit outer normal vectors to $\partial K$ at $x$ ), by $\sigma_{K}=\left(v_{K}\right)_{\sharp}\left(\mathcal{H}^{n-1}\llcorner\partial K)\right.$ its surface area measure, and by $P(K)=\int_{S^{n-1}} d \sigma_{K}=\mathcal{H}^{n-1}(\partial K)$ its perimeter. We say that $K$ is $\mathcal{C}_{+}^{2}$ if its boundary $\partial K$ is of class $\mathcal{C}^{2}$ with strictly positive Gaussian curvature.

For any $K \in \mathcal{K}^{n}$, we adopt the standard notation $h_{K}$ for the support function of $K$, defined by

$$
h_{K}(x):=\sup _{y \in K}\langle x, y\rangle \quad \forall x \in \mathbb{R}^{n} .
$$

We recall that the polar body $K^{o}$ of $K$ is given by

$$
K^{o}:=\left\{y \in \mathbb{R}^{n}:\langle x, y\rangle \leq 1 \forall x \in K\right\}
$$

if $0 \in \operatorname{int}(K)$, the support function of $K$ agrees with the gauge function of $K^{o}$, namely

$$
h_{K}(x)=p_{K^{o}}(x):=\inf \left\{t \geq 0: x \in t K^{o}\right\} .
$$

We denote by $I_{K}$ and $\chi_{K}$ the indicatrix function and characteristic function of $K$, defined respectively by

$$
I_{K}(x):=\left\{\begin{array}{ll}
0 & \text { if } x \in K \\
+\infty & \text { if } x \notin K,
\end{array} \quad \chi_{K}(x):= \begin{cases}1 & \text { if } x \in K \\
0 & \text { if } x \notin K\end{cases}\right.
$$

Let $u: \mathbb{R}^{n} \rightarrow \mathbb{R} \cup\{+\infty\}$ be a convex function. We set

$$
\operatorname{dom}(u)=\left\{x \in \mathbb{R}^{n}: u(x) \in \mathbb{R}\right\}
$$

By the convexity of $u, \operatorname{dom}(u)$ is a convex set. We say that $u$ is proper if $\operatorname{dom}(u) \neq \emptyset$. We say that $u$ is of class $\mathcal{C}_{+}^{2}$ if it is twice differentiable on $\operatorname{int}(\operatorname{dom}(u))$, with a positive definite Hessian matrix. We denote by epi $(u)$ the epigraph of $u$.

We recall that the Fenchel conjugate of $u$ is the convex function defined by:

$$
u^{*}(y)=\sup _{x \in \mathbb{R}^{n}}\langle x, y\rangle-u(x) \quad \forall y \in \mathbb{R}^{n} .
$$

On the class of convex functions from $\mathbb{R}^{n}$ to $\mathbb{R} \cup\{+\infty\}$, we consider the operation of infimal convolution, defined by

$$
\begin{equation*}
u \square v(x):=\inf _{y \in \mathbb{R}^{n}}\{u(x-y)+v(y)\} \quad \forall x \in \mathbb{R}^{n}, \tag{2.1}
\end{equation*}
$$

and the following right scalar multiplication by a nonnegative real number $\alpha$ :

$$
(u \alpha)(x):=\left\{\begin{array}{ll}
\alpha u\left(\frac{x}{\alpha}\right) & \text { if } \alpha>0  \tag{2.2}\\
I_{\{0\}} & \text { if } \alpha=0
\end{array} \quad \forall x \in \mathbb{R}^{n} .\right.
$$

Notice that these operations are convexity preserving, and that the function $I_{\{0\}}$ acts as the identity element in (2.1).

The proposition below gathers some elementary properties of the Fenchel conjugate, in particular about its behaviour with respect to the operations defined above.

Proposition 2.1. Let $u: \mathbb{R}^{n} \rightarrow \mathbb{R} \cup\{+\infty\}$ be a convex function. Then:
(i) it holds $u^{*}(0)=-\inf (u)$; in particular, $\inf (u)>-\infty$ implies $u^{*}$ proper;
(ii) if $u$ is proper, then $u^{*}(y)>-\infty \forall y \in \mathbb{R}^{n}$;
(iii) $\operatorname{dom}(u \square v)=\operatorname{dom}(u)+\operatorname{dom}(v)$;
(iv) $(u \square v)^{*}=u^{*}+v^{*}$;
(v) $(u \alpha)^{*}=\alpha u^{*}$.in

Proof. Items (i), (ii) and (v) are straightforward consequences of the definition of Fenchel conjugate; items (iii) and (iv) can be found in [28, p. 34] and [28, Theorem 16.4], respectively.

Given a differentiable real valued function $u$ on an open subset $C$ of $\mathbb{R}^{n}$, the Legendre conjugate of the pair $(C, u)$ is defined to be the pair $(D, v)$, where $D$ is the image of $C$ through the gradient mapping $\nabla u$, and

$$
v(y)=\left\langle\nabla u^{-1}(y), y\right\rangle-u\left(\nabla u^{-1}(y)\right) \quad \forall y \in D,
$$

where $\nabla u^{-1}(y):=\{x: \nabla u(x)=y\}$. The above definition of $v$ is well posed whenever, for any $y \in D$, the value of $\langle x, y\rangle-u(x)$ turns out to be independent from the choice of the point $x \in \nabla u^{-1}(y)$.

Following [28, Section 26], we say that a pair $(C, u)$ is a convex function of Legendre type if:
(a) $C$ is a nonempty open convex set;
(b) $u$ is differentiable and strictly convex on $C$;
(c) $\lim _{i}\left\|\nabla u\left(x_{i}\right)\right\| \rightarrow+\infty$ whenever $\left\{x_{i}\right\} \subset C$ is a sequence converging to some $x \in \partial C$.

Within the class of convex functions of Legendre type, Fenchel and Legendre conjugates may be identified according to proposition below [28, Theorem 26.5].

Proposition 2.2. Let $u: \mathbb{R}^{n} \rightarrow \mathbb{R} \cup\{+\infty\}$ be a closed convex function, and set $C:=$ $\operatorname{int}(\operatorname{dom}(u)), C^{*}:=\operatorname{int}\left(\operatorname{dom}\left(u^{*}\right)\right)$. Then $(C, u)$ is a convex function of Legendre type if and only if $\left(C^{*}, u^{*}\right)$ is. In this case, $\left(C^{*}, u^{*}\right)$ is the Legendre conjugate of $(C, u)$ (and conversely). Moreover, $\nabla u: C \rightarrow C^{*}$ is a continuous bijection, and the inverse map of $\nabla u$ is precisely $\nabla u^{*}$.

### 2.2. Functional setting

Let us introduce the classes of functions we deal with throughout the paper.
Definition 2.3. We set:

$$
\begin{aligned}
& \mathcal{L}:=\left\{u: \mathbb{R}^{n} \rightarrow \mathbb{R} \cup\{+\infty\} \mid u \text { proper, convex, } \lim _{\|x\| \rightarrow+\infty} u(x)=+\infty\right\}, \\
& \mathcal{A}:=\left\{f: \mathbb{R}^{n} \rightarrow \mathbb{R} \mid f=e^{-u}, u \in \mathcal{L}\right\} .
\end{aligned}
$$

Below, we give some examples and basic properties of functions in $\mathcal{L}$; we show that, consequently, the class of log-concave functions $\mathcal{A}$ can be endowed with an algebraic structure which extends in a natural way the usual Minkowski's structure on $\mathcal{K}^{n}$.

Example 2.4. (i) For any $K \in \mathcal{K}^{n}$, the function $u=I_{K}$ belongs to $\mathcal{L}$. Notice that $u^{*}=h_{K}$ belongs to $\mathcal{L}$ if and only if $0 \in \operatorname{int}(K)$, which shows that the class $\mathcal{L}$ is not closed under Fenchel transform.
(ii) For any $K \in \mathcal{K}^{n}$ with $0 \in \operatorname{int}(K)$, and any $p \in[1,+\infty)$, the function $u=\frac{1}{p} h_{K}^{p}$ belongs to $\mathcal{L}$. In particular, for any $p \in[1,+\infty)$, the function $u(x)=\frac{1}{p}\|x\|^{p}$ belongs to $\mathcal{L}$.

Lemma 2.5. Let $u \in \mathcal{L}$. Then there exist constants $a$ and $b$, with $a>0$, such that

$$
\begin{equation*}
u(x) \geq a\|x\|+b \quad \forall x \in \mathbb{R}^{n} . \tag{2.3}
\end{equation*}
$$

Moreover $u^{*}$ is proper, and satisfies $u^{*}(y)>-\infty \forall y \in \mathbb{R}^{n}$.
Proof. In order to show (2.3), assume first that $0 \in \operatorname{dom}(u)$. Let $r>0$ be such that $u(x) \geq$ $1+u(0)$ if $\|x\| \geq r$; for $\|x\| \geq r$, the convexity of $u$ implies

$$
u(x) \geq u(0)+\left(u\left(\frac{r x}{\|x\|}\right)-u(0)\right) \frac{\|x\|}{r} \geq u(0)+\frac{\|x\|}{r} .
$$

Then, setting $m:=\inf (u)$, it holds

$$
u(x) \geq m-1+\frac{\|x\|}{r} \quad \text { for }\|x\| \geq r
$$

Since the above inequality is verified for $\|x\| \leq r$ as well, it holds in $\mathbb{R}^{n}$. This shows that (2.3) is satisfied by taking $a=r^{-1}$ and $b=(m-1)$. In the general case, since $u$ is proper, one can choose $x_{0} \in \operatorname{dom}(u)$, and apply the above argument to the function $u\left(x-x_{0}\right)$, which yields

$$
u(x) \geq a\left\|x+x_{0}\right\|+b \geq a\|x\|+b-a\left\|x_{0}\right\| .
$$

The properties of $u^{*}$ follow from Proposition 2.1(i) and (ii).
We now use Lemma 2.5 in order to prove that $\mathcal{L}$ is closed under the operations of infimal convolution and right scalar multiplication defined in (2.1) and (2.2).

Proposition 2.6. Let $u, v \in \mathcal{L}$ and $\alpha, \beta \geq 0$. Then $(u \alpha) \square(v \beta) \in \mathcal{L}$.
Proof. From definition (2.2) it is immediate that $(u \alpha) \in \mathcal{L}$ for any $u \in \mathcal{L}$ and $\alpha \geq 0$. So we have just to show that $u \square v$ belongs to $\mathcal{L}$ for any $u, v \in \mathcal{L}$. Set for brevity $w:=u \square v$. Clearly, $w$ is a
convex function defined in $\mathbb{R}^{n}$. Let us prove that $w$ takes values into $\mathbb{R} \cup\{+\infty\}$, is proper, and satisfies $\lim _{\|x\| \rightarrow+\infty} w(x)=+\infty$.

By Proposition 2.1(i) and (iv), we have

$$
\inf w=-w^{*}(0)=-u^{*}(0)-v^{*}(0)=\inf (u)+\inf (v)
$$

Since $\inf (u), \inf (v)>-\infty$, we infer that $\inf (w)>-\infty$, which shows that $w$ takes values into $\mathbb{R} \cup\{+\infty\}$.

By Proposition 2.1(iii), $\operatorname{dom}(w)=\operatorname{dom}(u)+\operatorname{dom}(v)$, hence the properness of both $u$ and $v$ implies the same property for $w$.

Let $u(x) \geq a\|x\|+b$ and $v(x) \geq a^{\prime}\|x\|+b^{\prime}$ according to Lemma 2.5, and set $c:=$ $\min \left\{a, a^{\prime}\right\}>0, d:=b+b^{\prime}$. Then, by using the definition of $w$ and the lower bounds satisfied by $u$ and $v$, we get

$$
w(x) \geq \inf _{y \in \mathbb{R}^{n}}\left\{a\|x-y\|+b+a^{\prime}\|y\|+b^{\prime}\right\} \geq c\|x\|+d
$$

In particular, this implies that $\lim _{\|x\| \rightarrow+\infty} w(x)=+\infty$.
We are now in a position to endow the class $\mathcal{A}$ with an addition and a multiplication by nonnegative scalars. These operations are internal to $\mathcal{A}$ thanks to Proposition 2.6.

Definition 2.7. Let $f=e^{-u}, g=e^{-v} \in \mathcal{A}$, and let $\alpha, \beta \geq 0$. We define

$$
\begin{equation*}
\alpha \cdot f \oplus \beta \cdot g=e^{-[(u \alpha) \square(v \beta)]} . \tag{2.4}
\end{equation*}
$$

Recalling (2.1) and (2.2), the explicit form of (2.4) when $\alpha$ and $\beta$ are strictly positive reads

$$
(\alpha \cdot f \oplus \beta \cdot g)(x):=\sup _{y \in \mathbb{R}^{n}} f\left(\frac{x-y}{\alpha}\right)^{\alpha} g\left(\frac{y}{\beta}\right)^{\beta}
$$

In the particular case when $\alpha=0$ and $\beta>0$, we have $(\alpha \cdot f \oplus \beta \cdot g)(x)=g\left(\frac{x}{\beta}\right)^{\beta}$. Similarly, for $\alpha>0$ and $\beta=0,(\alpha \cdot f \oplus \beta \cdot g)(x)=f\left(\frac{x}{\alpha}\right)^{\alpha}$. Finally, if $\alpha=\beta=0$ we simply have $(\alpha \cdot f \oplus \beta \cdot g)=I_{\{0\}}$.

Remark 2.8. In view of the identities

$$
\begin{aligned}
& u \square v(x)=\inf \{\mu:(x, \mu) \in \operatorname{epi}(u)+\operatorname{epi}(v)\} \\
& (u \alpha)(x)=\inf \{\mu:(x, \mu) \in \alpha \operatorname{epi}(u)\},
\end{aligned}
$$

the functional operation in (2.4) has the following geometrical interpretation: it corresponds to the Minkowski's combination with coefficients $\alpha$ and $\beta$ of the epigraphs of $u$ and $v$ (as subsets of $\mathbb{R}^{n+1}$ ).

Next proposition shows that, when restricted to suitable subclasses of $\mathcal{A}$, Definition 2.7 allows to recover different algebraic structures on convex bodies. Recall that (see [21]), for a fixed $p \in[1,+\infty)$, the $p$-sum of two convex bodies $K$ and $L$ with coefficients $\alpha$ and $\beta$ is the convex body $\alpha \cdot{ }_{p} K+{ }_{p} \beta \cdot{ }_{p} L$ defined by the equality

$$
h_{\alpha{ }_{p} K+{ }_{p} \beta \cdot{ }_{p} L}^{p}=\alpha h_{K}^{p}+\beta h_{L}^{p} .
$$

Proposition 2.9. Set

$$
\begin{aligned}
\mathcal{L}_{1} & :=\left\{h_{K^{o}}: K \in \mathcal{K}^{n}, 0 \in \operatorname{int}(K)\right\} \\
\mathcal{L}_{q} & :=\left\{\frac{1}{q}\left(h_{K^{o}}\right)^{q}: K \in \mathcal{K}^{n}, 0 \in \operatorname{int}(K)\right\}, \quad q \in(1,+\infty), \\
\mathcal{L}_{\infty} & :=\left\{I_{K}: K \in \mathcal{K}^{n}\right\} .
\end{aligned}
$$

The above subclasses of $\mathcal{L}$ are closed with respect to the operations defined in (2.1) and (2.2).
More precisely, for any $\alpha, \beta \geq 0$, and any $u$, v belonging to the same class $\mathcal{L}_{q}$, it holds

Proof. Let $u \in \mathcal{L}_{q}$. We have

$$
u^{*}= \begin{cases}I_{K^{o}} & \text { if } q=1  \tag{2.5}\\ \frac{1}{p} h_{K}^{p} & \text { if } q \in(1,+\infty) \\ h_{K} & \text { if } q=\infty\end{cases}
$$

In particular, in order to check the above expression of $u^{*}$ in case $q \in(1,+\infty)$, one can apply with $\phi(s)=\frac{s^{q}}{q}$ the following identity holding for every increasing convex function $\phi$ (see e.g. [15]):

$$
\left(\phi\left(h_{K^{o}}\right)\right)^{*}(x)=\inf _{t \geq 0}\left\{\phi^{*}(t)+t h_{K^{o}}^{*}\left(\frac{x}{t}\right)\right\} ;
$$

this yields

$$
\left(\frac{1}{q}\left(h_{K^{o}}\right)^{q}\right)^{*}(x)=\inf _{\left\{t \geq 0: x \in t K^{o}\right\}}\left\{\frac{t^{p}}{p}\right\}=\frac{1}{p} \rho_{K^{o}}^{p}(x)=\frac{1}{p} h_{K}^{p}(x) .
$$

Now, the statement of the proposition follows easily from the computation of $((u \alpha) \square(v \beta))^{*}$. Indeed, by Proposition 2.1(iv)-(v), it holds $((u \alpha) \square(v \beta))^{*}=\alpha u^{*}+\beta v^{*}$. According to (2.5), one has

$$
\alpha u^{*}+\beta v^{*}=\left\{\begin{array}{l}
\alpha I_{K^{o}}+\beta I_{L^{o}}=I_{K^{o} \cap L^{o}}=\left(h_{K^{o} \cap L^{o}}\right)^{*} \quad \text { if } q=1 \\
\frac{1}{p}\left[\alpha h_{K}^{p}+\beta h_{K}^{p}\right]=\frac{1}{p}\left[h_{\alpha \cdot p} K+_{p} \beta \cdot{ }_{p} L\right]^{p} \\
\left.\left\{\frac{1}{q}\left[h_{(\alpha \cdot p} K+_{p} \beta \cdot{ }_{p} L\right)^{o}\right]^{q}\right\}^{*} \quad \text { if } q \in(1,+\infty) \\
\alpha h_{K}+\beta h_{L}=h_{\alpha K+\beta L}=\left(I_{\alpha K+\beta L)^{*}} \quad \text { if } q=\infty .\right.
\end{array}\right.
$$

## 3. Differentiability of the total mass functional

Definition 3.1. We call total mass functional the following integral

$$
J(f)=\int_{\mathbb{R}^{n}} f(x) d x \quad \forall f \in \mathcal{A} .
$$

Remark 3.2. (i) The growth condition from below satisfied by functions in $\mathcal{L}$ according to Lemma 2.5 (see (2.3)) ensures that $J(f) \in[0,+\infty$ ) for every $f \in \mathcal{A}$.
(ii) Clearly, when $f=\chi_{K}$, one has $J(f)=V(K)$.
(iii) If $f=e^{-u}$ is such that $J(f)=0$, then $f=0 \mathcal{H}^{n}$-a.e. in $\mathbb{R}^{n}$. This implies that the convex set $\operatorname{dom}(u)$ is Lebesgue negligible, and hence its dimension does not exceed $(n-1)$.

Remark 3.3. By the Prékopa-Leindler inequality, for every $f, g \in \mathcal{A}$ and for every $t \in[0,1]$, it holds

$$
J((1-t) \cdot f \oplus t \cdot g) \geq J(f)^{1-t} J(g)^{t}
$$

with equality sign if and only if there exists $x_{0} \in \mathbb{R}^{n}$ such that $g(x)=f\left(x-x_{0}\right) \forall x \in \mathbb{R}^{n}$ (see [9,12]). Consequently, for every fixed $f, g \in \mathcal{A}$, the functions $t \mapsto \log J(f \oplus t \cdot g)$ and $t \mapsto \log J((1-t) \cdot f \oplus t \cdot g)$ turn out to be concave respectively on $[0,+\infty)$ and on $[0,1]$. We shall repeatedly exploit this concavity property in the sequel.

We are going to study the first variation of the total mass functional, with respect to the algebraic structure introduced in Definition 2.7.

Definition 3.4. Let $f, g \in \mathcal{A}$. Whenever the following limit exists

$$
\lim _{t \rightarrow 0^{+}} \frac{J(f \oplus t \cdot g)-J(f)}{t}
$$

we denote it by $\delta J(f, g)$, and we call it the first variation of $J$ at $f$ along $g$.
Remark 3.5. Let $f=\chi_{K}$ and $g=\chi_{L}$, with $K, L \in \mathcal{K}^{n}$. In this case $J(f \oplus t \cdot g)=V(K+t L)$ is a polynomial in $t$; its derivative at $t=0^{+}$is equal to $n$ times the mixed volume $V_{1}(K, L)$, and admits the integral representation

$$
\begin{equation*}
\frac{d}{d t} V(K+t L)_{\left.\right|_{t=0^{+}}}=n V_{1}(K, L)=\int_{S^{n-1}} h_{L} d \sigma_{K} \tag{3.1}
\end{equation*}
$$

Notice in particular that $\delta J\left(\chi_{K}, \chi_{L}\right)$ is nonnegative and finite, which is not always true in general for $\delta J(f, g)$ ( $c f$. the examples given in Remark 3.8 below).

Section 3.1 below is devoted to prove that $\delta J(f, g)$ exists under the fairly weak hypothesis that $J(f)$ is strictly positive. Then in Section 3.2 we show the explicit expression of $\delta J(f, g)$ in some relevant cases.

### 3.1. Existence of the first variation

Theorem 3.6. Let $f, g \in \mathcal{A}$, and assume that $J(f)>0$. Then $J$ is differentiable at $f$ along $g$, and it holds

$$
\begin{equation*}
\delta J(f, g) \in[-k,+\infty], \tag{3.2}
\end{equation*}
$$

being $k:=[\inf (-\log g)]_{+} J(f)$. In dimension $n=1$, the same conclusions continue to hold also when $J(f)=0$.

Remark 3.7. We point out that the assumption $J(f)>0$ is somehow technical; we believe that, when $J(f)=0$, Theorem 3.6 is likely true not only in dimension $n=1$ but also in higher dimensions (as it is suggested by the fact that the mixed volume $V_{1}(K, L)$ exists regardless of the condition $V(K)>0)$.

Remark 3.8. Estimate (3.2) cannot be improved, as the following examples show.
(i) Let $f=e^{-u} \in \mathcal{A}$ with $J(f)>0$, and $g=e^{-v}$, where $v(0)=1$ and $v \equiv+\infty$ in $\mathbb{R}^{n} \backslash\{0\}$. Then $u \square(v t)(x)=u(x)+t$, which implies

$$
\delta J(f, g)=J(f) \cdot \lim _{t \rightarrow 0^{+}} \frac{e^{-t}-1}{t}=-J(f)<0
$$

(ii) Let $K, L \in \mathcal{K}^{n}$ with the origin in their interior, so that $u=h_{K}, v=h_{L} \in \mathcal{L}$, and take $f=e^{-u}, g=e^{-v}$. Then $u \square(v t)=h_{K \cap L}(c f$. Proposition 2.9), and therefore

$$
\delta J(f, g)=\lim _{t \rightarrow 0^{+}}\left[\frac{1}{t} \int_{\mathbb{R}^{n}}\left(e^{-h_{K \cap L}}-e^{-h_{L}}\right) d x\right]= \begin{cases}0 & \text { if } L \subseteq K \\ +\infty & \text { otherwise }\end{cases}
$$

Prior to the proof of Theorem 3.6, we state a preliminary lemma, which will be heavily exploited also in the next section.

Lemma 3.9. Let $f=e^{-u}, g=e^{-v} \in \mathcal{A}$. For $t \geq 0$, set $u_{t}=u \square(v t)$ and $f_{t}=e^{-u_{t}}$. Assume that $v(0)=0$. Then, for every fixed $x \in \mathbb{R}^{n}, u_{t}(x)$ and $f_{t}(x)$ are respectively pointwise decreasing and increasing with respect to $t$; in particular it holds

$$
u_{1}(x) \leq u_{t}(x) \leq u(x) \quad \text { and } \quad f(x) \leq f_{t}(x) \leq f_{1}(x) \quad \forall x \in \mathbb{R}^{n}, \forall t \in[0,1]
$$

Proof. Given $t \geq 0$ and $\delta>0$, let us show that $u_{t+\delta} \leq u_{t}$, i.e.

$$
u \square(v(t+\delta)) \leq u \square(v t)
$$

If $t=0$, the above inequality reduces to $u \square(v \delta) \leq u$. This is readily checked: recalling definitions (2.1) and (2.2), from the assumption $v(0)=0$ we deduce

$$
u \square(v \delta)(x)=\inf _{y \in \mathbb{R}^{n}}\left\{u(x-y)+\delta v\left(\frac{y}{\delta}\right)\right\} \leq u(x) \quad \forall x \in \mathbb{R}^{n} .
$$

If $t>0$, for every $x \in \mathbb{R}^{n}$ we have

$$
\begin{aligned}
u \square(v(t+\delta))(x) & =\inf _{\xi \in \mathbb{R}^{n}}\left\{u(x-\xi)+(t+\delta) v\left(\frac{\xi}{t+\delta}\right)\right\} \\
& =\inf _{\xi \in \mathbb{R}^{n}}\left\{u(x-\xi)+\inf _{y \in \mathbb{R}^{n}}\left[t v\left(\frac{\xi-y}{t}\right)+\delta v\left(\frac{y}{\delta}\right)\right]\right\} \\
& =\inf _{y, z \in \mathbb{R}^{n}}\left\{u(x-y-z)+t v\left(\frac{z}{t}\right)+\delta v\left(\frac{y}{\delta}\right)\right\} \\
& =(u \square(v t)) \square(v \delta)(x) \leq u \square(v t)(x) .
\end{aligned}
$$

Thus $u_{t}$ is monotone decreasing with respect to $t$, which immediately implies that $f_{t}=e^{-u_{t}}$ is monotone increasing.
Proof of Theorem 3.6. We set

$$
\begin{equation*}
u:=-\log f, \quad v:=-\log g, \quad f_{t}:=f \oplus t \cdot g, \tag{3.3}
\end{equation*}
$$

and

$$
\begin{equation*}
d:=v(0), \quad \tilde{v}(x):=v(x)-d, \quad \tilde{g}(x):=e^{-\tilde{v}(x)}, \quad \tilde{f}_{t}:=f \oplus t \cdot \tilde{g} \tag{3.4}
\end{equation*}
$$

Up to a translation of coordinates, we may also assume without loss of generality that $\inf (v)=$ $v(0)$.

Since by construction $\tilde{v}(0)=0$, by Lemma 3.9 for every $x \in \mathbb{R}^{n}$ there exists $\tilde{f}(x):=$ $\lim _{t \rightarrow 0^{+}} \tilde{f}_{t}(x)$ and it holds $\tilde{f}(x) \geq f(x)$. As $t \rightarrow 0_{\tilde{f}}^{+}, \tilde{f}_{t}$ is pointwise decreasing by Lemma 3.9; moreover, by Lemma 2.5 and Proposition 2.6, $J\left(\tilde{f}_{1}\right)<\infty$. Hence, by monotone convergence, we have $\lim _{t \rightarrow 0^{+}} J\left(\tilde{f_{t}}\right)=J(\tilde{f})$.

Since $f_{t}(x)=e^{-d t} \tilde{f}_{t}(x)$, we have

$$
\begin{equation*}
\frac{J\left(f_{t}\right)-J(f)}{t}=J(f) \frac{e^{-d t}-1}{t}+e^{-d t} \frac{J\left(\tilde{f}_{t}\right)-J(f)}{t} \tag{3.5}
\end{equation*}
$$

Let us consider separately the two cases $J(\tilde{f})>J(f)$ and $J(\tilde{f})=J(f)$.
If $J(\tilde{f})>J(f)$, then

$$
\lim _{t \rightarrow 0^{+}} \frac{J\left(f_{t}\right)-J(f)}{t}=\lim _{t \rightarrow 0^{+}} \frac{J\left(\tilde{f}_{t}\right)-J(f)}{t}=+\infty
$$

and the statement of the theorem holds true.
If $J(\tilde{f})=J(f)$, we further distinguish the following two subcases:

$$
\exists t_{0}>0: J\left(\tilde{f}_{t_{0}}\right)=J(f) \quad \text { or } \quad J\left(\tilde{f}_{t}\right)>J(f) \quad \forall t>0
$$

In the former subcase, since by Lemma $3.9 J\left(\tilde{f}_{t}\right)$ is a monotone increasing function of $t$, necessarily it holds $J\left(\tilde{f}_{t}\right)=J(f)$ for every $t \in\left[0, t_{0}\right]$. Hence the second term in the r.h.s. of (3.5) tends to 0 , so that

$$
\lim _{t \rightarrow 0^{+}} \frac{J\left(f_{t}\right)-J(f)}{t}=-d J(f)
$$

and the statement of the theorem holds true.
In the latter subcase, we can write

$$
\begin{equation*}
\frac{J\left(\tilde{f}_{t}\right)-J(f)}{t}=\frac{\log \left(J\left(\tilde{f}_{t}\right)\right)-\log (J(f))}{t} \cdot \frac{J\left(\tilde{f}_{t}\right)-J(f)}{\log \left(J\left(\tilde{f}_{t}\right)\right)-\log (J(f))} \tag{3.6}
\end{equation*}
$$

Since $\log \left(J\left(\tilde{f}_{t}\right)\right)$ is an increasing concave function of $t$ (respectively by Lemma 3.9 and by the Prékopa-Leindler inequality, $c f$. Remark 3.2),

$$
\begin{equation*}
\exists \lim _{t \rightarrow 0^{+}} \frac{\log \left(J\left(\tilde{f}_{t}\right)\right)-\log (J(f))}{t} \in[0,+\infty] . \tag{3.7}
\end{equation*}
$$

On the other hand,

$$
\begin{equation*}
\exists \lim _{t \rightarrow 0^{+}} \frac{J\left(\tilde{f}_{t}\right)-J(f)}{\log \left(J\left(\tilde{f}_{t}\right)\right)-\log (J(f))}=J(f)>0 \tag{3.8}
\end{equation*}
$$

From (3.6)-(3.8), we infer that

$$
\begin{equation*}
\exists \lim _{t \rightarrow 0^{+}} \frac{J\left(\tilde{f}_{t}\right)-J(f)}{t} \in[0,+\infty] . \tag{3.9}
\end{equation*}
$$

Combining (3.5) and (3.9), we deduce that

$$
\begin{equation*}
\exists \lim _{t \rightarrow 0^{+}} \frac{J\left(f_{t}\right)-J(f)}{t} \in[-\max \{d, 0\} J(f),+\infty] \tag{3.10}
\end{equation*}
$$

Finally, let us show that in the one-dimensional case $\delta J(f, g)$ exists also when $J(f)=0$. We keep definitions (3.3) and (3.4). Since by assumption $\operatorname{dom}(u)$ is a Lebesgue negligible convex set, it consists of exactly one point $x_{0}$. Then

$$
u \square(\tilde{v} t)(x)=u\left(x_{0}\right)+t \tilde{v}\left(\frac{x-x_{0}}{t}\right) \quad \forall x \in \mathbb{R}, \forall t>0 .
$$

Hence

$$
\begin{align*}
\lim _{t \rightarrow 0^{+}} \frac{J\left(\tilde{f}_{t}\right)-J(f)}{t} & =\lim _{t \rightarrow 0^{+}} e^{-u\left(x_{0}\right)} \int_{\mathbb{R}} e^{-t \tilde{v}\left(x-x_{0}\right)} d x \\
& =e^{-u\left(x_{0}\right)} \mathcal{H}^{1}(\operatorname{dom}(v)) \in[0,+\infty] \tag{3.11}
\end{align*}
$$

where the last equality holds true by monotone convergence. Combining (3.5) and (3.11), we see that (3.10) remains true.

### 3.2. Computation of the first variation in some special cases

Firstly, we analyze the case $f=g$, and we show that $\delta J(f, f)$ admits a very simple representation in terms of the mass and the entropy of $f$, which is defined as follows (cf. [18]).

Definition 3.10. For every $f \in \mathcal{A}$ with $J(f)>0$, we call entropy of $f$ the following quantity:

$$
\operatorname{Ent}(f)=\int_{\mathbb{R}^{n}} f \log f d x-J(f) \log J(f) \quad \forall f \in \mathcal{A}
$$

Proposition 3.11. For every $f \in \mathcal{A}$ with $J(f)>0$, it holds $\operatorname{Ent}(f) \in(-\infty,+\infty)$ and

$$
\begin{equation*}
\delta J(f, f)=n J(f)+\int_{\mathbb{R}^{n}} f \log f d x=(n+\log J(f)) J(f)+\operatorname{Ent}(f) \tag{3.12}
\end{equation*}
$$

Proof. Since $J(f) \in(0,+\infty)$ for every $f \in \mathcal{A}$, to prove the finiteness of $\operatorname{Ent}(f)$ we have just to show that

$$
\int_{\mathbb{R}^{n}} f \log f d x \in(-\infty,+\infty)
$$

We set $u:=-\log f$ and $\Omega:=\left\{x \in \mathbb{R}^{n}: u(x) \leq 0\right\}$ (which is possibly an empty set). It holds

$$
\int_{\Omega} f \log f d x=-\int_{\Omega} f u d x<-\inf _{\Omega}(u) \int_{\Omega} f<+\infty,
$$

where in the last inequality we have used the boundedness of $u$ from below on $\Omega$ and the finiteness of $J(f)$. On the other hand, we have

$$
\int_{\mathbb{R}^{n} \backslash \Omega} f \log f d x=-\int_{\mathbb{R}^{n} \backslash \Omega} f u d x \geq-m \int_{\mathbb{R}^{n} \backslash \Omega} e^{-u(x) / 2} d x>-\infty
$$

where we have used the elementary inequality $s e^{-s / 2} \leq m:=2 / e$ holding for every $s \in \mathbb{R}_{+}$and Lemma 2.5. So we have $J(f \log f) \in(-\infty,+\infty)$.

In order to prove the representation formula (3.12), assume first that $u \geq 0$. Since $u \square(u t)=$ $u(1+t)$, we have

$$
\begin{aligned}
\frac{J(f \oplus t \cdot f)-J(f)}{t} & =\frac{1}{t}\left[(1+t)^{n} \int_{\mathbb{R}^{n}} e^{-(1+t) u} d x-\int_{\mathbb{R}^{n}} e^{-u} d x\right] \\
& =\left[\frac{(1+t)^{n}-1}{t}\right] \int_{\mathbb{R}^{n}} e^{-(1+t) u} d x+\int_{\mathbb{R}^{n}} e^{-u}\left(\frac{e^{-t u}-1}{t}\right) d x
\end{aligned}
$$

Now (3.12) follows by passing to the limit as $t \rightarrow 0^{+}$(notice indeed that by the assumption $u \geq 0$ one can apply the monotone convergence theorem).

In the general case when the assumption $u \geq 0$ is removed, we consider the function $\tilde{f}=e^{-\tilde{u}}$, where $\tilde{u}=u+c$ and $c=-\inf (u)$. One can easily check that $u \square(u t)=-c(1+t)+\tilde{u} \square(\tilde{u} t)$ and consequently $J(f \oplus t \cdot f)=e^{c(1+t)} J(\tilde{f} \oplus t \cdot \tilde{f})$. As $\tilde{u} \geq 0$, we know that $\delta J(\tilde{f}, \tilde{f})$ exists and it is finite, so the same is true for $\delta J(f, f)$. Moreover,

$$
\begin{aligned}
\delta J(f, f) & =c e^{c} J(\tilde{f})+e^{c} \delta J(\tilde{f}, \tilde{f})=c J(f)+e^{c}\left[n J(\tilde{f})-\int_{\mathbb{R}^{n}} e^{-(u+c)}(u+c) d x\right] \\
& =n J(f)+\int_{\mathbb{R}^{n}} f \log f d x .
\end{aligned}
$$

Next we show that, when $-\log f$ and $-\log g$ belong to the class $\mathcal{L}_{q}$ introduced in Proposition $2.9, \delta J(f, g)$ can be written explicitly in integral form, by using the representation formula for $p$-mixed volumes given in [21].

Proposition 3.12. Let $q \in(1,+\infty)$, and let $p:=q /(q-1)$. Let $K, L \in \mathcal{K}^{n}$ with the origin in their interior, let $u:=\frac{1}{q}\left(h_{K^{o}}\right)^{q}, v:=\frac{1}{q}\left(h_{L^{o}}\right)^{q}$, and $f:=e^{-u}, g:=e^{-v}$. Setting $c(n, q):=q^{\frac{n}{q}} \Gamma\left(\frac{n+q}{q}\right)$ (where $\Gamma$ denotes the Euler Gamma-function), there holds

$$
\begin{align*}
& J(f)=c(q, n) V(K)  \tag{3.13}\\
& \delta J(f, g)=\frac{c(q, n)}{n} \int_{S^{n-1}} h_{L}(\xi)^{p}\left(h_{K}(\xi)\right)^{1-p} d \sigma_{K}(\xi) . \tag{3.14}
\end{align*}
$$

Proof. We set for brevity $a(t)=t^{p} / p$, so that $a^{*}(t)=t^{q} / q$. We have:

$$
\begin{aligned}
J(f) & =\int_{\mathbb{R}^{n}} e^{-a^{*}\left(h_{K^{o}}\right)} d x=\int_{0}^{1} \mathcal{H}^{n}\left(\left\{x: e^{-a^{*}\left(h_{K^{o}}\right)(x)}>t\right\}\right) d t \\
& =\int_{0}^{1} \mathcal{H}^{n}\left(\left\{x: h_{K^{o}}(x)<\left(a^{*}\right)^{-1}(-\log t)\right\}\right) d t \\
& =\int_{0}^{1} \mathcal{H}^{n}\left(\left\{x: h_{K^{o}}\left(\frac{x}{\left(a^{*}\right)^{-1}(-\log t)}\right)<1\right\}\right) d t
\end{aligned}
$$

$$
\begin{aligned}
& =\int_{0}^{1}\left(\left(a^{*}\right)^{-1}(-\log t)\right)^{n} \mathcal{H}^{n}\left(\left\{y: h_{K^{o}}(y)<1\right\}\right) d t \\
& =\left\{\int_{0}^{1}\left(\left(a^{*}\right)^{-1}(-\log t)\right)^{n} d t\right\} V(K)
\end{aligned}
$$

which proves (3.13) with

$$
c(q, n):=\int_{0}^{1}\left(\left(a^{*}\right)^{-1}(-\log t)\right)^{n} d t=q^{\frac{n}{q}} \Gamma\left(\frac{n+q}{q}\right)
$$

Now we recall from Proposition 2.9 that

$$
f \oplus t \cdot g=e^{-\frac{1}{q}\left(h_{\left.(K+p t \cdot p L)^{o}\right)^{q}}\right.},
$$

which combined with (3.13) implies

$$
\delta J(f, g)=c(q, n) \lim _{t \rightarrow 0^{+}} \frac{V\left(K+{ }_{p} t \cdot{ }_{p} L\right)-V(K)}{t}
$$

Then (3.14) follows from the representation formula for $p$-mixed volumes given in [21, (IIIp)].

## 4. Integral representation of the first variation

In view of the examples in Section 3.2, it is natural to ask whether, in general, $\delta J(f, g)$ admits some kind of integral representation. In this section we show that this is true when both $f$ and $g$ belong to suitable subclasses of $\mathcal{A}$.

Let us begin by introducing the measures which intervene in the representation formulae for $\delta J(f, g)$. Such measures can be viewed as the "first variation" of $J$ in the class of log-concave functions, since they play for $f$ the same role as the surface area measure for the volume in Convex Geometry. This fact emerges in a clear way by comparing the first variation of volume in (3.1) with Theorems 4.5 and 4.6 below.

Definition 4.1. Let $f=e^{-u} \in \mathcal{A}$, and consider the gradient map $\nabla u: \operatorname{dom}(u) \rightarrow \mathbb{R}^{n}$. We set $\mu(f)$ the Borel measure on $\mathbb{R}^{n}$ defined by

$$
\mu(f):=(\nabla u)_{\sharp}\left(f \mathcal{H}^{n}\right) .
$$

When $\operatorname{dom}(u)=: K \in \mathcal{K}^{n}$, we also set $\sigma(f)$ the Borel measure on $S^{n-1}$ defined by

$$
\sigma(f):=\left(v_{K}\right)_{\sharp}\left(f \mathcal{H}^{n-1}\llcorner\partial K) .\right.
$$

Next, we define the subclasses of $\mathcal{A}$ where our integral representation formulae are settled.
Definition 4.2. We set $\mathcal{A}^{\prime}, \mathcal{A}^{\prime \prime}$ the subclasses of $\mathcal{A}$ given by functions $f$ such that $u=-\log f$ belongs respectively to

$$
\begin{aligned}
\mathcal{L}^{\prime}:= & \left\{u \in \mathcal{L}: \operatorname{dom}(u)=\mathbb{R}^{n}, u \in \mathcal{C}_{+}^{2}\left(\mathbb{R}^{n}\right), \lim _{\|x\| \rightarrow+\infty} \frac{u(x)}{\|x\|}=+\infty\right\} \\
\mathcal{L}^{\prime \prime}:= & \left\{u \in \mathcal{L}: \operatorname{dom}(u)=K \in \mathcal{K}^{n} \cap \mathcal{C}_{+}^{2}, u \in \mathcal{C}_{+}^{2}(\operatorname{int}(K)) \cap \mathcal{C}^{0}(K),\right. \\
& \left.\lim _{x \rightarrow \partial K}\|\nabla u(x)\|=+\infty\right\} .
\end{aligned}
$$

Remark 4.3. Notice that, for any $u$ belonging to $\mathcal{L}^{\prime}$ or $\mathcal{L}^{\prime \prime},(\operatorname{int}(\operatorname{dom}(u)), u)$ is a convex function of Legendre type, and $u$ is cofinite, i.e. the domain of its Fenchel conjugate is the whole $\mathbb{R}^{n}$.

Finally, we introduce the concept of an admissible perturbation.
Definition 4.4. We say that $g=e^{-v}$ is an admissible perturbation for $f=e^{-u}$ if

$$
\begin{equation*}
\exists c>0: \varphi-c \psi \text { is convex, } \quad \text { where } \varphi=u^{*} \text { and } \psi=v^{*} . \tag{4.1}
\end{equation*}
$$

Our integral representation results read as follows.
Theorem 4.5. Let $f, g \in \mathcal{A}^{\prime}$, and assume that $g$ is an admissible perturbation for $f$. Then $\delta J(f, g)$ is finite and is given by

$$
\begin{equation*}
\delta J(f, g)=\int_{\mathbb{R}^{n}} \psi d \mu(f) \tag{4.2}
\end{equation*}
$$

where $\psi=v^{*}$.
Theorem 4.6. Let $f, g \in \mathcal{A}^{\prime \prime}$, and assume that $g$ is an admissible perturbation for $f$. Then $\delta J(f, g)$ is finite and is given by

$$
\begin{equation*}
\delta J(f, g)=\int_{\mathbb{R}^{n}} \psi d \mu(f)+\int_{S^{n-1}} h_{L} d \sigma(f), \tag{4.3}
\end{equation*}
$$

where $\psi=v^{*}$ and $L=\operatorname{dom}(v)$.
Remark 4.7. For $n=1$, (4.2) and (4.3) continue to hold, possibly as an equality $+\infty=+\infty$, if the assumption that $g$ is an admissible perturbation for $f$ is removed (see the Appendix for a proof).

Remark 4.8. Under the assumptions of Theorem 4.5 or Theorem 4.6, by using the definition of push-forward measure and the change of variables $\nabla u(x)=y$, one obtains

$$
\begin{aligned}
\int_{\mathbb{R}^{n}} \psi d \mu(f) & =\int_{\operatorname{dom}(u)} \psi(\nabla u(x)) f(x) d x \\
& =\int_{\mathbb{R}^{n}} \psi(y) e^{-\langle y, \nabla \varphi(y)\rangle+\varphi(y)} \operatorname{det}\left(\nabla^{2} \varphi(y)\right) d y
\end{aligned}
$$

Similarly, under the assumptions of Theorem 4.6, it holds

$$
\begin{aligned}
\int_{S^{n-1}} h_{L} d \sigma(f) & =\int_{\partial K} h_{L}\left(v_{K}(x)\right) f(x) d \mathcal{H}^{n-1}(x) \\
& =\int_{S^{n-1}} h_{L}(\xi) f\left(v_{K}^{-1}(\xi)\right) \operatorname{det}\left(\nabla v_{K}^{-1}(\xi)\right) d \mathcal{H}^{n-1}(\xi)
\end{aligned}
$$

The proof of Theorems 4.5 and 4.6 is quite delicate and requires several preliminary lemmas, whose proof is postponed to the Appendix.

The first one establishes the closure of the two subclasses of $\mathcal{L}$ introduced in Definition 4.2 with respect to the operations of infimal convolution and right scalar multiplication.

Lemma 4.9. Let $u$ and $v$ belong both to the same class $\mathcal{L}^{\prime}$ or $\mathcal{L}^{\prime \prime}$ and, for any $t>0$, set $u_{t}:=u \square(v t)$. Then $u_{t}$ belongs to the same class as $u$ and $v$.

We now turn attention to the behaviour of the functions $u_{t}=u \square(v t)$ with respect to the parameter $t$, more precisely regarding their pointwise convergence as $t \rightarrow 0^{+}$(Lemma 4.10), and their differentiability in $t$ (Lemma 4.11).

Lemma 4.10. Let $u$ and $v$ belong both to the same class $\mathcal{L}^{\prime}$ or $\mathcal{L}^{\prime \prime}$ and, for any $t>0$, set $u_{t}:=u \square(v t)$. Assume that $v(0)=0$. Then
(i) $\forall x \in \operatorname{dom}(u), \lim _{t \rightarrow 0^{+}} u_{t}(x)=u(x)$;
(ii) $\forall E \subset \subset \operatorname{dom}(u), \lim _{t \rightarrow 0^{+}} \nabla u_{t}(x)=\nabla u$ uniformly on $E$.

The following result is a key point in the proof of Theorems 4.5 and 4.6 ; it contains an explicit expression of the pointwise derivative of $u \square(v t)$ with respect to $t$.

Lemma 4.11. Let $u$ and $v$ belong both to the same class $\mathcal{L}^{\prime}$ or $\mathcal{L}^{\prime \prime}$ and, for any $t>0$, let $u_{t}:=u \square(v t)$. Then

$$
\forall x \in \operatorname{int}\left(\operatorname{dom}\left(u_{t}\right)\right), \forall t>0, \quad \frac{d}{d t} u_{t}(x)=-\psi\left(\nabla u_{t}(x)\right), \text { where } \psi:=v^{*}
$$

Next lemma provides a summability property of the Fenchel conjugate of $u=-\log f$ with respect to the measure $\mu(f)$ introduced in Definition 4.1.

Lemma 4.12. Let $f=e^{-u} \in \mathcal{A}$, with $\varphi=u^{*} \geq 0$. Then $\varphi \in L^{1}(d \mu(f))$, namely

$$
\int_{\mathbb{R}^{n}} \varphi(\nabla u(x)) f(x) d x<+\infty
$$

Finally, when $u, v \in \mathcal{L}^{\prime \prime}$, we need an estimate for $u_{t}=u \square(v t)$ which will be exploited to deal with the boundary term in Theorem 4.6.

Lemma 4.13. Let $u, v \in \mathcal{L}^{\prime \prime}$ and, for any $t>0$, let $u_{t}=u \square(v t)$. Set $K:=\operatorname{dom}(u), L:=$ $\operatorname{dom}(v), v_{\max }:=\max _{L} v$, and $v_{\min }:=\min _{L} v$. Then, for every $x \in K+t L$, there exists $y=y(x, t) \in K \cap(x-t L)$ such that

$$
t v_{\min }+u(y) \leq u_{t}(x) \leq t v_{\max }+u(y)
$$

Proof of Theorems 4.5 and 4.6. We assume that either the hypotheses of Theorem 4.5 or the hypotheses of Theorem 4.6 are satisfied.

Throughout the proof we set

$$
\begin{aligned}
& f=e^{-u}, \quad g=e^{-v}, \quad \varphi=u^{*}, \\
& \psi=v^{*}, \quad E=\operatorname{dom}(u), \quad F=\operatorname{dom}(v),
\end{aligned}
$$

and, for every $t \geq 0$,

$$
f_{t}=f \oplus t \cdot g, \quad u_{t}=u \square(v t), \quad \varphi_{t}=\varphi+t \psi, \quad E_{t}=E+t F
$$

Let us point out that, under the assumptions of Theorem 4.5, we have $E=F=\mathbb{R}^{n}$, whereas, under the assumptions of Theorem $4.6, E$ and $F$ are convex bodies that will be named respectively $K$ and $L$.

Further, we need to 'localize' our total mass functional: for every measurable set $A \subseteq \mathbb{R}^{n}$ and any function $h \in \mathcal{A}$, we set

$$
J_{A}(h):=\int_{A} h d x .
$$

For convenience, we divide the proof into several steps.
Step 1. Decomposition.
With the notation introduced above, we can write

$$
J\left(f_{t}\right)-J(f)=J_{E}\left(f_{t}\right)-J_{E}(f)+J_{E_{t} \backslash E}\left(f_{t}\right)
$$

We are going to prove the integral representation formulae (4.2) and (4.3) by showing that:

- under the assumptions of one among Theorems 4.5 and 4.6, it holds

$$
\begin{equation*}
\lim _{t \rightarrow 0^{+}} \frac{J_{E}\left(f_{t}\right)-J_{E}(f)}{t}=\int_{\mathbb{R}^{n}} \psi d \mu(f) \tag{4.4}
\end{equation*}
$$

- under the assumptions of Theorem 4.6, it holds

$$
\begin{equation*}
\lim _{t \rightarrow 0^{+}} \frac{J_{E_{t} \backslash E}\left(f_{t}\right)}{t}=\int_{S^{n-1}} h_{L} d \sigma(f) \tag{4.5}
\end{equation*}
$$

Step 2. Reduction to the case $0 \in \operatorname{int}(F), v(0)=0, v \geq 0, \varphi \geq 0, \psi \geq 0$.
Assume that equalities (4.4) and (4.5) hold true (respectively under the assumptions of Theorems 4.5 and 4.6), when all the conditions $0 \in \operatorname{int}(F), v(0)=0, v \geq 0, \varphi \geq 0, \psi \geq 0$ are satisfied.

In the general case, up to a translation of coordinates (which does not affect $J$ ), we may assume that $\inf v=v(0)$. Since by assumption $v$ belongs to $\mathcal{L}^{\prime}$ or $\mathcal{L}^{\prime \prime}$, its minimum is necessarily attained in the interior of its domain, so we have $0 \in \operatorname{int}(F)$. If $c:=u(0)$ and $d:=v(0)$, we set

$$
\tilde{u}(x):=u(x)-c, \quad \tilde{v}(x):=v(x)-d, \quad \tilde{\varphi}(y):=(\tilde{u})^{*}(y), \quad \tilde{\psi}(y):=(\tilde{v})^{*}(y)
$$

and

$$
\tilde{f}=e^{-\tilde{u}}, \quad \tilde{g}=e^{-\tilde{v}}, \quad \tilde{f}_{t}:=\tilde{f} \oplus t \cdot \tilde{g}
$$

By construction it holds $\operatorname{dom}(\tilde{v})=F, \tilde{v}(0)=0, \tilde{v} \geq 0, \tilde{\varphi} \geq 0, \tilde{\psi} \geq 0$. Then, taking also into account that $\operatorname{dom}(\tilde{u})=E, \tilde{\psi}(y)=\psi(y)+d$, and $\tilde{f}=e^{c} f$, it holds

$$
\begin{equation*}
\lim _{t \rightarrow 0^{+}} \frac{J_{E}\left(\tilde{f}_{t}\right)-J_{E}(\tilde{f})}{t}=\int_{\mathbb{R}^{n}} \tilde{\psi} d \mu(\tilde{f})=e^{c} \int_{\mathbb{R}^{n}} \psi d \mu(f)+d e^{c} J_{E}(f) \tag{4.6}
\end{equation*}
$$

and

$$
\begin{equation*}
\lim _{t \rightarrow 0^{+}} \frac{J_{E_{t} \backslash E}\left(\tilde{f_{t}}\right)}{t}=\int_{S^{n-1}} h_{L} d \sigma(\tilde{f})=e^{c} \int_{S^{n-1}} h_{L} d \sigma(f) \tag{4.7}
\end{equation*}
$$

Now, since

$$
f \oplus t \cdot g=e^{-(c+d t)}(\tilde{f} \oplus t \cdot \tilde{g})
$$

we may compute the left hand sides of (4.4) and (4.5) as derivatives of a product.

Using (4.6), we get

$$
\begin{aligned}
\lim _{t \rightarrow 0^{+}} \frac{J_{E}\left(f_{t}\right)-J_{E}(f)}{t} & =-d e^{-c} J_{E}(\tilde{f})+e^{-c}\left[e^{c} \int_{\mathbb{R}^{n}} \psi d \mu(f)+d e^{c} J_{E}(f)\right] \\
& =\int_{\mathbb{R}^{n}} \psi d \mu(f)
\end{aligned}
$$

Similarly, using (4.7), we get

$$
\lim _{t \rightarrow 0^{+}} \frac{J_{E_{t} \backslash E}\left(f_{t}\right)}{t}=e^{-c} \cdot e^{c} \int_{S^{n-1}} h_{L} d \sigma(f)=\int_{S^{n-1}} h_{L} d \sigma(f) .
$$

Thus, in the remaining of the proof, we assume that all the following conditions hold true: $0 \in \operatorname{int}(F), v(0)=0, v \geq 0, \varphi \geq 0, \psi \geq 0$.

Step 3. For every $t>0$, it holds

$$
\begin{equation*}
J_{E}\left(f_{t}\right)-J_{E}(f)=\int_{0}^{t} \Psi(s) d s \tag{4.8}
\end{equation*}
$$

where

$$
\begin{equation*}
\Psi(s):=\int_{E} \psi d \mu\left(f_{s}\right) \quad \forall s \geq 0 \tag{4.9}
\end{equation*}
$$

Let $t>0$ be fixed, and take $C \subset \subset E$. Thanks to the reduction $0 \in \operatorname{int}(F)$ made in Step 2, we have $C \subset \subset E_{t}$. Then by Lemma 4.11 it holds

$$
\begin{equation*}
\lim _{h \rightarrow 0} \frac{f_{t+h}(x)-f_{t}(x)}{h}=\psi\left(\nabla u_{t}(x)\right) f_{t}(x) \quad \forall x \in C . \tag{4.10}
\end{equation*}
$$

Moreover, thanks to the reduction $v(0)=0$ made in Step 2, we can apply Lemmas 3.9 and 4.10 (ii) to infer that, for every $s \in[0,1]$, the nonnegative functions $\psi\left(\nabla u_{s}(x)\right) f_{s}(x)$ are bounded above on $C$ by some continuous function independent of $s$. Then, by the pointwise convergence in (4.10), Lagrange mean value theorem, and dominated convergence we infer

$$
\lim _{h \rightarrow 0} \frac{J_{C}\left(f_{t+h}\right)-J_{C}\left(f_{t}\right)}{h}=\lim _{h \rightarrow 0} \int_{C} \frac{f_{t+h}-f_{t}}{h} d x=\int_{C} \psi\left(\nabla u_{t}\right) f_{t} d x .
$$

So we have

$$
J_{C}\left(f_{t}\right)-J_{C}(f)=\int_{0}^{t}\left\{\int_{C} \psi d \mu\left(f_{s}\right)\right\} d s
$$

which implies (4.8) by letting $C \uparrow E$.
Step 4. The function $\Psi$ defined in (4.9) takes finite values at every $s \geq 0$.
Let $s>0$. By the reduction $\varphi \geq 0$ made in Step 2, we have

$$
s \Psi(s) \leq \int_{\mathbb{R}^{n}}(\varphi+s \psi) d \mu\left(f_{s}\right)=\int_{\mathbb{R}^{n}} u_{s}^{*}\left(\nabla u_{s}\right) f_{s} d x<+\infty
$$

where the last inequality follows from Lemma 4.12 (which applies thanks to the conditions $\varphi, \psi \geq 0$ ).

Let now $s=0$. Since by assumption $g$ is an admissible perturbation for $f$, by (4.1) it holds

$$
(\varphi-c \psi)(y) \geq(\varphi-c \psi)(0)+\langle y, \nabla \varphi(0)-c \nabla \psi(0)\rangle,
$$

so that

$$
\psi(y) \leq c_{1}+c_{2} \varphi(y)+c_{3}\|y\|,
$$

with

$$
c_{1}:=\psi(0)-c^{-1} \varphi(0), \quad c_{2}:=c^{-1}, \quad c_{3}:=c^{-1}\|\nabla \varphi(0)-c \nabla \psi(0)\| .
$$

Therefore

$$
\begin{aligned}
\int_{\mathbb{R}^{n}} \psi(\nabla u(x)) f(x) d x \leq & c_{1} \int_{\mathbb{R}^{n}} f(x) d x+c_{2} \int_{\mathbb{R}^{n}} \varphi(\nabla u(x)) f(x) d x \\
& +c_{3} \int_{\mathbb{R}^{n}}\|\nabla u(x)\| f(x) d x \\
= & c_{1} I_{1}+c_{2} I_{2}+c_{3} I_{3} .
\end{aligned}
$$

Let us show separately that each of the integrals $I_{j}, j=1,2,3$, is finite. As already noticed in Remark 3.2(i), the integral $I_{1}$ is finite for every $f \in \mathcal{A}$. The integral $I_{2}$ is finite by Lemma 4.12. Finally, in order to estimate the integral $I_{3}$, we use the coarea formula: if $m:=\max _{\mathbb{R}^{n}} f$ it holds

$$
\begin{equation*}
I_{3}=\int_{\mathbb{R}^{n}}\|\nabla f\| d x=\int_{0}^{m} \mathcal{H}^{n-1}(\partial\{f \geq s\}) d s \tag{4.11}
\end{equation*}
$$

According to Lemma 2.5, there exist constant $a, b$, with $a>0$ such that

$$
f(x) \leq g(x):=e^{-a\|x\|-b},
$$

which implies $\{f \geq s\} \subseteq\{g \geq s\}$, and in turn,

$$
\begin{equation*}
\mathcal{H}^{n-1}(\partial\{f \geq s\}) \leq \mathcal{H}^{n-1}(\partial\{g \geq s\})=c(n)\left(\frac{-\log s-b}{a}\right)^{n-1} . \tag{4.12}
\end{equation*}
$$

The finiteness of $I_{3}$ follows from (4.11) and (4.12).
Step 5. The function $\Psi$ defined in (4.9) is continuous at every $s>0$, and it is continuous from the right at $s=0$.

Through the change of variable $\nabla u_{s}(x)=y$, we obtain

$$
\Psi(s)=\int_{E} \psi\left(\nabla u_{s}(x)\right) f_{s}(x) d x=\int_{\mathbb{R}^{n}} h(s, y) d y
$$

with

$$
h(s, y):=\psi(y) e^{\varphi_{s}(y)-\left\langle y, \nabla \varphi_{s}(y)\right\rangle} \operatorname{det}\left(\nabla^{2} \varphi_{s}\right)(y) \chi_{Q_{s}}(y), \quad Q_{s}:=\nabla u_{s}(E) .
$$

We now use the expansion

$$
\operatorname{det}\left(\nabla^{2} \varphi_{s}\right)=\operatorname{det}\left(\nabla^{2} \varphi+s \nabla^{2} \psi\right)=\sum_{j=0}^{n} s^{j} D_{j}(\varphi, \psi),
$$

where the mixed determinants $D_{i}(\varphi, \psi)$ are nonnegative functions of $y$ independent of $s$. We infer that

$$
\begin{equation*}
\Psi(s)=\sum_{j=0}^{n} s^{j} \Psi_{j}(s), \tag{4.13}
\end{equation*}
$$

where

$$
\Psi_{j}(s):=\int_{\mathbb{R}^{n}} h_{j}(s, y) d y \quad h_{j}(s, y):=\psi(y) e^{\varphi_{s}(y)-\left\langle y, \nabla \varphi_{s}(y)\right\rangle} D_{j}(\varphi, \psi)(y) \chi_{Q_{s}}(y)
$$

In the sequel, when no ambiguity may arise, for brevity we omit to indicate the variable $y$ in the expressions $\psi(y), \varphi_{s}(y), \nabla \varphi_{s}(y), D_{j}(\varphi, \psi)(y)$, and $\chi Q_{s}(y)$.

Let us prove the continuity of $\Psi$ at a fixed $s_{0}>0$. In view of (4.13) it is enough to show that, for any fixed index $i \in\{0,1, \ldots, n\}$, the function $\Psi_{i}$ is continuous at $s_{0}$.

We begin by noticing that

$$
\begin{equation*}
\lim _{s \rightarrow s_{0}} \chi_{Q_{s}}(y)=\lim _{s \rightarrow s_{0}} \chi_{Q_{s_{0}}}(y) \quad \text { for a.e. } y \in \mathbb{R}^{n} . \tag{4.14}
\end{equation*}
$$

Indeed, when $E=F=\mathbb{R}^{n}$, (4.14) is trivially true since $Q_{s}=\mathbb{R}^{n}$ for every $s \geq 0$. Assume $E=K$ and $F=L$, with $K, L \in \mathcal{K}^{n}$. The reduction $0 \in \operatorname{int}(F)$ made in Step 2 ensures that $K \subset \subset K_{s_{0}}$, and hence by Lemma 4.10(ii), we know that $\nabla u_{s}$ converge uniformly to $\nabla u_{s_{0}}$ on $K$. Therefore, the compact sets $Q_{s}$ converge to $Q_{s_{0}}$ in Hausdorff distance, which implies that the characteristic functions $\chi Q_{s}$ converge to $\chi Q_{s_{0}}$ in $L^{1}\left(\mathbb{R}^{n}\right)$, which in turn implies (4.14).

Using (4.14), we deduce that we have the pointwise convergence

$$
\lim _{s \rightarrow s_{0}} h_{i}(s, y)=h_{i}\left(s_{0}, y\right) \quad \text { for a.e. } y \in \mathbb{R}^{n} .
$$

We claim that, as a consequence, $\Psi_{i}(s)$ tends to $\Psi_{i}\left(s_{0}\right)$ as $s \rightarrow s_{0}$ by dominated convergence. Indeed, let us show that $h_{i}(s, y)$ are bounded from above by a function in $L^{1}\left(\mathbb{R}^{n}\right)$ independent of $s$. By the reduction $v \geq 0$ made in Step 2 and by Lemma 3.9, for any fixed $y \in \mathbb{R}^{n}$ the map

$$
\begin{equation*}
s \mapsto e^{\varphi_{s}(y)-\left\langle y, \nabla \varphi_{s}(y)\right\rangle} \tag{4.15}
\end{equation*}
$$

is pointwise decreasing. Therefore, if we fix $\bar{s} \in\left(0, s_{0}\right)$, for any $s \geq \bar{s}$ it holds

$$
\begin{aligned}
h_{i}(s, y) & \leq \psi e^{\varphi_{\bar{s}}-\left\langle y, \nabla \varphi_{\bar{s}}\right\rangle} D_{i}(\varphi, \psi) \\
& =\frac{1}{\bar{s}^{i}} \psi e^{\varphi_{\bar{s}}-\left\langle y, \nabla \varphi_{\bar{s}}\right\rangle} \bar{s}^{i} D_{i}(\varphi, \psi) \\
& \leq \frac{1}{\overline{\bar{s}}^{i}} \psi e^{\varphi_{\bar{s}}-\left\langle y, \nabla \varphi_{\bar{s}}\right\rangle} \sum_{j=0}^{n} \bar{s}^{j} D_{j}(\varphi, \psi) \\
& =\frac{1}{\bar{s}^{i}} \psi e^{\varphi_{\bar{s}}-\left\langle y, \nabla \varphi_{\bar{s}}\right\rangle} \operatorname{det}\left(\nabla^{2} \varphi_{\bar{s}}\right) \\
& \leq \frac{1}{\bar{s}^{i+1}} \varphi_{\bar{s}} e^{\varphi_{\bar{s}}-\left\langle y, \nabla \varphi_{\bar{s}}\right\rangle} \operatorname{det}\left(\nabla^{2} \varphi_{\bar{s}}\right),
\end{aligned}
$$

and the function in the last line belongs to $L^{1}\left(\mathbb{R}^{n}\right)$ by Lemma 4.12.
Let us now prove the continuity from the right of $\Psi$ at $s=0$. To that aim, in view of (4.13) it is enough to show that

$$
\begin{align*}
& \lim _{s \rightarrow 0^{+}} \Psi_{0}(s)=\Psi(0)  \tag{4.16}\\
& \limsup _{s \rightarrow 0^{+}} \Psi_{i}(s)<+\infty \quad \forall i \in\{1, \ldots, n\} \tag{4.17}
\end{align*}
$$

To prove equality (4.16), we begin by noticing that, as $s \rightarrow 0^{+}$, the sets $Q_{s}$ invade $\mathbb{R}^{n}$, meaning

$$
\begin{equation*}
\forall r>0, \quad \exists \bar{s}>0: Q_{s} \supseteq B_{r} \forall s \leq \bar{s} \tag{4.18}
\end{equation*}
$$

Indeed, when $E=F=\mathbb{R}^{n}$, (4.18) is trivially true since $Q_{s}=\mathbb{R}^{n}$ for every $s \geq 0$. Assume $E=K$ and $F=L$, with $K, L \in \mathcal{K}^{n}$, and let $r>0$ be fixed. We have

$$
\begin{equation*}
Q_{s}=\nabla u_{s}(K) \supseteq \nabla u_{s}(C), \quad \text { with } C:=\nabla u^{-1}\left(B_{2 r}\right) \tag{4.19}
\end{equation*}
$$

Since $C \subset \subset K$ and $K \subset \subset K_{s}$ (the latter thanks to the reduction $0 \in \operatorname{int}(L)$ made in Step 2), by Lemma 4.10(ii) we know that $\nabla u_{s}$ converge uniformly to $\nabla u$ on $C$. Therefore, the compact sets $\nabla u_{s}(C)$ converge to $B_{2 r}$ in Hausdorff distance, so that they contain $B_{r}$ for $s$ sufficiently small. Combined with (4.19), this implies (4.18).

Using (4.18), we deduce that we have the pointwise convergence

$$
\begin{equation*}
\lim _{s \rightarrow 0} h_{0}(s, y)=h_{0}(0, y) \quad \text { for a.e. } y \in \mathbb{R}^{n} . \tag{4.20}
\end{equation*}
$$

Now, by the monotonicity of the map (4.15), for any $s \geq 0$ it holds

$$
\begin{equation*}
h_{0}(s, y) \leq h_{0}(0, y)=\psi e^{\varphi-\langle y, \nabla \varphi\rangle} \operatorname{det}\left(\nabla^{2} \varphi\right), \tag{4.21}
\end{equation*}
$$

and the last expression is in $L^{1}\left(\mathbb{R}^{n}\right)$ because we have proved in Step 4 that $\Psi(0)$ is finite.
In view of (4.20) and (4.21), (4.16) holds true by dominated convergence.
To prove (4.17) we notice that assumption (4.1) implies $\nabla^{2} \psi \leq c^{-1} \nabla^{2} \varphi$ and hence

$$
D_{i}(\varphi, \psi) \leq D_{i}\left(\varphi, c^{-1} \varphi\right)
$$

This, combined with the monotonicity of the map (4.15), implies

$$
h_{i}(s, y) \leq \psi e^{\varphi-\langle y, \nabla \varphi\rangle} D_{i}\left(\varphi, c^{-1} \varphi\right)=\psi e^{\varphi-\langle y, \nabla \varphi\rangle} \gamma_{i}(c) \operatorname{det}\left(\nabla^{2} \varphi\right),
$$

where the coefficients $\gamma_{i}(c)$ depend only on $c$. The last expression is in $L^{1}\left(\mathbb{R}^{n}\right)$ again by the finiteness of $\Psi(0)$, and (4.17) follows.

Step 6. Equality (4.4) holds.
The equality (4.8) proved in Step 3, together with the finiteness and continuity of $\Psi(s)$ for $s>0$ proved respectively in Steps 4 and 5, gives

$$
\begin{equation*}
\Psi(s)=\left.\frac{d}{d t} J_{E}\left(f_{t}\right)\right|_{t=s} \quad \forall s>0 \tag{4.22}
\end{equation*}
$$

Moreover, the continuity from the right of $\Psi$ at $s=0$ proved in Step 5 implies

$$
\begin{equation*}
\lim _{s \rightarrow 0^{+}} \Psi(s)=\Psi(0)=\int_{\mathbb{R}^{n}} \psi d \mu(f) \tag{4.23}
\end{equation*}
$$

Therefore,

$$
\begin{align*}
\lim _{t \rightarrow 0^{+}} \frac{J_{E}\left(f_{t}\right)-J_{E}(f)}{t} & =\left.\frac{d}{d t} J_{E}\left(f_{t}\right)\right|_{t=0^{+}}=\left.\lim _{s \rightarrow 0^{+}} \frac{d}{d t} J_{E}\left(f_{t}\right)\right|_{t=s} \\
& =\lim _{s \rightarrow 0^{+}} \Psi(s)=\int_{\mathbb{R}^{n}} \psi d \mu(f) \tag{4.24}
\end{align*}
$$

Step 7. Under the assumptions of Theorem 4.6, equality (4.5) holds.
We define the map $m: S^{n-1} \times(0, t] \rightarrow K_{t} \backslash K$ by

$$
m(\xi, s):=v_{K_{s}}^{-1}(\xi)=v_{K}^{-1}(\xi)+s v_{L}^{-1}(\xi)
$$

By the area formula [13, Section 3.1.5], we have

$$
\begin{equation*}
\int_{K_{t} \backslash K} f_{t}=\int_{0}^{t} \int_{S^{n-1}} f_{t}(m(\xi, s))|\operatorname{det} \operatorname{Jm}(\xi, s)| d \mathcal{H}^{n-1}(\xi) d s \tag{4.25}
\end{equation*}
$$

Let $(\xi, s) \in S^{n-1} \times[0, t]$ be fixed and let us compute $|\operatorname{det} \operatorname{Jm}(\xi, s)|$. We choose an orthonormal basis $\left\{e_{1}, \ldots, e_{n}\right\}$ of the tangent space $\xi^{\perp} \times \mathbb{R}$ to $S^{n-1} \times[0, t]$ given by

$$
e_{i}=\left(v_{i}, 0\right) \quad i=1, \ldots, n-1, \quad e_{n}=(0, \ldots, 0,1),
$$

where $v_{i}$ are eigenvectors of the reverse Weingarten operator $\nabla v_{K_{s}}^{-1}(\xi)$. Then, denoting by $\rho_{i}(\xi, s)$ the corresponding eigenvalues (namely the principal radii of curvature of $\partial K_{s}$ at $\xi$ ), it holds

$$
\partial_{e_{i}} m(\xi, s)=\rho_{i}(\xi, s) e_{i} \quad i=1, \ldots, n-1, \quad \partial_{e_{n}} m(\xi, s)=v_{L}^{-1}(\xi) .
$$

Hence

$$
\begin{aligned}
|\operatorname{det} \operatorname{Jm}(\xi, s)| & =\left\|\partial_{e_{1}} m(\xi, s) \wedge \cdots \wedge \partial_{e_{n}} m(\xi, s)\right\| \\
& =\left|\left\langle\xi, v_{L}^{-1}(\xi)\right\rangle\right| \cdot \prod_{i=1}^{n-1} \rho_{i}(\xi, s)=h_{L}(\xi) \operatorname{det}\left(\nabla v_{K_{s}}^{-1}(\xi)\right),
\end{aligned}
$$

where the last equality holds because, by the reduction $0 \in \operatorname{int}(L)$ made in Step 2, we have $h_{L} \geq 0$.

Now we recall that the reverse Weingarten operator of $K_{S}$ is given by

$$
\nabla v_{K_{s}}^{-1}=\left(h_{K_{s}}\right)_{i j}+h_{K_{s}} \delta_{i j},
$$

where indices $i$ and $j$ denote second order covariant derivation with respect to an orthonormal frame on $S^{n-1}$. Therefore, as $h_{K_{s}}=h_{K}+s h_{L}$, we have

$$
\nabla v_{K_{s}}^{-1}=\left(h_{K}\right)_{i j}+h_{K} \delta_{i j}+s\left[\left(h_{L}\right)_{i j}+h_{L} \delta_{i j}\right],
$$

and hence

$$
\begin{equation*}
|\operatorname{det} \operatorname{Jm}(\xi, s)|=h_{L}(\xi)\left[\operatorname{det}\left(\nabla v_{K}^{-1}(\xi)\right)+\sum_{i=1}^{n-1} \gamma_{i}(\xi) s^{i}\right] \tag{4.26}
\end{equation*}
$$

where $\gamma_{i}(\xi)$ are continuous functions depending on the principal curvatures and on the principal directions of $\partial K$ and $\partial L$ at $\xi$.

Inserting (4.26) into (4.25) and dividing by $t$ we obtain

$$
\begin{align*}
\frac{1}{t} \int_{K_{t} \backslash K} f_{t} d x= & \frac{1}{t} \int_{0}^{t} \int_{S^{n-1}} f_{t}(m(\xi, s)) h_{L}(\xi) \operatorname{det}\left(\nabla v_{K}^{-1}(\xi)\right) d \mathcal{H}^{n-1}(\xi) d s \\
& +\sum_{i=1}^{n-1} \frac{1}{t} \int_{0}^{t} s^{i}\left\{\int_{S^{n-1}} f_{t}(m(\xi, s)) h_{L}(\xi) \gamma_{i}(\xi) d \mathcal{H}^{n-1}(\xi)\right\} d s \tag{4.27}
\end{align*}
$$

We observe that

$$
\begin{equation*}
\lim _{t \rightarrow 0^{+}} \sum_{i=1}^{n-1} \frac{1}{t} \int_{0}^{t} s^{i}\left\{\int_{S^{n-1}} f_{t}(m(\xi, s)) h_{L}(\xi) \gamma_{i}(\xi) d \mathcal{H}^{n-1}(\xi)\right\}=0 \tag{4.28}
\end{equation*}
$$

Indeed, for every $i=1, \ldots, n-1$, we have

$$
\begin{aligned}
& \int_{0}^{t} s^{i}\left\{\int_{S^{n-1}} f_{t}(m(\xi, s)) h_{L}(\xi) \gamma_{i}(\xi) d \mathcal{H}^{n-1}(\xi)\right\} \\
& \leq\left(\sup _{\mathbb{R}^{n}} f_{1}\right) \int_{S^{n-1}} h_{L} \gamma_{i} d \mathcal{H}^{n-1} \int_{0}^{t} s^{i} d s,
\end{aligned}
$$

where we used the inequality $f_{t}(x) \leq f_{1}(x)$ holding for every $x \in \mathbb{R}^{n}$ and every $t \in[0,1]$ by Lemma 3.9 (which applies thanks to the reduction $v(0)=0$ made in Step 2).

By (4.27) and (4.28), to conclude the proof of Step 7 it is enough to show that

$$
\lim _{t \rightarrow 0^{+}} \frac{1}{t} \int_{0}^{t} \int_{S^{n-1}} f_{t}(m(\xi, s)) h_{L}(\xi) \operatorname{det}\left(\nabla v_{K}^{-1}(\xi)\right) d \mathcal{H}^{n-1}(\xi) d s=\int_{S^{n-1}} h_{L} d \sigma(f)
$$

or equivalently

$$
\lim _{t \rightarrow 0^{+}} \frac{1}{t} \int_{0}^{t} \int_{S^{n-1}}\left[f_{t}(m(\xi, s))-f(m(\xi, 0))\right] h_{L}(\xi) \operatorname{det}\left(\nabla v_{K}^{-1}(\xi)\right) d \mathcal{H}^{n-1}(\xi)=0
$$

Such equality is clearly satisfied if

$$
\begin{equation*}
\lim _{t \rightarrow 0^{+}} \sup _{s \in[0, t], \xi \in S^{n-1}}\left|f_{t}(m(\xi, s))-f(m(\xi, 0))\right|=0 \tag{4.29}
\end{equation*}
$$

Let $s \in[0, t]$ and $\xi \in S^{n-1}$. By Lemma 4.13 applied at the point $x:=m(\xi, s) \in \partial K_{s} \subset K_{t}$, there exists $y \in K \cap(x-t L)$ such that

$$
t v_{\min }+u(y) \leq u_{t}(m(\xi, s)) \leq t v_{\max }+u(y)
$$

Hence

$$
\begin{align*}
t v_{\min }+u(y)-u(m(\xi, 0)) & \leq u_{t}(m(\xi, s))-u(m(\xi, 0)) \\
& \leq t v_{\max }+u(y)-u(m(\xi, 0)) . \tag{4.30}
\end{align*}
$$

As $x \in m(\xi, 0)+s L \subseteq m(\xi, 0)+t L$, we have $m(\xi, 0) \in K \cap(x-t L)$, and therefore

$$
\begin{equation*}
\|m(\xi, 0)-y\| \leq \operatorname{diam}(K \cap(x-t L)) \leq t \operatorname{diam}(L) \tag{4.31}
\end{equation*}
$$

By (4.30), (4.31) and the uniform continuity of $u$ on $K$, we infer that

$$
\lim _{t \rightarrow 0^{+}} \sup _{s \in[0, t], \xi \in S^{n-1}}\left|u_{t}(m(\xi, s))-u(m(\xi, 0))\right|=0
$$

and (4.29) follows.
Step 8: Conclusion.
Equalities (4.2) and (4.3) follow from Steps 1, 6, and 7. Moreover, the finiteness of $\Psi(0)$ proved in Step 4 implies that $\int_{\mathbb{R}^{n}} \psi d \mu(f)<+\infty$; on the other hand, for any $K, L \in \mathcal{K}^{n}$, one has $\int_{S^{n-1}} h_{L} d \sigma(f)<+\infty$. Therefore $\delta J(f, g)$ is finite.

## 5. A functional form of Minkowski's first inequality

Minkowski's first inequality states that

$$
\begin{equation*}
\lim _{t \rightarrow 0^{+}} \frac{V(K+t L)-V(K)}{t}=n V_{1}(K, L) \geq n V(K)^{1-\frac{1}{n}} V(L)^{\frac{1}{n}} \quad \forall K, L \in \mathcal{K}_{0}^{n} \tag{5.1}
\end{equation*}
$$

with equality sign if and only if $K$ and $L$ are homothetic (see [31, Theorem 6.2.1]).
The main result of this section provides a functional version of such inequality:
Theorem 5.1. Let $f, g \in \mathcal{A}$, and assume that $J(f)>0$. Then

$$
\begin{equation*}
\delta J(f, g) \geq J(f)[\log J(g)+n]+\operatorname{Ent}(f) \tag{5.2}
\end{equation*}
$$

with equality sign if and only if there exists $x_{0} \in \mathbb{R}^{n}$ such that $g(x)=f\left(x-x_{0}\right) \forall x \in \mathbb{R}^{n}$.

Remark 5.2. We point out that, by choosing $f=\gamma_{n}$, Theorem 5.1 allows to recover the Urysohn-type inequality for the mean width of a log-concave function proved in [17, Proposition 3.2] and [30, Theorem 1.4].

Before giving the proof of Theorem 5.1, let us present a straightforward consequence of it, which will be exploited in Section 7 in order to get uniqueness in the functional form of the Minkowski's problem.

Corollary 5.3. Let $f_{1}, f_{2} \in \mathcal{A}$, with $J\left(f_{1}\right)=J\left(f_{2}\right)>0$, and assume that

$$
\begin{equation*}
\delta J\left(f_{2}, f_{1}\right)=\delta J\left(f_{1}, f_{1}\right) \quad \text { and } \quad \delta J\left(f_{1}, f_{2}\right)=\delta J\left(f_{2}, f_{2}\right) \tag{5.3}
\end{equation*}
$$

Then there exists $x_{0} \in \mathbb{R}^{n}$ such that $f_{2}(x)=f_{1}\left(x-x_{0}\right) \forall x \in \mathbb{R}^{n}$.
Proof. By the assumption $J\left(f_{i}\right)>0$, we may apply inequality (5.2) (once with $f=f_{1}$ and $g=f_{2}$ and once with $f=f_{2}$ and $\left.g=f_{1}\right)$; since $J\left(f_{1}\right)=J\left(f_{2}\right)$, we get

$$
\begin{align*}
& \delta J\left(f_{1}, f_{2}\right) \geq n J\left(f_{1}\right)+\int_{\mathbb{R}^{n}} f_{1} \log f_{1} d x \text { and }  \tag{5.4}\\
& \delta J\left(f_{2}, f_{1}\right) \geq n J\left(f_{2}\right)+\int_{\mathbb{R}^{n}} f_{2} \log f_{2} d x .
\end{align*}
$$

By assumption (5.3) and Proposition 3.11, the two inequalities in (5.4) may be rewritten respectively as

$$
\delta J\left(f_{2}, f_{2}\right) \geq \delta J\left(f_{1}, f_{1}\right) \quad \text { and } \quad \delta J\left(f_{1}, f_{1}\right) \geq \delta J\left(f_{2}, f_{2}\right),
$$

which implies that both hold with equality sign. Then $f_{1}$ and $f_{2}$ are translates of each other by Theorem 5.1.

We now turn to the proof of Theorem 5.1. We need the following
Lemma 5.4. Let $f, g \in \mathcal{A}$, and assume that $J(f)>0$. Then

$$
\lim _{t \rightarrow 0^{+}} \frac{J((1-t) \cdot f \oplus t \cdot g)-J(f)}{t}=\delta J(f, g)-\delta J(f, f) .
$$

Proof. For $t \in(0,1)$, we set

$$
\alpha(t):=\frac{t}{1-t} \quad \text { and } \quad f_{\alpha(t)}:=f \oplus \alpha(t) \cdot g
$$

Let us write

$$
\begin{equation*}
\frac{J((1-t) \cdot f \oplus t \cdot g)-J(f)}{t}=\frac{J\left((1-t) \cdot f_{\alpha(t)}\right)-J\left(f_{\alpha(t)}\right)}{t}+\frac{J\left(f_{\alpha(t)}\right)-J(f)}{t} \tag{5.5}
\end{equation*}
$$

and let us focus attention on the first term in the r.h.s. of (5.5).
For every fixed $t \in(0,1)$, we have

$$
\frac{J\left((1-t) \cdot f_{\alpha(t)}\right)-J\left(f_{\alpha(t)}\right)}{t}=\frac{\gamma_{t}(t)-\gamma_{t}(0)}{t}
$$

where the function $\gamma_{t}$ is defined by

$$
\gamma_{t}(s):=J\left((1-s) \cdot f_{\alpha(t)}\right) \quad \forall s \in(0,1) .
$$

In more explicit terms

$$
\gamma_{t}(s)=\int_{\mathbb{R}^{n}} f_{\alpha(t)}^{1-s}\left(\frac{x}{1-s}\right) d x=(1-s)^{n} \int_{\mathbb{R}^{n}} f_{\alpha(t)}^{1-s}(y) d y .
$$

Proceeding as in the proof of Proposition 3.11 we can differentiate this expression with respect to $s$ and obtain:

$$
\gamma_{t}^{\prime}(s)=-n(1-s)^{n-1} J\left(f_{\alpha(t)}^{1-s}\right)-(1-s)^{n} \int_{\mathbb{R}^{n}} f_{\alpha(t)}^{1-s} \log \left(f_{\alpha(t)}^{1-s}\right) d x
$$

Then, for every fixed $t \in(0,1)$, we can apply Lagrange mean value theorem to infer that there exists $\bar{s} \in(0, t)$ such that

$$
\begin{align*}
\frac{J\left((1-t) \cdot f_{\alpha(t)}\right)-J\left(f_{\alpha(t)}\right)}{t}= & \gamma_{t}^{\prime}(\bar{s})=-n(1-\bar{s})^{n-1} J\left(f_{\alpha(t)}^{1-\bar{s}}\right) \\
& -(1-\bar{s})^{n} \int_{\mathbb{R}^{n}} f_{\alpha(t)}^{1-\bar{s}} \log \left(f_{\alpha(t)}^{1-\bar{s}}\right) d x \tag{5.6}
\end{align*}
$$

We are now ready to pass to the limit as $t \rightarrow 0^{+}$in the r.h.s. of (5.5).
Concerning the first term, assume for a moment that the function $v:=-\log g$ satisfies the condition $v(0)=0$. In this case, by Lemma 3.9, as $t \rightarrow 0^{+}$the functions $f_{\alpha(t)}(x)$ converge decreasingly to some pointwise limit $\tilde{f}(x)$ (which is bounded above and below by some functions in $\mathcal{A}$ ). Then, by monotone convergence, taking also into account that $\bar{s} \rightarrow 0^{+}$as $t \rightarrow 0^{+}$, we infer from (5.6) that

$$
\begin{equation*}
\lim _{t \rightarrow 0^{+}} \frac{J\left((1-t) \cdot f_{\alpha(t)}\right)-J\left(f_{\alpha(t)}\right)}{t}=-n J(\tilde{f})-\int_{\mathbb{R}^{n}} \tilde{f} \log \tilde{f} d x \in(-\infty,+\infty) \tag{5.7}
\end{equation*}
$$

Concerning the second term, differentiating a composition of functions shows immediately that

$$
\begin{equation*}
\lim _{t \rightarrow 0^{+}} \frac{J\left(f_{\alpha(t)}\right)-J(f)}{t}=\delta J(f, g) \tag{5.8}
\end{equation*}
$$

By combining (5.7) and (5.8), it is straightforward to conclude. Indeed, similarly as in the proof of Theorem 3.6, we may distinguish the two cases $J(\tilde{f})>J(f)$ and $J(\tilde{f})=J(f)$.

If $J(\tilde{f})>J(f)$, the limit in (5.7) remains finite, whereas the limit in (5.8) becomes $+\infty$. Hence it holds

$$
\lim _{t \rightarrow 0^{+}} \frac{J((1-t) \cdot f \oplus t \cdot g)-J(f)}{t}=\delta J(f, g)=+\infty
$$

and the statement of the lemma holds true.
If $J(\tilde{f})=J(f)$, then $\tilde{f}=f \mathcal{H}^{n}$-a.e., so that the r.h.s. of (5.7) agrees with $-\delta J(f, f)$, and the lemma follows summing up (5.7) and (5.8).

It remains to get rid of the assumption $v(0)=0$. In the general case, we set as usual

$$
d:=v(0), \quad \tilde{v}(x):=v(x)-d, \quad \tilde{g}(x):=e^{-\tilde{v}(x)}
$$

Since

$$
(1-t) \cdot f \oplus t \cdot g=e^{-d t}((1-t) \cdot f \oplus t \cdot \tilde{g})
$$

we have

$$
\begin{aligned}
\frac{J((1-t) \cdot f \oplus t \cdot g)-J(f)}{t}= & J(f) \frac{e^{-d t}-1}{t} \\
& +e^{-d t} \frac{J((1-t) \cdot f \oplus t \cdot \tilde{g})-J(f)}{t}
\end{aligned}
$$

By passing to the limit as $t \rightarrow 0^{+}$, since $\tilde{v}(0) \geq 0$ by construction, we obtain

$$
\lim _{t \rightarrow 0^{+}} \frac{J((1-t) \cdot f \oplus t \cdot g)-J(f)}{t} \geq-d J(f)+\delta J(f, \tilde{g})-\delta J(f, f)
$$

To conclude, it is enough to observe that $-d J(f)+\delta J(f, \tilde{g})=\delta J(f, g)(c f .(3.5))$.
Proof of Theorem 5.1. By the Prékopa-Leindler inequality, the function $\psi(t):=\log (J((1-$ $t) \cdot f \oplus t \cdot g))$ is concave on $[0,1](c f$. Remark 3.3). In particular, it holds

$$
\begin{equation*}
\psi(t) \geq \psi(0)+t[\psi(1)-\psi(0)] \quad \forall t \in[0,1] . \tag{5.9}
\end{equation*}
$$

As a consequence, the (right) derivative of the function $\psi$ at $t=0$ satisfies

$$
\begin{equation*}
\psi^{\prime}(0) \geq[\psi(1)-\psi(0)] . \tag{5.10}
\end{equation*}
$$

By Lemma 5.4, we have

$$
\psi^{\prime}(0)=\frac{\delta J(f, g)-\delta J(f, f)}{J(f)}
$$

Therefore (5.10) can be rewritten as

$$
\frac{\delta J(f, g)-\delta J(f, f)}{J(f)} \geq \log \left(\frac{J(g)}{J(f)}\right)
$$

Inserting (3.12) into the above inequality, (5.2) is proved.
Finally, assume that $g(x)=f\left(x-x_{0}\right)$ for some $x_{0} \in \mathbb{R}^{n}$. Then (5.2) holds with equality sign thanks to Proposition 3.11 and the invariance of $J$ by translation of coordinates. Conversely, assume that (5.2) holds with equality sign. By inspection of the above proof one sees immediately that also inequality (5.10), and hence inequality (5.9), must hold with equality sign. This entails that the Prékopa-Leindler inequality holds as an equality, and therefore $f$ and $g$ agree up to a translation.

## 6. Isoperimetric and log-Sobolev inequalities for log-concave functions

Let us now turn attention to some consequences of the results in Sections 4 and 5.
Motivated by the equality

$$
\lim _{t \rightarrow 0^{+}} \frac{V\left(K+t B_{1}\right)-V(K)}{t}=P(K)
$$

and having in mind that the Gaussian probability density

$$
\gamma_{n}(x):=c_{n} e^{-\frac{\|x\|^{2}}{2}}, \quad c_{n}:=(2 \pi)^{-\frac{n}{2}},
$$

plays within the class $\mathcal{A}^{\prime}$ the role of the unit ball in $\mathcal{K}^{n}$, we set the following

Definition 6.1. For any $f \in \mathcal{A}^{\prime}$ with $J(f)>0$, we define the perimeter of $f$ as

$$
P(f):=\delta J\left(f, \gamma_{n}\right) .
$$

Similarly as Minkowski’s first inequality (5.1) (when applied with $L$ equal to a ball $B$ ) implies the classical isoperimetric inequality

$$
V(K)^{\frac{1}{n}} P(K)^{-\frac{1}{n-1}} \leq V(B)^{\frac{1}{n}} P(B)^{-\frac{1}{n-1}} \quad \forall K \in \mathcal{K}_{0}^{n}
$$

Theorem 5.1 (when applied with $g=\gamma_{n}$ and combined with Theorem 4.5) yields the following functional version of the isoperimetric inequality:

Proposition 6.2. Let $f=e^{-u} \in \mathcal{A}^{\prime}$ Then

$$
\begin{equation*}
P(f) \geq n J(f)+\operatorname{Ent}(f), \tag{6.1}
\end{equation*}
$$

with equality sign if and only if there exists $x_{0} \in \mathbb{R}^{n}$ such that $f(x)=\gamma_{n}\left(x-x_{0}\right) \forall x \in \mathbb{R}^{n}$. In particular, if $\varphi:=u^{*}$ is uniformly strictly convex, namely

$$
\begin{equation*}
\exists c>0: \nabla^{2} \varphi(y) \geq c \text { Id } \quad \forall y \in \mathbb{R}^{n} \tag{6.2}
\end{equation*}
$$

inequality (6.1) reads

$$
\begin{equation*}
\frac{1}{2} \int_{\mathbb{R}^{n}} \frac{\|\nabla f\|^{2}}{f} d x+\left(\log c_{n}\right) J(f) \geq n J(f)+\operatorname{Ent}(f) \tag{6.3}
\end{equation*}
$$

Proof. Inequality in (6.1) is obtained by applying Theorem 5.1 (simply take into account that $J\left(\gamma_{n}\right)=1$ ). If (6.2) holds, then $\gamma_{n}$ is an admissible perturbation for $f$ according to Definition 4.4. In this case, by applying Theorem 4.5, one gets

$$
P(f)=\delta J\left(f, \gamma_{n}\right)=\int_{\mathbb{R}^{n}}\left(\frac{1}{2}\|\nabla u\|^{2}+\log c_{n}\right) f d x
$$

and (6.3) follows.
For related functional versions of the isoperimetric inequality, we refer to [5, Section 6].
As a further application of our results, we now provide a generalized logarithmic Sobolev inequality for log-concave measures. After the pioneering result by Gross concerning the Gaussian measure [14], the validity of logarithmic Sobolev inequalities for more general probability measures, having in particular a log-concave density, has been investigated by several authors. We refer in particular to the paper [4] by Bobkov, where necessary and sufficient conditions are discussed.

Proposition 6.3. Let $v=g \mathcal{H}^{n}=e^{-v} \mathcal{H}^{n}$ be a log-concave probability measure such that $g \in \mathcal{A}^{\prime}$ and

$$
\begin{equation*}
\nabla^{2} v \geq c \text { Id } \quad \text { for some } c>0 \tag{6.4}
\end{equation*}
$$

Let $a: \mathbb{R}_{+} \rightarrow \mathbb{R}_{+}$be a $\mathcal{C}^{2}$ increasing function with $a(0)=0$.
Let $h$ be a positive function of class $\mathcal{C}^{2}\left(\mathbb{R}^{n}\right)$ which satisfies the conditions

$$
\begin{align*}
& \lim _{\|x\| \rightarrow+\infty} \frac{-\log (a(h))+v}{\|x\|}=+\infty \quad \text { and }  \tag{6.5}\\
& -c^{\prime} \nabla^{2} v \leq \nabla^{2} \log (a(h))<\nabla^{2} v \quad \text { for some } c^{\prime}>0
\end{align*}
$$

Then it holds

$$
\begin{align*}
& \int_{\mathbb{R}^{n}} a(h) \log a(h) d v-\left(\int_{\mathbb{R}^{n}} a(h) d \nu\right) \log \left(\int_{\mathbb{R}^{n}} a(h) d \nu\right) \\
& \quad \leq \frac{1}{c} \int_{\mathbb{R}^{n}} \frac{\left(a^{\prime}(h)\right)^{2}}{a(h)}\|\nabla h\|^{2} d \nu \tag{6.6}
\end{align*}
$$

Example 6.4. Assume that $g=\gamma_{n}$ is the density of the Gaussian measure and $a$ is defined by $a(t)=t^{p}$, for some $p \geq 1$. If $h$ is of the form $h=e^{-w}$, where $w \in C^{2}\left(\mathbb{R}^{n}\right)$ is such that $w(x) \rightarrow 0$ as $\|x\| \rightarrow \infty$ and there exist $c, C>0$ such that

$$
c \operatorname{Id} \leq D^{2} w(x) \leq C \text { Id } \quad \forall x \in \mathbb{R}^{n}
$$

then the assumptions of Proposition 6.3 are verified.
Remark 6.5. The constant $\frac{1}{c}$ in the r.h.s. of (6.6) is non-optimal. Indeed, consider for instance the case when $g=\gamma_{n}$ (so that $c=1$ ), and $a(t)=t^{2}$. Then (6.6) becomes

$$
\begin{equation*}
\int_{\mathbb{R}^{n}} h^{2} \log \left(h^{2}\right) d v-\left(\int_{\mathbb{R}^{n}} h^{2} d v\right) \log \left(\int_{\mathbb{R}^{n}} h^{2} d v\right) \leq 4 \int_{\mathbb{R}^{n}}\|\nabla h\|^{2} d v \tag{6.7}
\end{equation*}
$$

and it is known that (6.7) holds true with 2 in place of 4 at the r.h.s. This assertion can be recovered by inspection of the proof below, since in this case the number $t$ appearing in (6.11) equals $\frac{1}{2}$.

Remark 6.6. It is not surprising that, in order to have an inequality of logarithmic Sobolev type for the measure $v$, condition (6.4) is needed; indeed, (6.4) can be related to the so-called Herbst necessary condition (see [4] for a more detailed discussion).

Proof of Proposition 6.3. Set $f:=a(h) g$. Since $\int_{\mathbb{R}^{n}} g d x=1$, inequality (5.2) reads

$$
\begin{equation*}
\delta J(f, g) \geq n J(f)+\operatorname{Ent}(f) . \tag{6.8}
\end{equation*}
$$

The computation of $J(f)$ and $\int_{\mathbb{R}^{n}} f \log f d x$ is straightforward:

$$
J(f)=\int_{\mathbb{R}^{n}} a(h) d v, \quad \int_{\mathbb{R}^{n}} f \log f d x=\int_{\mathbb{R}^{n}}(-v+\log a(h)) a(h) d \nu
$$

On the other hand, by the hypotheses made on $h$ and $g$, the functions $f=a(h) g$ and $g$ turn out to satisfy the assumptions of Theorem 4.5. In particular, we point out that the upper and lower bound for $\nabla^{2} \log (a(h))$ in (6.5) ensure respectively that $f \in \mathcal{A}^{\prime}$ and that $g$ is an admissible perturbation for $f$. Then, setting $\psi=v^{*}$, we have

$$
\delta J(f, g)=\int_{\mathbb{R}^{n}} \psi(\nabla v-\nabla \log a(h)) a(h) d v .
$$

Inserting the above expressions of $J(f), \int_{\mathbb{R}^{n}} f \log f d x$, and $\delta J(f, g)$ into (6.8) leads to

$$
\begin{equation*}
\int_{\mathbb{R}^{n}} a(h) \log a(h) d v-\left(\int_{\mathbb{R}^{n}} a(h) d v\right) \log \left(\int_{\mathbb{R}^{n}} a(h) d v\right) \leq R(h) \tag{6.9}
\end{equation*}
$$

with

$$
R(h)=\int_{\mathbb{R}^{n}}[\psi(\nabla v-\nabla \log a(h))+v-n] a(h) d v
$$

Using the identity $v(x)=\langle x, \nabla v(x)\rangle-\psi(\nabla v(x))$, we may rewrite $R(h)$ as

$$
R(h)=\int_{\mathbb{R}^{n}}[\psi(\nabla v-\nabla \log a(h))-\psi(\nabla v)+\langle x, \nabla v\rangle-n] a(h) d v .
$$

Now we observe that

$$
\langle x, \nabla v\rangle a(h) g=-\langle x, \nabla g\rangle a(h)=-\operatorname{div}(x a(h) g)+\langle x, \nabla a(h)\rangle g+n a(h) g,
$$

and

$$
\int_{\mathbb{R}^{n}} \operatorname{div}(x a(h) g) d x=\lim _{r \rightarrow+\infty} r \int_{\partial B_{r}} a(h) g d \mathcal{H}^{n-1}=\lim _{r \rightarrow+\infty} r \int_{\partial B_{r}} f d \mathcal{H}^{n-1}=0
$$

(where the last equality is satisfied by the exponential decay of $f$ at infinity, $c f$. Lemma 2.5). Therefore,

$$
\begin{equation*}
R(h)=\int_{\mathbb{R}^{n}}[\psi(\nabla v-\nabla \log a(h))-\psi(\nabla v)+\langle x, \nabla \log a(h)\rangle] a(h) d v . \tag{6.10}
\end{equation*}
$$

In view of (6.9) and (6.10), the statement is proved if the following pointwise inequality holds:

$$
\psi(\nabla v-\nabla \log a(h))-\psi(\nabla v)+\langle x, \nabla \log a(h)\rangle \leq \frac{1}{c}\|\nabla \log a(h)\|^{2} .
$$

This is readily checked: indeed, setting $y:=-\nabla \log a(h)$, by Lagrange mean value theorem and assumption (6.4), there exist $t, s \in(0,1)$ such that

$$
\begin{align*}
\psi(\nabla v+y)-\psi(\nabla v)-\langle x, y\rangle & =\langle\nabla \psi(\nabla v+t y), y\rangle-\langle\nabla \psi(\nabla v), y\rangle \\
& =\left\langle\nabla^{2} \psi(\nabla v+\text { sty }) t y, y\right\rangle \leq \frac{1}{c}\|y\|^{2} \tag{6.11}
\end{align*}
$$

and the proof is achieved.

## 7. About the Minkowski's problem

In this concluding section we move the first steps towards the solution of the functional Minkowski's problem. In view of Theorems 4.5 and 4.6, its formulation within the class $\mathcal{A}^{\prime}$ or $\mathcal{A}^{\prime \prime}$ reads as follows: find $f \in \mathcal{A}^{\prime}$ such that

$$
\begin{equation*}
\mu(f)=m \tag{7.1}
\end{equation*}
$$

where $m$ is a given positive Borel measure on $\mathbb{R}^{n}$, or find $f \in \mathcal{A}^{\prime \prime}$ such that

$$
\begin{equation*}
(\mu(f), \sigma(f))=(m, \eta), \tag{7.2}
\end{equation*}
$$

where $(m, \eta)$ are given positive Borel measures respectively on $\mathbb{R}^{n}$ and $S^{n-1}$. Here the measures $\mu(f)$ and $\sigma(f)$ are intended according to Definition 4.1.

We begin by the following simple observation.

Remark 7.1. We have the following finiteness necessary condition on the measures $m$ and $\eta$, in order to solve the Minkowski's problem with datum $m$ or $(m, \eta)$ :

$$
\int_{\mathbb{R}^{n}} d m<+\infty, \quad \int_{S^{n-1}} d \eta<+\infty .
$$

Indeed, if $f$ belongs to $\mathcal{A}^{\prime}$ or to $\mathcal{A}^{\prime \prime}$, we have

$$
\int_{\mathbb{R}^{n}} d \mu(f)=J(f)<+\infty
$$

while, if $f \in \mathcal{A}^{\prime \prime}$ we have

$$
\int_{S^{n-1}} d \sigma(f) \leq\left(\max _{K} f\right) \mathcal{H}^{n-1}(\partial K)<+\infty
$$

where $K=\operatorname{dom}(-\log f)$.
Next, we show that, for the solvability of (7.1), $m$ must satisfy an equilibrium condition, which is completely analogous to the null barycentre property well-known in the classical Minkowski's problem for convex bodies. The same holds true, for the solvability of (7.2), replacing $m$ by the pair $(m, \eta)$.

Proposition 7.2. (i) For any $f \in \mathcal{A}^{\prime}$, the measure $\mu(f)$ verifies

$$
\int_{\mathbb{R}^{n}} y d \mu(f)(y)=0 .
$$

(ii) For any $f \in \mathcal{A}^{\prime \prime}$, the measures $\mu(f)$ and $\sigma(f)$ verify

$$
\int_{\mathbb{R}^{n}} y d \mu(f)(y)+\int_{S^{n-1}} y d \sigma(f)(y)=0 .
$$

Proof. Given a point $x_{0} \in \mathbb{R}^{n}$ and a function $v \in \mathcal{L}$, we denote by $[v]_{x_{0}}$ the translated function $x \mapsto v\left(x-x_{0}\right)$. With this notation it is straightforward to check that, for any $u, v \in \mathcal{L}$, it holds

$$
\begin{equation*}
u \square[v]_{x_{o}}=[u \square v]_{x_{0}} . \tag{7.3}
\end{equation*}
$$

Assume now that $f=e^{-u}$ belongs either to $\mathcal{A}^{\prime}$ or to $\mathcal{A}^{\prime \prime}$. For any fixed $x_{0} \in \mathbb{R}^{n}$ and any $\varepsilon>0$, let us compute $\delta J\left(f, g_{\varepsilon}\right)$, where $g_{\varepsilon}=e^{-v_{\varepsilon}}$, being

$$
v_{\varepsilon}(x):=\varepsilon u\left(\frac{x-x_{0}}{\varepsilon}\right)=[u \varepsilon]_{\frac{x_{0}}{\varepsilon}}(x) \quad \forall x \in \mathbb{R}^{n} .
$$

For any $t>0$ one has

$$
\left(v_{\varepsilon} t\right)(x)=t \varepsilon u\left(\frac{x-t x_{0}}{t \varepsilon}\right)=[u(t \varepsilon)]_{\frac{x_{0}}{\varepsilon}}(x),
$$

and hence, in view of (7.3),

$$
u \square\left(v_{\varepsilon} t\right)=[u \square u(t \varepsilon)]_{\frac{x_{0}}{\varepsilon}} .
$$

Therefore,

$$
\delta J\left(f, g_{\varepsilon}\right)=\lim _{t \rightarrow 0^{+}} \frac{J\left(e^{-u \square u(t \varepsilon)}\right)-J(f)}{t}
$$

$$
\begin{equation*}
=\varepsilon \lim _{t \rightarrow 0^{+}} \frac{J\left(e^{-u \square u(t \varepsilon)}\right)-J(f)}{t \varepsilon}=\varepsilon \delta J(f, f) . \tag{7.4}
\end{equation*}
$$

On the other hand, we observe that

$$
v_{\varepsilon}^{*}(y)=\left\langle x_{0}, y\right\rangle+\varepsilon u^{*}(y) \quad \text { and } \quad \operatorname{dom}\left(v_{\varepsilon}\right)=x_{0}+\varepsilon \operatorname{dom}(u) .
$$

Therefore, if $f \in \mathcal{A}^{\prime}$, by applying Theorem 4.5 we get

$$
\begin{equation*}
\delta J\left(f, g_{\varepsilon}\right)=\int_{\mathbb{R}^{n}}\left\langle x_{0}, y\right\rangle d \mu(f)(y)+\varepsilon \int_{\mathbb{R}^{n}} u^{*}(y) d \mu(f)(y) ; \tag{7.5}
\end{equation*}
$$

similarly, if $f \in \mathcal{A}^{\prime \prime}$, by applying Theorem 4.6 we get

$$
\begin{align*}
\delta J\left(f, g_{\varepsilon}\right)= & \int_{\mathbb{R}^{n}}\left\langle x_{0}, y\right\rangle d \mu(f)(y)+\varepsilon \int_{\mathbb{R}^{n}} u^{*}(y) d \mu(f)(y) \\
& +\int_{S^{n-1}}\left\langle x_{0}, y\right\rangle d \sigma(f)(y)+\varepsilon \int_{S^{n-1}} h_{\operatorname{dom}(u)}(y) d \sigma(f)(y) \tag{7.6}
\end{align*}
$$

We now observe that the following terms, which appear multiplied by $\varepsilon$ in (7.4)-(7.6), are finite:

$$
\delta J(f, f), \quad \int_{\mathbb{R}^{n}} u^{*}(y) d \mu(f)(y), \quad \int_{S^{n-1}} h_{\operatorname{dom}(u)}(y) d \sigma(f)(y)
$$

(recall in particular Proposition 3.11 and Lemma 4.12). Then the statement follows by combining (7.4) with (7.5) or (7.6), in the limit as $\varepsilon \rightarrow 0^{+}$.

Remark 7.3. We observe that the conditions expressed by Remark 7.1 and Proposition 7.2 are in general not sufficient for the solvability of the Minkowski's problem within one of the classes $\mathcal{A}^{\prime}$ or $\mathcal{A}^{\prime \prime}$. Indeed, assume for instance that $n=1$ and consider the Minkowski's problem in $\mathcal{A}^{\prime}$ : given an absolutely continuous measure on $\mathbb{R}$ with a positive continuous density $m$, satisfying the necessary conditions $\int_{\mathbb{R}} m(y) d y<+\infty$ and $\int_{\mathbb{R}} y m(y) d y=0$, it amounts to finding a function $\varphi \in \mathcal{C}_{+}^{2}(\mathbb{R})$, with $u=\varphi^{*} \in \mathcal{L}^{\prime}$, solving the second order o.d.e.

$$
\begin{equation*}
e^{\varphi(y)-y \varphi^{\prime}(y)} \varphi^{\prime \prime}(y)=m(y) \quad \forall y \in \mathbb{R} \tag{7.7}
\end{equation*}
$$

We observe that, if $\varphi$ is a solution to (7.7), for any $\alpha \in \mathbb{R}$, also $\varphi+\alpha y$ is a solution. Therefore, we may assume with no loss of generality that $\varphi^{\prime}(0)=0$, and write the unique solution to (7.7) with initial datum at $y=0$ as

$$
\begin{equation*}
\varphi(y)=\varphi(0)-y \int_{0}^{y} \frac{\log \left(e^{\varphi(0)}-M(t)\right)-\varphi(0)}{t^{2}} d t, \quad \text { where } M(t):=\int_{0}^{t} \operatorname{sm}(s) d s \tag{7.8}
\end{equation*}
$$

Now, let $u=\varphi^{*} \in \mathcal{L}^{\prime}$. In particular, this implies that $\operatorname{dom}\left(\varphi^{*}\right)=\mathbb{R}^{n}$. Since the condition of being cofinite is equivalent to the condition of being supercoercive (see [6, Proposition 3.5.4]), we have to impose that $\frac{\varphi(y)}{y}$ diverges as $|y| \rightarrow+\infty$. Such condition can be satisfied (by inspection of (7.8)) only if

$$
\begin{equation*}
e^{\varphi(0)}=M_{\infty}:=\int_{0}^{+\infty} \operatorname{sm}(s) d s \tag{7.9}
\end{equation*}
$$

By (7.8) and (7.9), it holds

$$
\lim _{y \rightarrow+\infty} \frac{\varphi(y)}{y}=\lim _{y \rightarrow+\infty} \int_{0}^{y} \frac{\log \left(M_{\infty}-M(t)\right)-\log M_{\infty}}{t^{2}} d t
$$

It is quite easy to construct explicit examples of positive continuous functions $m$, with finite integral and zero barycentre, such that limit at the r.h.s. of the above equality remains finite. For such a datum $m$, the Minkowski's problem does not admit solutions in $\mathcal{A}^{\prime}$.

In view of the above remark, and since in higher dimensions equality (7.1) does not correspond any longer to an o.d.e., but rather to a Monge-Ampère type equation, proving a general existence result for the functional Minkowski's problem seems to be a quite delicate task. On the other hand, as a consequence of Corollary 5.3, we are able to prove that uniqueness (up to translations) holds true, in both the cases of $\mathcal{A}^{\prime}$ and $\mathcal{A}^{\prime \prime}$.

Proposition 7.4. Let $f_{1}, f_{2} \in \mathcal{A}$ satisfy one of the following conditions:

$$
\begin{equation*}
f_{i} \in \mathcal{A}^{\prime} \quad i=1,2, \quad \text { and } \quad \mu\left(f_{1}\right)=\mu\left(f_{2}\right) \tag{7.10}
\end{equation*}
$$

or

$$
\begin{equation*}
f_{i} \in \mathcal{A}^{\prime \prime} \quad i=1,2, \quad \text { and } \quad \mu\left(f_{1}\right)=\mu\left(f_{2}\right), \quad \sigma\left(f_{1}\right)=\sigma\left(f_{2}\right) \tag{7.11}
\end{equation*}
$$

Then there exists $x_{0} \in \mathbb{R}^{n}$ such that $f_{2}(x)=f_{1}\left(x-x_{0}\right)$.
Proof. Firstly notice that the equality $\mu\left(f_{1}\right)=\mu\left(f_{2}\right)$ implies $J\left(f_{1}\right)=J\left(f_{2}\right)$. Moreover the assumption $f_{i} \in \mathcal{A}^{\prime}$ (or $f_{i} \in \mathcal{A}^{\prime \prime}$ ) implies that $J\left(f_{i}\right)>0$. If (7.10) holds, by Theorem 4.5 one has

$$
\delta J\left(f_{1}, g\right)=\delta J\left(f_{2}, g\right) \quad \forall g \in \mathcal{A}^{\prime \prime}
$$

In particular, taking $g=f_{1}$ or $g=f_{2}$, one sees that condition (5.3) is satisfied. Therefore, we are in a position to apply Corollary 5.3, and the statement follows. If (7.11) holds, the proof is exactly the same by using Theorem 4.6 in place of Theorem 4.5.

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## Appendix

This appendix contains the proofs of some results stated in Section 4, precisely all the preliminary lemmas used in the proof of Theorems 4.5 and 4.6 , and the claim made in Remark 4.7.

Proof of Lemma 4.9. It is immediate to check that the classes $\mathcal{L}^{\prime}$ and $\mathcal{L}^{\prime \prime}$ are closed by right multiplication by a positive scalar. Let us show that each of them is closed also by infimal convolution.
(i) Let $u, v \in \mathcal{L}^{\prime}$, set $\varphi:=u^{*}, \psi:=v^{*}$, and $w:=u \square v$.

By Proposition 2.1(iii), it holds $\operatorname{dom}(w)=\operatorname{dom}(u)+\operatorname{dom}(v)=\mathbb{R}^{n}$.
The condition of having a superlinear growth at infinity is equivalent to the condition of being cofinite [6, Proposition 3.5.4], and the latter is clearly closed by infimal convolution in view of the equality $w^{*}=\varphi+\psi$ holding by Proposition 2.1(iv). Therefore, $w$ has superlinear growth at infinity.

Since $\left(\mathbb{R}^{n}, u\right)$ and $\left(\mathbb{R}^{n}, v\right)$ are convex functions of Legendre type, with $u, v \in \mathcal{C}_{+}^{2}$, the mappings $\nabla u$ and $\nabla v$ are $\mathcal{C}^{1}$ bijections from $\mathbb{R}^{n}$ to $\mathbb{R}^{n}$, with a nonsingular Jacobian. Therefore also their inverse maps, which by Proposition 2.2 are precisely $\nabla \varphi$ and $\nabla \psi$, are $\mathcal{C}^{1}$ bijections from $\mathbb{R}^{n}$ to $\mathbb{R}^{n}$, and the same holds true for their sum. Hence $\left(\mathbb{R}^{n}, \varphi+\psi\right)$ is a convex function of Legendre type, with $\varphi+\psi$ of class $\mathcal{C}_{+}^{2}$. In turn, this implies that the Legendre conjugate of $\left(\mathbb{R}^{n}, \varphi+\psi\right)$, namely $\left(\mathbb{R}^{n}, w\right)$, is a convex function of Legendre type, with $w$ of class $\mathcal{C}_{+}^{2}$.
(ii) Let $u, v \in \mathcal{L}^{\prime \prime}$, and set $K:=\operatorname{dom}(u), L:=\operatorname{dom}(v), \varphi, \psi$, and $w$ as above.

By Proposition 2.1(iii), it holds $\operatorname{dom}(w)=K+L \in \mathcal{K}^{n} \cap \mathcal{C}_{+}^{2}$.
Since $u$ and $v$ are of class $\mathcal{C}_{+}^{2}$, and their gradients diverge at the boundary of their domains, (int $(K), u)$ and $(\operatorname{int}(L), v)$ are convex functions of Legendre type, and the mappings $\nabla u$ and $\nabla v$ are $\mathcal{C}^{1}$ bijections respectively from $K$ and $L$ onto $\mathbb{R}^{n}$. Hence, similarly as above, we may apply Proposition 2.2 to infer that $\left(\mathbb{R}^{n}, \varphi\right)$, $\left(\mathbb{R}^{n}, \psi\right)$, and hence $\left(\mathbb{R}^{n}, \varphi+\psi\right)$, are convex functions of Legendre type, with $\varphi+\psi$ of class $\mathcal{C}_{+}^{2}$. This yields that $\left(\mathbb{R}^{n}, w\right)$ is a convex function of Legendre type, with $w$ of class $\mathcal{C}_{+}^{2}$.

It remains to check that $w$ is continuous up to $\partial(K+L)$. To this end we are going to use as a crucial tool the identity

$$
\begin{align*}
u \square v(x) & =\inf _{x_{1}+x_{2}=x}\left\{u\left(x_{1}\right)+v\left(x_{2}\right)\right\} \\
& =u\left(v_{K}^{-1}\left(v_{K+L}(x)\right)\right)+v\left(v_{L}^{-1}\left(v_{K+L}(x)\right)\right) \quad \forall x \in \partial(K+L), \tag{A.1}
\end{align*}
$$

which follows from the definition of infimal convolution and the assumption $\partial K, \partial L \in \mathcal{C}_{+}^{2}$.
Let $\bar{x} \in \partial(K+L)$, and let us show that for every sequence of points $x^{h} \in K+L$ such that $x^{h} \rightarrow \bar{x}$, it holds

$$
\begin{equation*}
\lim _{h} u \square v\left(x^{h}\right)=u \square v(\bar{x}) . \tag{A.2}
\end{equation*}
$$

Up to passing to a (not relabelled) subsequence, we may assume that one of the following two cases occurs:

$$
x^{h} \in \partial(K+L) \forall h \quad \text { or } \quad x^{h} \in \operatorname{int}(K+L) \forall h .
$$

Consider first the case $x^{h} \in \partial(K+L) \forall h$. Let us write the identity (A.1) at $x^{h}$

$$
u \square v\left(x^{h}\right)=u\left(v_{K}^{-1}\left(v_{K+L}\left(x^{h}\right)\right)\right)+v\left(v_{L}^{-1}\left(v_{K+L}\left(x^{h}\right)\right)\right) \quad \forall h,
$$

and then let us pass to the limit in $h$. Since by hypothesis the Gauss maps $v_{K}, \nu_{L}$ and their inverse are continuous, and $u, v$ are continuous up to $\partial K, \partial L$, we get

$$
\lim _{h} u \square v\left(x^{h}\right)=u\left(v_{K}^{-1}\left(v_{K+L}(\bar{x})\right)\right)+v\left(v_{L}^{-1}\left(v_{K+L}(\bar{x})\right)\right) .
$$

In view of the identity (A.1), the r.h.s. of the above equality equals $u \square v(\bar{x})$, and (A.2) is proved.
Consider now the case $x^{h} \in \operatorname{int}(K+L) \forall h$. We set

$$
y^{h}:=\nabla w\left(x^{h}\right)=(\nabla(\varphi+\psi))^{-1}\left(x^{h}\right)
$$

and we decompose $x^{h}$ as $x_{1}^{h}+x_{2}^{h}$, with

$$
x_{1}^{h}:=\nabla \varphi\left(y^{h}\right) \in \operatorname{int}(K) \quad \text { and } \quad x_{2}^{h}:=\nabla \psi\left(y^{h}\right) \in \operatorname{int}(L) .
$$

Then we have

$$
u \square v\left(x^{h}\right)=\left[\left\langle x_{1}^{h}, y^{h}\right\rangle-\varphi\left(y^{h}\right)\right]+\left[\left\langle x_{2}^{h}, y^{h}\right\rangle-\psi\left(y^{h}\right)\right]=u\left(x_{1}^{h}\right)+v\left(x_{2}^{h}\right) .
$$

Let us now pass the limit in $h$. By compactness, after possibly selecting a (not relabelled) subsequence, there exist $\lim _{h} x_{1}^{h}=: \bar{x}_{1} \in \partial K$ and $\lim _{h} x_{2}^{h}=: \bar{x}_{2} \in \partial L$. Since by assumption $u \in \mathcal{C}^{0}(K)$ and $v \in \mathcal{C}^{0}(L)$, we infer

$$
\lim _{h} u \square v\left(x^{h}\right)=u\left(\bar{x}_{1}\right)+v\left(\bar{x}_{2}\right) .
$$

In view of the identity (A.1), the above equality implies (A.2) provided

$$
\bar{x}_{1}=v_{K}^{-1}\left(v_{K+L}(\bar{x})\right) \quad \text { and } \quad \bar{x}_{2}=v_{L}^{-1}\left(v_{K+L}(\bar{x})\right) .
$$

In turn, by the $\mathcal{C}_{+}^{2}$ assumption on $\partial K$, $\partial L$, such conditions are satisfied provided the normal vectors $v_{K}\left(\bar{x}_{1}\right)$ and $v_{L}\left(\bar{x}_{2}\right)$ coincide. Let us show that in fact each of them agrees with

$$
\bar{\xi}:=\lim _{h} \frac{y^{h}}{\left\|y^{h}\right\|}
$$

Since $y^{h}=\nabla u\left(x_{1}^{h}\right)$, and $\left\|y^{h}\right\| \rightarrow+\infty$ (being $y^{h}=\nabla w_{h}\left(x^{h}\right)$ and $x^{h} \rightarrow \bar{x} \in \partial(K+L)$ ), by passing to the limit in the inequality

$$
\frac{u(x)}{\left\|y^{h}\right\|} \geq \frac{u\left(x_{1}^{h}\right)}{\left\|y^{h}\right\|}+\left\langle\frac{y^{h}}{\left\|y^{h}\right\|}, x-x_{1}^{h}\right\rangle
$$

we infer that any cluster point of the sequence $y^{h} /\left\|y^{h}\right\|$ belongs to the normal cone to $\partial K$ at $\bar{x}_{1}$, which is reduced to $\nu_{K}\left(\bar{x}_{1}\right)$. In the same way we obtain $\bar{\xi}=\nu_{L}\left(\bar{x}_{2}\right)$, and the proof is achieved.

Proof of Lemma 4.10. (i) Let $x \in \operatorname{dom}(u)$ be fixed. By the assumption $v(0)=0$, we have $u_{t}(x) \leq u(x)$ for every $t>0$, so that $\lim \sup _{t \rightarrow 0^{+}} u_{t}(x) \leq u(x)$. Let us prove that we also have

$$
\begin{equation*}
\liminf _{t \rightarrow 0^{+}} u_{t}(x) \geq u(x) \tag{A.3}
\end{equation*}
$$

Assume $u, v \in \mathcal{L}^{\prime}$, and set $\varphi:=u^{*}, \psi:=v^{*}$. We choose $r>\|\nabla u(x)\|$ and we set $c:=\sup _{B_{r}} \psi$ (notice that $c$ is finite because $\psi$ is bounded on bounded sets [6, Theorem 4.4.13]). Then

$$
\begin{aligned}
u_{t}(x) & =\sup _{y \in \mathbb{R}^{n}}\{\langle x, y\rangle-\varphi(y)-t \psi(y)\} \geq \sup _{y \in B_{r}}\{\langle x, y\rangle-\varphi(y)\}-t c \\
& =\langle x, \nabla u(x)\rangle-\varphi(\nabla u(x))-t c=u(x)-t c
\end{aligned}
$$

and (A.3) follows by passing to the inferior limit as $t \rightarrow 0^{+}$.
Assume $u, v \in \mathcal{L}^{\prime \prime}$. Setting $L:=\operatorname{dom}(v)$ and $m:=\min v$, it holds $v \geq I_{L}+m$. Then

$$
\begin{aligned}
u_{t}(x) & =\inf _{x_{1}+x_{2}=x}\left\{u\left(x_{1}\right)+t v\left(x_{2} / t\right)\right\} \geq \inf _{x_{1}+x_{2}=x}\left\{u\left(x_{1}\right)+t I_{L}\left(x_{2} / t\right)\right\}+t m \\
& =\inf _{x_{1}+x_{2}=x}\left\{u\left(x_{1}\right)+t I_{t L}\left(x_{2}\right)\right\}+t m=\inf _{x_{1} \in K \cap(x-t L)}\left\{u\left(x_{1}\right)\right\}+t m
\end{aligned}
$$

and, thanks to the continuity of $u$ at $x$, (A.3) follows by passing to the inferior limit as $t \rightarrow 0^{+}$.
Statement (ii) is an immediate consequence of the convexity of the functions $u_{t}$ and of the differentiability of their pointwise limit $u$ in the interior of its domain.

Proof of Lemma 4.11. Set $K_{t}:=\operatorname{dom}\left(u_{t}\right)$. First we claim that, for every fixed $x \in \operatorname{int}\left(K_{t}\right)$,
the map $t \mapsto \nabla u_{t}(x)$ is differentiable on $(0,+\infty)$.

Indeed, as noticed in the proof of Lemma 4.9, the Fenchel conjugates $\varphi:=u^{*}$ and $\psi:=v^{*}$ are both of class $\mathcal{C}_{+}^{2}$ on $\mathbb{R}^{n}$. Therefore, the function $F: \mathbb{R}^{n} \times \mathbb{R}^{n} \times(0,+\infty) \rightarrow \mathbb{R}^{n}$ defined by

$$
F(x, y, t):=\nabla \varphi(y)+t \nabla \psi(y)-x,
$$

is of class $\mathcal{C}^{1}$ on $\mathbb{R}^{n} \times \mathbb{R}^{n} \times(0,+\infty)$, and $\frac{\partial F}{\partial y}=\nabla^{2} \varphi+t \nabla^{2} \psi$ is nonsingular for every $y \in \mathbb{R}^{n}$. Consequently, by the implicit function theorem, the equation $F(x, y, t)=0$ locally defines a map $y=y(x, t)$ which is of class $\mathcal{C}^{1}$ in its arguments. By Lemma 4.9, $\left(\operatorname{int}\left(K_{t}\right), u_{t}\right)$ is a convex function of Legendre type, hence by Proposition $2.2 \nabla u_{t}$ is the inverse map of $\nabla \varphi_{t}$, namely

$$
F\left(x, \nabla u_{t}(x), t\right)=\nabla \varphi_{t}\left(\nabla u_{t}(x)\right)-x=0 .
$$

Therefore, for every $x \in \operatorname{int}\left(K_{t}\right)$ and every $t>0, y(x, t)=\nabla u_{t}(x)$, and (A.4) is proved.
Next, we apply again to Proposition 2.2 in order to write the identity

$$
\begin{equation*}
u_{t}(x)=\left\langle x, \nabla u_{t}(x)\right\rangle-\varphi_{t}\left(\nabla u_{t}(x)\right) \quad \forall x \in \operatorname{int}\left(K_{t}\right) . \tag{A.5}
\end{equation*}
$$

By (A.4) and (A.5) we obtain that, for every fixed $x \in \operatorname{int}\left(K_{t}\right)$, the map $t \mapsto u_{t}(x)$ is differentiable on $(0,+\infty)$, with

$$
\begin{aligned}
\frac{d}{d t} u_{t}(x) & =\left\langle x, \frac{d}{d t}\left(\nabla u_{t}(x)\right)\right\rangle-\psi\left(\nabla u_{t}(x)\right)-\left\langle\nabla \varphi_{t}\left(\nabla u_{t}(x)\right), \frac{d}{d t}\left(\nabla u_{t}(x)\right)\right\rangle \\
& =-\psi\left(\nabla u_{t}(x)\right) .
\end{aligned}
$$

Proof of Lemma 4.12. We have

$$
\begin{aligned}
\int_{\mathbb{R}^{n}} \varphi(\nabla u(x)) f(x) d x & =\int_{\mathbb{R}^{n}}(\langle x, \nabla u\rangle-u) f d x \\
& =-\int_{\mathbb{R}^{n}}\langle x, \nabla f\rangle d x+\int_{\mathbb{R}^{n}} f \log f d x . \\
& =-\int_{\mathbb{R}^{n}} \operatorname{div}(f x) d x+n J(f)+\int_{\mathbb{R}^{n}} f \log f d x .
\end{aligned}
$$

We observe that

$$
\int_{\mathbb{R}^{n}} \operatorname{div}(f x) d x=\lim _{r \rightarrow+\infty} \int_{B_{r}} \operatorname{div}(f x) d x=\lim _{r \rightarrow+\infty} r \int_{\partial B_{r}} f d \mathcal{H}^{n-1}=0
$$

where the last equality holds true by Lemma 2.5. Therefore we have

$$
\int_{\mathbb{R}^{n}} \varphi(\nabla u(x)) f(x) d x=n J(f)+\int_{\mathbb{R}^{n}} f \log f d x
$$

and the lemma follows recalling that both $J(f)$ and $\int_{\mathbb{R}^{n}} f \log f d x$ are finite ( $c f$. respectively Lemma 2.5 and Proposition 3.11).

Proof of Lemma 4.13. By definition we have

$$
u_{t}(x)=\inf _{x_{1}+x_{2}=x}\left\{u\left(x_{1}\right)+t v\left(\frac{x_{2}}{t}\right)\right\} .
$$

Since

$$
v_{\min }+I_{L}(x) \leq v(x) \leq v_{\max }+I_{L}(x) \quad \forall x \in \mathbb{R}^{n},
$$

it holds

$$
\begin{aligned}
\inf _{x_{1}+x_{2}=x}\left\{u\left(x_{1}\right)+t v_{\min }+t I_{L}\left(\frac{x_{2}}{t}\right)\right\} & \leq u_{t}(x) \\
& \leq \inf _{x_{1}+x_{2}=x}\left\{u\left(x_{1}\right)+t v_{\max }+t I_{L}\left(\frac{x_{2}}{t}\right)\right\},
\end{aligned}
$$

namely

$$
t v_{\min }+\inf _{x_{1} \in K \cap(x-t L)}\left\{u\left(x_{1}\right)\right\} \leq u_{t}(x) \leq t v_{\max }+\inf _{x_{1} \in K \cap(x-t L)}\left\{u\left(x_{1}\right)\right\} .
$$

Therefore the statement is satisfied by taking $y$ as a point where $u$ attains its minimum on $K \cap(x-t L)$.

Proof of Remark 4.7. By inspection of the proof of Theorems 4.5 and 4.6 , one can see that assumption (4.1) is used only in Step 4 (in order to prove that $\Psi(0)<+\infty$ ) and in Step 5 (in order to prove that $\left.\lim _{s \rightarrow 0^{+}} \Psi(s)=\Psi(0)\right)$. Assume now $n=1$, and drop assumption (4.1): let us indicate how Steps 4 and 5 (and consequently also Step 6) have to be modified in order to show that (4.3) continues to hold, possibly as an equality $+\infty=+\infty$.

In Step 4, we limit ourselves to prove that $\Psi$ takes finite values at every $s>0$.
In Step 5, the proof of the continuity of $\Psi$ at every $s>0$ remains unchanged, whereas for $s \rightarrow 0^{+}$we make the following claim (whose proof is postponed below):

$$
\begin{equation*}
\text { if } \Psi(0)<+\infty \text {, then } \Psi \text { is continuous from the right at } s=0 \text {. } \tag{A.6}
\end{equation*}
$$

Consequently, in Step 6 we must distinguish two cases. In case $\Psi(0)<+\infty$, thanks to (A.6) equality (4.4) can be proved exactly as before. In case $\Psi(0)=+\infty$, (4.4) continues to hold as an equality $+\infty=+\infty$, and it can be proved by slight modifications of the case $\Psi(0)<\infty$. More precisely, (4.22) and (4.24) in Step 6 remain unchanged, whereas (4.23) has to be replaced by

$$
\begin{equation*}
\liminf _{s \rightarrow 0^{+}} \Psi(s) \geq \sup _{C \subset \subset E} \liminf _{s \rightarrow 0^{+}} \int_{C} \psi d \mu\left(f_{s}\right)=\sup _{C \subset \subset E} \int_{C} \psi d \mu(f)=+\infty \tag{A.7}
\end{equation*}
$$

(notice that the second equality in (A.7) holds by dominated convergence, since by Lemma 4.10 we have $\psi\left(\nabla u_{s}\right) f_{s} \rightarrow \psi(\nabla u) f$ as $s \rightarrow 0^{+}$, and by Lemmas 3.9 and 4.10(ii) the nonnegative functions $\psi\left(\nabla u_{s}\right) f_{s}$ are bounded above on $C$ by some continuous function independent of $s$ ).

Let us finally prove (A.6). Assume

$$
\begin{equation*}
\Psi(0)=\int_{\mathbb{R}} \psi \varphi^{\prime \prime} e^{\varphi-y \varphi^{\prime}} d y<+\infty \tag{A.8}
\end{equation*}
$$

Since $n=1$, (4.13) simplifies into

$$
\Psi(s)=\Psi_{0}(s)+s \Psi_{1}(s)
$$

where

$$
\begin{array}{ll}
\Psi_{0}(s):=\int_{\mathbb{R}} h_{0}(s, y) d y & h_{0}(s, y):=\psi e^{\varphi_{s}-y \varphi_{s}^{\prime}} \varphi^{\prime \prime} \chi_{Q_{s}} \\
\Psi_{1}(s):=\int_{\mathbb{R}} h_{1}(s, y) d y & h_{1}(s, y):=\psi e^{\varphi_{s}-y \varphi_{s}^{\prime}} \psi^{\prime \prime} \chi_{Q_{s}} .
\end{array}
$$

To get (A.6) it suffices to show that

$$
\begin{align*}
& \lim _{s \rightarrow 0^{+}} \Psi_{0}(s)=\Psi(0)  \tag{A.9}\\
& \lim _{s \rightarrow 0^{+}} s \Psi_{1}(s)=0 \tag{A.10}
\end{align*}
$$

Thanks to assumption (A.8), (A.9) can be proved exactly as before (cf. the proof of (4.16)). To prove (A.10), we write

$$
s \Psi_{1}(s)=I_{+}(s)+I_{-}(s):=\int_{\mathbb{R}_{+}} s h_{1}(s, y) d y+\int_{\mathbb{R}_{-}} s h_{1}(s, y) d y
$$

and we show that both $I_{ \pm}(s)$ tend to 0 as $s \rightarrow 0^{+}$. Let us consider $I_{+}(s)$ (the case of $I_{-}(s)$ is completely analogous).

We observe that

$$
0 \leq s h_{1}(s, y)=s \psi e^{\varphi_{s}-y \varphi_{s}^{\prime}} \psi^{\prime \prime} \chi_{Q_{s}} \leq-F(y) G_{s}^{\prime}(y) \quad \forall y, s>0
$$

where we have set

$$
F(y):=\frac{\psi}{y} e^{\varphi-y \varphi^{\prime}} \quad \text { and } \quad G_{s}(y):=e^{s\left(\psi-y \psi^{\prime}\right)}
$$

Then an integration by parts gives

$$
0 \leq I_{+}(s) \leq \lim _{\varepsilon \rightarrow 0^{+}, r \rightarrow+\infty}\left\{\int_{\varepsilon}^{r} F^{\prime}(y) G_{s}(y) d y+F(\varepsilon) G_{s}(\varepsilon)-F(r) G_{s}(r)\right\}
$$

Since $\psi(0)=\psi^{\prime}(0)=0$ (respectively because $v \geq 0$ and $\psi \geq 0$ ), passing to the limit in $\varepsilon$ gives

$$
\begin{equation*}
0 \leq I_{+}(s) \leq \lim _{r \rightarrow+\infty}\left\{\int_{0}^{r} F^{\prime}(y) G_{s}(y) d y-F(r) G_{s}(r)\right\} \tag{A.11}
\end{equation*}
$$

Next we observe that the following limit exists:

$$
\alpha:=\lim _{r \rightarrow+\infty} F(r)=\lim _{r \rightarrow+\infty} \int_{0}^{r} F^{\prime}(y) d y
$$

Indeed a straightforward computation gives

$$
F^{\prime}(y)=-\psi \varphi^{\prime \prime} e^{\varphi-y \varphi^{\prime}}+\frac{e^{\varphi-y \varphi^{\prime}}}{y^{2}}\left(\psi^{\prime} y-\psi\right)
$$

and both the functions at the right and side are integrable on $(0,+\infty)$ (the former by assumption (A.8), the latter because it is nonnegative).

Let us show that $\alpha>0$ cannot occur. Indeed in such case, for some constants $\bar{c}$ and $\bar{r}$, it would be $F(r) \geq \bar{c} \forall r \geq \bar{r}$. This would contradict (A.8), since

$$
\begin{aligned}
\Psi(0) & \geq \int_{\bar{r}}^{+\infty} \psi \varphi^{\prime \prime} e^{\varphi-y \varphi^{\prime}} d y \geq \bar{c} \int_{\bar{r}}^{+\infty} y \varphi^{\prime \prime} d y \\
& =\bar{c}\left\{\lim _{r \rightarrow+\infty}\left[r \varphi^{\prime}(r)-\varphi(r)\right]-\left[\bar{r} \varphi^{\prime}(\bar{r})-\varphi(\bar{r})\right]\right\}=+\infty .
\end{aligned}
$$

Taking into account that $\alpha=0$ (and also that $\lim _{r \rightarrow+\infty} G_{s}(r)=0$ ), we may rewrite (A.11) as

$$
\begin{equation*}
0 \leq I_{+}(s) \leq \lim _{r \rightarrow+\infty} \int_{0}^{r} F^{\prime}(y) G_{s}(y) d y \tag{A.12}
\end{equation*}
$$

Moreover, since $\alpha=0$, we have in particular $\int_{0}^{+\infty} F^{\prime}(y) d y<+\infty$, which implies $F^{\prime} \in$ $L^{1}(0,+\infty)$. Therefore, for every fixed $s>0$, the functions $F^{\prime} G_{s}$ satisfy

$$
\left|F^{\prime}(y) G_{s}(y)\right| \leq\left|F^{\prime}(y)\right| \in L^{1}(0,+\infty)
$$

We deduce that (A.12) can be rewritten as

$$
\begin{equation*}
0 \leq I_{+}(s) \leq \int_{0}^{+\infty} F^{\prime}(y) G_{s}(y) d y \tag{A.13}
\end{equation*}
$$

Finally, passing to the limit as $s \rightarrow 0^{+}$in the right hand side of (A.13) we obtain

$$
\lim _{s \rightarrow 0^{+}} \int_{0}^{+\infty} F^{\prime}(y) G_{s}(y) d y=\int_{0}^{+\infty} \lim _{s \rightarrow 0^{+}} F^{\prime}(y) G_{s}(y) d y=\int_{0}^{+\infty} F^{\prime}(y) d y=\alpha=0
$$

This implies that $I_{+}(s)$ tends to 0 as $s \rightarrow 0^{+}$and the proof is achieved.

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[^0]:    * Corresponding author.

    E-mail address: colesant@math.unifi.it (A. Colesanti).

